# Introduction to Microlocal Analysis 

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## Preface

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## Introduction

I shall assume some familiarity with distribution theory, with basic analysis knowledge of the theory of manifolds would also be useful. Any one or two of these prerequisites can be easily picked up along the way, but the prospective student with none of them should perhaps do some preliminary reading: ${ }_{\text {Friedlander2 }}$

Distributions: I good introduction is Friedlander's book firm For a more exhaustive treatment see Volume I of Hörmander's treatise $\frac{8}{\mathbf{8}] .}$

Analysis on manifolds: Most of what we need can be picked up from Munkres' book $\left[\frac{M u n k r e s 1}{[y]}\right.$ or Spivak's little book [12].

## CHAPTER 1

## Tempered

Tempered distributions and the Fourier transform

Microlocal analysis is a geometric theory of distributions, or a theory of geometric distributions. Rather than study general distributions - which are like general continuous functions but worse - we consider more specific types of distributions which actually arise in the study of differential and integral equations. Distributions are usually defined by duality, starting from very "good" test functions; correspondingly a general distribution is everywhere "bad". The conormal distributions we shall study implicitly for a long time, and eventually explicitly, are usually good, but like (other) people have a few interesting faults, i.e. singularities. These singularities are our principal target of study. Nevertheless we need the general framework of distribution theory to work in, so I will start with a brief introduction. This is designed either to remind you of what you already know or else to send you off to work it out. ${ }^{1}$ Proofs of some of the main theorems are outlined in the problems at the end of the chapter.

### 1.1. Schwartz test functions

To fix matters at the beginning we shall work in the space of tempered distributions. These are defined by duality from the space of Schwartz functions, also called the space of test functions of rapid decrease. We can think of analysis as starting off from algebra, which gives us the polynomials. Thus in $\mathbb{R}^{n}$ we have the coordinate functions, $x_{1}, \ldots, x_{n}$ and the constant functions and then the polynomials are obtained by taking (finite) sums and products:

$$
\begin{align*}
\phi(x)=\sum_{|\alpha| \leq k} p_{\alpha} x^{\alpha}, p_{\alpha} \in \mathbb{C}, \alpha \in \mathbb{N}_{0}^{n}, \alpha=\left(\alpha_{1}, \ldots, \alpha_{n}\right)  \tag{1.1}\\
\text { where } x^{\alpha}=x_{1}^{\alpha_{1}} \ldots x_{n}^{\alpha_{n}}=\prod_{j=1}^{n} x_{j}^{\alpha_{j}} \text { and } \mathbb{N}_{0}=\{0,1,2, \ldots\} .
\end{align*}
$$

A general function $\phi: \mathbb{R}^{n} \longrightarrow \mathbb{C}$ is differentiable at $\bar{x}$ if there is a linear function $\ell_{\bar{x}}(x)=c+\sum_{j=1}^{n}\left(x_{j}-\bar{x}_{j}\right) d_{j}$ such that for every $\epsilon>0$ there exists $\delta>0$ such that

$$
\begin{equation*}
\left|\phi(x)-\ell_{\bar{x}}(x)\right| \leq \epsilon|x-\bar{x}| \quad \forall|x-\bar{x}|<\bar{\delta} . \tag{1.2}
\end{equation*}
$$

The coefficients $d_{j}$ are the partial derivative of $\phi$ at the point $\bar{x}$. Then, $\phi$ is said to be differentiable on $\mathbb{R}^{n}$ if it is differentiable at each point $\bar{x} \in \mathbb{R}^{n}$; the partial derivatives are then also functions on $\mathbb{R}^{n}$ and $\phi$ is twice differentiable if the partial

[^0]derivatives are differentiable. In general it is $k$ times differentiable if its partial derivatives are $k-1$ times differentiable.

If $\phi$ is $k$ times differentiable then, for each $\bar{x} \in \mathbb{R}^{n}$, there is a polynomial of degree $k$,

$$
p_{k}(x ; \bar{x})=\sum_{|\alpha| \leq k} a_{\alpha} i^{|\alpha|}(x-\bar{x})^{\alpha} / \alpha!,|\alpha|=\alpha_{1}+\cdots+\alpha_{n}
$$

such that for each $\epsilon>0$ there exists $\delta>0$ such that

$$
\left|\phi(x)-p_{k}(x, \bar{x})\right| \leq \epsilon|x-\bar{x}|^{k} \quad \text { if }|x-\bar{x}|<\delta
$$

Then we set

$$
\begin{equation*}
D^{\alpha} \phi(\bar{x})=a_{\alpha} . \tag{tabular}
\end{equation*}
$$

If $\phi$ is infinitely differentiable all the $D^{\alpha} \phi$ are infinitely differentiable (hence continuous!) functions.
1.5 Definition 1.1. The space of Schwartz test functions of rapid decrease consists of those $\phi: \mathbb{R}^{n} \longrightarrow \mathbb{C}$ such that for every $\alpha, \beta \in \mathbb{N}_{0}^{n}$

$$
\begin{equation*}
\sup _{x \in \mathbb{R}^{n}}\left|x^{\beta} D^{\alpha} \phi(x)\right|<\infty ; \tag{tabular}
\end{equation*}
$$

it is denoted $\mathcal{S}\left(\mathbb{R}^{n}\right)$.
From $\left(\frac{1.6}{1.5}\right)$ we construct norms on $\mathcal{S}\left(\mathbb{R}^{n}\right)$ :

$$
\begin{equation*}
\|\phi\|_{k}=\max _{|\alpha|+|\beta| \leq k} \sup _{x \in \mathbb{R}^{n}}\left|x^{\alpha} D^{\beta} \phi(x)\right| . \tag{1.6}
\end{equation*}
$$

It is straightforward to check the conditions for a norm:
(1) $\|\phi\|_{k} \geq 0,\|\phi\|_{k}=0 \Longleftrightarrow \phi \equiv 0$
(2) $\|t \phi\|_{k}=|t|\|\phi\|_{k}, t \in \mathbb{C}$
(3) $\|\phi+\psi\|_{k} \leq\|\phi\|_{k}+\|\psi\|_{k} \forall \phi, \psi \in \mathcal{S}\left(\mathbb{R}^{n}\right)$.

The topology on $\mathcal{S}\left(\mathbb{R}^{n}\right)$ is given by the metric

$$
\begin{equation*}
d(\phi, \psi)=\sum_{k} 2^{-k} \frac{\|\phi-\psi\|_{k}}{1+\|\phi-\psi\|_{k}} \tag{1.7}
\end{equation*}
$$

31.1 .2000 .260

See Problem $\frac{1.4 .}{11.4}$
1.9 Proposition 1.1. With the distance function $\left(\frac{1.8}{1.7}\right), \mathcal{S}\left(\mathbb{R}^{n}\right)$ becomes a complete metric space (in fact it is a Fréchet space).

Of course one needs to check that $\mathcal{S}\left(\mathbb{R}^{n}\right)$ is non-trivial; however one can easily see that

$$
\begin{equation*}
\exp \left(-|x|^{2}\right) \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{1.8}
\end{equation*}
$$

In fact there are lots of smooth functions of compact support and
1.2 .2000 .266

$$
\begin{equation*}
\mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)=\left\{u \in \mathcal{S}\left(\mathbb{R}^{n}\right) ; u=0 \text { in }|x|>R=R(u)\right\} \subset \mathcal{S}\left(\mathbb{R}^{n}\right) \text { is dense. } \tag{1.9}
\end{equation*}
$$

The two elementary operations of differentiation and coordinate multiplication give continuous linear operators:
31.1 .2000 .263

$$
\begin{aligned}
& x_{j}: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}\left(\mathbb{R}^{n}\right) \\
& D_{j}: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}\left(\mathbb{R}^{n}\right)
\end{aligned}
$$

Other important operations we shall encounter include the exterior product,

$$
\begin{gather*}
\mathcal{S}\left(\mathbb{R}^{n}\right) \times \mathcal{S}\left(\mathbb{R}^{m}\right) \ni(\phi, \psi) \mapsto \phi \boxtimes \psi \in \mathcal{S}\left(\mathbb{R}^{n+m}\right)  \tag{1.11}\\
\phi \boxtimes \psi(x, y)=\phi(x) \psi(y) .
\end{gather*}
$$

and pull-back or restriction. If $\mathbb{R}^{k} \subset \mathbb{R}^{n}$ is identified as the subspace $x_{j}=0, j>k$, then the restriction map

$$
\begin{equation*}
L: \mathbb{R}^{n} \longrightarrow \mathbb{R}^{n},(L x)_{j}=\sum_{k=1}^{n} L_{j k} x_{k} \tag{1.13}
\end{equation*}
$$

The Lie group of invertible linear transformations, $\mathrm{GL}(n, \mathbb{R})$ is fixed by several equivalent conditions

$$
\begin{array}{r}
L \in \operatorname{GL}(n, \mathbb{R}) \Longleftrightarrow \operatorname{det}(L) \neq 0 \\
\Longleftrightarrow \exists L^{-1} \text { s.t. }\left(L^{-1}\right) L x=x \forall x \in \mathbb{R}^{n}  \tag{1.14}\\
\Longleftrightarrow \exists c>0 \text { s.t. } c|x| \leq|L x| \leq c^{-1}|x| \forall x \in \mathbb{R}^{n} .
\end{array}
$$

Pull-back of functions is defined by

$$
L^{*} \phi(x)=\phi(L x)=(\phi \circ L)(x)
$$

The chain rule for differentiation shows that if $\phi$ is diffferentiable then ${ }^{3}$
1.2 .2000 .269

$$
\begin{align*}
& D_{j} L^{*} \phi(x)=D_{j} \phi(L x)=\sum_{k=1}^{n} L_{k j}\left(D_{k} \phi\right)(L x)=L^{*}\left(\left(L_{*} D_{j}\right) \phi\right)(x)  \tag{1.15}\\
& L_{*} D_{j}=\sum_{k=1}^{n} L_{k j} D_{k}
\end{align*}
$$

From this it follows that
1.2 .2000 .270

$$
\begin{equation*}
L^{*}: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}\left(\mathbb{R}^{n}\right) \text { is an isomorphism for } L \in \mathrm{GL}(n, \mathbb{R}) \tag{1.16}
\end{equation*}
$$

To characterize the action of $L \in \mathrm{GL}(n, \mathbb{R})$ on $\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ consider, as usual, the distribution associated to $L^{*} \phi$ :
1.2 .2000 .271

$$
\begin{equation*}
\pi_{k}^{*}: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}\left(\mathbb{R}^{k}\right), \pi_{k}^{*} f(y)=f\left(y_{1}, \ldots, y_{k}, 0, \ldots, 0\right) \tag{1.12}
\end{equation*}
$$

is continuous (and surjective).

A linear transformation acts on $\mathbb{R}^{n}$ as a matrix ${ }^{2}$

Pull-ack of funtions is

$$
\begin{aligned}
T_{L^{*} \phi}(\psi)=\int_{\mathbb{R}^{n}} & \phi(L x) \psi(x) d x \\
& =\int_{\mathbb{R}^{n}} \phi(y) \psi\left(L^{-1} y\right)|\operatorname{det} L|^{-1} d y=T_{\phi}\left(|\operatorname{det} L|^{-1}\left(L^{-1}\right)^{*} \psi\right)
\end{aligned}
$$

[^1]Since the operator $|\operatorname{det} L|^{-1}\left(L^{-1}\right)^{*}$ is an ismorphism of $\mathcal{S}\left(\mathbb{R}^{n}\right)$ it follows that if we take the definition by duality

$$
\begin{align*}
& L^{*} u(\psi)=u\left(|\operatorname{det} L|^{-1}\left(L^{-1}\right)^{*} \psi\right), u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right), \psi \in \mathcal{S}\left(\mathbb{R}^{n}\right), L \in \operatorname{GL}(n, \mathbb{R})  \tag{1.18}\\
& \Longrightarrow L^{*}: \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)
\end{align*}
$$

is an isomorphism which extends $\left(\frac{1.2 .2000 .270}{1.16)}\right.$ and satisfies

$$
\begin{equation*}
D_{j} L^{*} u=L^{*}\left(\left(L_{*} D_{j}\right) u\right), L^{*}\left(x_{j} u\right)=\left(L^{*} x_{j}\right)\left(L^{*} u\right), u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right), L \in \operatorname{GL}(n, \mathbb{R}) \tag{1.19}
\end{equation*}
$$ as in (11.2.2000.269

### 1.3. Tempered distributions

As well as exterior multiplication $\left(\frac{1.10}{1.11}\right)$ there is the even more obvious multiplication operation

$$
\begin{gather*}
\mathcal{S}\left(\mathbb{R}^{n}\right) \times \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}\left(\mathbb{R}^{n}\right) \\
(\phi, \psi) \mapsto \phi(x) \psi(x) \tag{1.20}
\end{gather*}
$$

which turns $\mathcal{S}\left(\mathbb{R}^{n}\right)$ into a commutative algebra without identity. There is also integration

$$
\begin{equation*}
\int: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathbb{C} \tag{1.21}
\end{equation*}
$$

Combining these gives a pairing, a bilinear map

$$
\begin{equation*}
\mathcal{S}\left(\mathbb{R}^{n}\right) \times \mathcal{S}\left(\mathbb{R}^{n}\right) \ni(\phi, \psi) \longmapsto \int_{\mathbb{R}^{n}} \phi(x) \psi(x) d x \tag{1.22}
\end{equation*}
$$

If we fix $\phi \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ this defines a continuous linear map:

$$
\begin{equation*}
T_{\phi}: \mathcal{S}\left(\mathbb{R}^{n}\right) \ni \psi \longmapsto \int \phi(x) \psi(x) d x \tag{1.23}
\end{equation*}
$$

Continuity becomes the condition:

$$
\begin{equation*}
\exists k, C_{k} \text { s.t. }\left|T_{\phi}(\psi)\right| \leq C_{k}\|\psi\|_{k} \forall \psi \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{1.24}
\end{equation*}
$$

We generalize this by denoting by $\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ the dual space, i.e. the space of all continuous linear functionals

$$
u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \Longleftrightarrow u: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathbb{C}
$$

$\exists k, C_{k}$ such that $|u(\psi)| \leq C_{k}\|\psi\|_{k} \forall \psi \in \mathcal{S}\left(\mathbb{R}^{n}\right)$.
Lemma 1.1. The map

$$
\begin{equation*}
\mathcal{S}\left(\mathbb{R}^{n}\right) \ni \phi \longmapsto T_{\phi} \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \tag{1.25}
\end{equation*}
$$

is an injection.
Proof. For any $\phi \in \mathcal{S}\left(\mathbb{R}^{n}\right), T_{\phi}(\phi)=\int|\phi(x)|^{2} d x$, so $T_{\phi}=0$ implies $\phi \equiv 0$.
If we wish to consider a topology on $\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ it will normally be the weak topology, that is the weakest topology with respect to which all the linear maps

$$
\begin{equation*}
\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \ni u \longmapsto u(\phi) \in \mathbb{C}, \quad \phi \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{1.26}
\end{equation*}
$$

are continuous. This just means that it is given by the seminorms

$$
\begin{equation*}
\mathcal{S}\left(\mathbb{R}^{n}\right) \ni u \longmapsto|u(\phi)| \in \mathbb{R} \tag{1.27}
\end{equation*}
$$

where $\phi \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ is fixed but arbitrary. The sets

$$
\begin{equation*}
\left\{u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) ;\left|u\left(\phi_{j}\right)\right|<\epsilon_{j}, \phi_{j} \in \Phi\right\} \tag{1.28}
\end{equation*}
$$

form a basis of the neighbourhoods of 0 as $\Phi \subset \mathcal{S}\left(\mathbb{R}^{n}\right)$ runs over finite sets and the $\epsilon_{j}$ are positive numbers.

### 1.17

Proposition 1.2. The continuous injection $\mathcal{S}\left(\mathbb{R}^{n}\right) \hookrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$, given by $\left(\frac{1.15}{1.25}\right)$, has dense range in the weak topology.
See Problem $\mathbb{P 1 . 4}_{1.8}$ for the outline of a proof.
The maps $x_{i}, D_{j}$ extend by continuity (and hence uniquely) to operators
1.18

$$
\begin{equation*}
x_{j}, D_{j}: \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \tag{1.29}
\end{equation*}
$$

This is easily seen by defining them by duality. Thus if $\phi \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ set $D_{j} T_{\phi}=T_{D_{j} \phi}$, then

$$
\begin{equation*}
T_{D_{j} \phi}(\psi)=\int D_{j} \phi \psi=-\int \phi D_{j} \psi \tag{1.30}
\end{equation*}
$$

the integration by parts formula. The definitions

$$
\begin{equation*}
D_{j} u(\psi)=u\left(-D_{j} \psi\right), x_{j} u(\psi)=u\left(x_{j} \psi\right), \quad u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right), \psi \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{1.31}
\end{equation*}
$$

satisfy all requirements, in that they give continuous maps $\left(\frac{1.18}{1.29}\right)$ which extend the standard definitions on $\mathcal{S}\left(\mathbb{R}^{n}\right)$.

### 1.4. Two big theorems

The association, by $\left(\frac{1.15}{1.25}\right)$, of a distribution to a function can be extended considerably. For example if $u: \mathbb{R}^{n} \longrightarrow \mathbb{C}$ is a bounded and continuous function then

$$
\begin{equation*}
T_{u}(\psi)=\int u(x) \psi(x) d x \tag{1.32}
\end{equation*}
$$

still defines a distribution which vanishes if and only if $u$ vanishes identically. Using the operations ( 1.189$)$ we conclude that for any $\alpha, \beta \in \mathbb{N}_{0}^{n}$

$$
\begin{equation*}
x^{\beta} D_{x}^{\alpha} u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \text { if } u: \mathbb{R}^{n} \longrightarrow \mathbb{C} \text { is bounded and continuous. } \tag{1.33}
\end{equation*}
$$

Conversely we have the Schwartz representation Theorem:
1.19 Theorem 1.1. For any $u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ there is a finite collection $u_{\alpha \beta}: \mathbb{R}^{n} \longrightarrow \mathbb{C}$ of bounded continuous functions, $|\alpha|+|\beta| \leq k$, such that

$$
\begin{equation*}
u=\sum_{|\alpha|+|\beta| \leq k} x^{\beta} D_{x}^{\alpha} u_{\alpha \beta} . \tag{1.34}
\end{equation*}
$$

Thus tempered distributions are just products of polynomials and derivatives of bounded continuous functions. This is important because it says that distributions are "not too bad".

The second important result (long considered very difficult to prove, but there is a relatively straightforward proof using the Fourier transform) is the the Schwartz kernel theorem. To show this we need to use the exterior product (1.11). If $K \in$ $\mathcal{S}^{\prime}\left(\mathbb{R}^{n+m}\right)$ this allows us to define a linear map

$$
\begin{equation*}
O_{K}: \mathcal{S}\left(\mathbb{R}^{m}\right) \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \tag{1.35}
\end{equation*}
$$

by

$$
\begin{equation*}
O_{K}(\psi)(\phi)=\int K \cdot \phi \boxtimes \psi d x d y \tag{1.36}
\end{equation*}
$$

THEOREM 1.2. There is a 1-1 correspondence between continuous linear operators

$$
\begin{equation*}
A: \mathcal{S}\left(\mathbb{R}^{m}\right) \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \tag{1.37}
\end{equation*}
$$

and $\mathcal{S}^{\prime}\left(\mathbb{R}^{n+m}\right)$ given by $A=O_{K}$.
Brief outlines of the proofs of these two results can be found in Problems ${ }^{\mathbb{P} 2.15}$ and 17.16.

### 1.5. Examples

Amongst tempered distributions we think of $\mathcal{S}\left(\mathbb{R}^{n}\right)$ as being the 'trivial' examples, since they are the test functions. One can say that the study of the singularities of tempered distributions amounts to the study of the quotient

$$
\begin{equation*}
\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) / \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{1.38}
\end{equation*}
$$

which could, reasonably, be called the space of tempered microfunctions.
The sort of distributions we are interested in are those like the Dirac delta "function"

$$
\begin{equation*}
\delta(x) \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right), \quad \delta(\phi)=\phi(0) \tag{1.39}
\end{equation*}
$$

The definition here shows that $\delta$ is just the Schwartz kernel of the operator

$$
\begin{equation*}
\mathcal{S}\left(\mathbb{R}^{n}\right) \ni \phi \longmapsto \phi(0) \in \mathbb{C}=\mathcal{S}\left(\mathbb{R}^{0}\right) \tag{1.40}
\end{equation*}
$$

This is precisely one reason it is interesting. More generally we can consider the maps

$$
\begin{equation*}
\mathcal{S}\left(\mathbb{R}^{n}\right) \ni \phi \longmapsto D^{\alpha} \phi(0), \quad \alpha \in \mathbb{N}_{0}^{n} . \tag{1.41}
\end{equation*}
$$

These have Schwartz kernels $(-D)^{\alpha} \delta$ since

$$
\begin{equation*}
(-D)^{\alpha} \delta(\phi)=\delta\left(D^{\alpha} \phi\right)=D^{\alpha} \phi(0) \tag{1.42}
\end{equation*}
$$

If we write the relationship $A=O_{K} \longleftrightarrow K$ as

$$
\begin{equation*}
(A \psi)(\phi)=\int K(x, y) \phi(x) \psi(y) d x d y \tag{1.43}
\end{equation*}
$$

then $\left(\frac{1.21}{1.42}\right)$ becomes

$$
\begin{equation*}
D^{\alpha} \phi(0)=\int(-D)^{\alpha} \delta(x) \phi(x) d x \tag{1.44}
\end{equation*}
$$

More generally, if $K(x, y)$ is the kernel of an operator $A$ then the kernel of $A \cdot D^{\alpha}$ is $(-D)_{y}^{\alpha} K(x, y)$ whereas the kernel of $D^{\alpha} \circ A$ is $D_{x}^{\alpha} K(x, y)$.

### 1.6. Two little lemmas

Above, some of the basic properties of tempered distributions have been outlined. The main "raison d'être" for $\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ is the Fourier transform which we proceed to discuss. We shall use the Fourier transform as an almost indispensable tool in the treatment of pseudodifferential operators. The description of differential operators, via their Schwartz kernels, using the Fourier transform is an essential motivation for the extension to pseudodifferential operators.

Partly as simple exercises in the theory of distributions, and more significantly as preparation for the proof of the inversion formula for the Fourier transform we consider two lemmas.

First recall that if $u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ then we have defined $D_{j} u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ by

$$
\begin{equation*}
D_{j} u(\phi)=u\left(-D_{j} \phi\right) \quad \forall \phi \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{1.45}
\end{equation*}
$$

In this sense it is a "weak derivative". Let us consider the simple question of the form of the solutions to

$$
\begin{equation*}
D_{j} u=0, u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \tag{1.46}
\end{equation*}
$$

Let $I_{j}$ be the integration operator:

$$
\begin{align*}
& I_{j}: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}\left(\mathbb{R}^{n-1}\right) \\
& I_{j}(\phi)\left(y_{1}, \ldots, y_{n-1}\right)=\int \phi\left(y_{1}, \ldots y_{j-1}, x, y_{j}, \ldots y_{n-1}\right) d x \tag{1.47}
\end{align*}
$$

Then if $\pi_{j}: \mathbb{R}^{n} \longrightarrow \mathbb{R}^{n-1}$ is the map $\pi_{j}(x)=\left(x_{1}, \ldots, x_{j-1}, x_{j+1} \ldots, x_{n}\right)$, we define, for $v \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n-1}\right)$,

$$
\begin{equation*}
\pi_{j}^{*} v(\phi)=v\left(I_{j} \phi\right) \forall \phi \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{1.48}
\end{equation*}
$$

It is clear from $\left(\frac{2.2}{1.47}\right)$ that $I_{j}: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}\left(\mathbb{R}^{n-1}\right)$ is continuous and hence $\pi_{j}^{*} v \in$ $\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ is well-defined for each $v \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n-1}\right)$.
2.3 Lemma 1.2. The equation (12.1 1.46 ) holds if and only if $u=\pi_{j}^{*} v$ for some $v \in$ $\mathcal{S}^{\prime}\left(\mathbb{R}^{n-1}\right)$.

Proof. If $\phi \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ and $\phi=D_{j} \psi$ with $\psi \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ then $I_{j} \phi=I_{j}\left(D_{j} \psi\right)=0$. Thus if $u=\pi_{j}^{*} v$ then

$$
\begin{equation*}
u\left(-D_{j} \phi\right)=\pi_{j}^{*} v\left(-D_{j} \phi\right)=v\left(I_{j}\left(-D_{j} \phi\right)\right)=0 \tag{1.49}
\end{equation*}
$$

Thus $u=\pi_{j}^{*} v$ does always satisfy (1.1.46).
Conversely suppose (11.46) holds. Choose $\rho \in \mathcal{S}(\mathbb{R})$ with the property

$$
\begin{equation*}
\int \rho(x) d x=1 \tag{1.50}
\end{equation*}
$$

Then each $\phi \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ can be decomposed as

$$
\begin{equation*}
\phi(x)=\rho\left(x_{j}\right) I_{j} \phi\left(x_{1}, \ldots, x_{j-1}, x_{j+1}, \ldots x_{n}\right)+D_{j} \psi, \quad \psi \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{1.51}
\end{equation*}
$$

Indeed this is just the statement

$$
\begin{gathered}
\zeta \in \mathcal{S}\left(\mathbb{R}^{n}\right), I_{j} \zeta=0 \Longrightarrow \psi(x) \in \mathcal{S}\left(\mathbb{R}^{n}\right) \text { where } \\
\begin{aligned}
\psi(x)= & \int_{-\infty}^{x_{j}} \zeta\left(x_{1}, \ldots, x_{j-1}, t, x_{j+1}, \ldots, x_{n}\right) d t \\
& =\int_{\infty}^{x_{j}} \zeta\left(x_{1}, \ldots, x_{j-1}, t, x_{j+1}, \ldots, x_{n}\right) d t
\end{aligned}
\end{gathered}
$$

Using ( $\frac{2.4}{1.51)}$ and (2.1 1.46$)$ we have

$$
\begin{equation*}
u(\phi)=u\left(\rho\left(x_{j}\right) I_{j} \phi\left(x_{1}, \ldots, x_{j-1}, x_{j+1}, \ldots x_{n}\right)\right) \tag{1.52}
\end{equation*}
$$

Thus if

$$
\begin{equation*}
v(\psi)=u\left(\rho\left(x_{j}\right) \psi\left(x_{1}, \ldots, x_{j-1}, x_{j+1}, \ldots x_{n}\right)\right) \forall \psi \in \mathcal{S}\left(\mathbb{R}^{n-1}\right) \tag{1.53}
\end{equation*}
$$

then $v \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n-1}\right)$ and $u=\pi_{j}^{*} v$. This proves the lemma.
Of course the notation $u=\pi_{j}^{*} v$ is much too heavy-handed. We just write $u(x)=v\left(x_{1}, \ldots, x_{j-1}, x_{j+1}, \ldots, x_{n}\right)$ and regard ' $v$ as a distribution in one additional variable'.

The second, related, lemma is just a special case of a general result of Schwartz concerning the support of a distribution. The particular result is:
2.5 Lemma 1.3. Suppose $u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ and $x_{j} u=0, j=1, \ldots n$ then $u=c \delta(x)$ for some constant $c$.

Proof. Again we use the definition of multiplication and a dual result for test functions. Namely, choose $\rho \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ with $\rho(x)=1$ in $|x|<\frac{1}{2}, \rho(x)=0$ in $|x| \geq 3 / 4$. Then any $\phi \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ can be written

$$
\begin{equation*}
\phi=\phi(0) \cdot \rho(x)+\sum_{j=1}^{n} x_{j} \psi_{j}(x), \quad \psi_{j} \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{1.54}
\end{equation*}
$$

This in turn can be proved using Taylor's formula as I proceed to show. Thus

$$
\begin{equation*}
\phi(x)=\phi(0)+\sum_{j=1}^{n} x_{j} \zeta_{j}(x) \text { in }|x| \leq 1, \text { with } \zeta_{j} \in \mathcal{C}^{\infty} \tag{1.55}
\end{equation*}
$$

Then,

$$
\begin{equation*}
\rho(x) \phi(x)=\phi(0) \rho(x)+\sum_{j=1}^{n} x_{j} \rho \zeta_{j}(x) \tag{1.56}
\end{equation*}
$$

and $\rho \zeta_{j} \in \mathcal{S}\left(\mathbb{R}^{n}\right)$. Thus it suffices to check $\left(\frac{2.6}{1.54} 4\right)$ for $(1-\rho) \phi$, which vanishes identically near 0 . Then $\zeta=|x|^{-2}(1-\rho) \phi \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ and so

$$
\begin{equation*}
(1-\rho) \phi=|x|^{2} \zeta=\sum_{j=1}^{n} x_{j}\left(x_{j} \zeta\right) \tag{1.57}
\end{equation*}
$$

finally gives $\left(\frac{2.6}{1.5} 4\right)$ with $\psi_{j}(x)=\rho(x) \zeta_{j}(x)+x_{j} \zeta(x)$. Having proved the existence of such a decomposition we see that if $x_{j} u=0$ for all $j$ then

$$
\begin{equation*}
u(\phi)=u(\phi(0) \rho(x))+\sum_{j=1}^{n} u\left(x_{j} \psi_{j}\right)=c \phi(0), \quad c=u(\rho(x)), \tag{1.58}
\end{equation*}
$$

i.e. $u=c \delta(x)$.

$$
\begin{equation*}
\mathcal{F} \phi(\xi)=\int e^{-i \xi \cdot x} \phi(x) d x \tag{1.59}
\end{equation*}
$$

As you all know the inverse Fourier transform is given by

$$
\begin{equation*}
\mathcal{G} \psi(x)=(2 \pi)^{-n} \int e^{i x \cdot \xi} \psi(\xi) d \xi \tag{1.60}
\end{equation*}
$$

Since it is so important here I will give a proof of this invertibility. First however, let us note some of the basic properties.

Both $\mathcal{F}$ and $\mathcal{G}$ give continuous linear maps

$$
\begin{equation*}
\mathcal{F}, \mathcal{G}: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{1.61}
\end{equation*}
$$

To see this observe first that the integrals in ( $\frac{2.7}{1.5} 9$ ) and (2.8 1.60 ) are absolutely convergent:

$$
\begin{equation*}
|\mathcal{F} \phi(\xi)| \leq \int|\phi(x)| d x \leq \int\left(1+|x|^{2}\right)^{-n} d x \times \sup _{x \in \mathbb{R}^{n}}\left(1+|x|^{2}\right)^{n}|\phi(x)| \tag{1.62}
\end{equation*}
$$

where we use the definition of $\mathcal{S}\left(\mathbb{R}^{n}\right)$. In fact this shows that $\sup |\mathcal{F} \phi|<\infty$ if $\phi \in$ $\mathcal{S}\left(\mathbb{R}^{n}\right)$. Formal differentiation under the integral sign gives an absolutely convergent integral:

$$
D_{j} \mathcal{F} \phi(\xi)=\int D_{\xi_{j}} e^{-i x \xi} \phi(x) d x=\int e^{-i x \cdot \xi}\left(-x_{j} \phi\right) d x
$$

since $\sup \left(1+|x|^{2}\right)^{n}\left|x_{j} \phi\right|<\infty$. Then it follows that $D_{j} \mathcal{F} \phi$ is also bounded, i.e. $\mathcal{F} \phi$ is differentiable, and $\left(\frac{2.10}{1.7}\right.$ holds. This argument can be extended to show that $\mathcal{F} \phi$ is $\mathcal{C}^{\infty}$,

$$
\begin{equation*}
D^{\alpha} \mathcal{F} \phi(\xi)=\mathcal{F}\left((-x)^{\alpha} \phi\right) \tag{1.63}
\end{equation*}
$$

$$
\frac{2.7}{1.59}
$$

Similarly, starting from (1.7.59), we can use integration by parts to show that

$$
\xi_{j} \mathcal{F} \phi(\xi)=\int e^{-i x \xi} \xi_{j} \phi(x) d x=\int e^{-i x \cdot \xi}\left(D_{j} \phi\right)(x) d x
$$

i.e. $\xi_{j} \mathcal{F} \phi=\mathcal{F}\left(D_{j} \phi\right)$. Combining this with (2.11 1.63$)$ gives

$$
\begin{equation*}
\xi^{\alpha} D_{\xi}^{\beta} \mathcal{F} \phi=\mathcal{F}\left(D^{\alpha} \cdot\left[(-x)^{\beta} \phi\right]\right) \tag{1.64}
\end{equation*}
$$

Since $D_{x}^{\alpha}\left((-x)^{\beta} \phi\right) \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ we conclude

$$
\begin{equation*}
\sup \left|\xi^{\alpha} D_{\zeta}^{\beta} \mathcal{F} \phi\right|<\infty \Longrightarrow \mathcal{F} \phi \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{1.65}
\end{equation*}
$$

This map is continuous since

$$
\begin{gathered}
\sup \left|\xi^{\alpha} D_{\xi}^{\beta} \mathcal{F} \phi\right| \leq C \cdot \sup _{x} \mid\left(1+|x|^{2}\right)^{n} D_{x}^{\alpha}\left[(-x)^{\beta} \phi\right] \\
\Longrightarrow\|\mathcal{F} \phi\|_{k} \leq C_{k}\|\phi\|_{k+2 n}, \quad \forall k
\end{gathered}
$$

The identity ( ${ }^{2.12} 1.64$ ), written in the form

$$
\begin{gather*}
\mathcal{F}\left(D_{j} \phi\right)=\xi_{j} \mathcal{F} \phi \\
\mathcal{F}\left(x_{j} \phi\right)=-D_{\xi_{j}} \mathcal{F} \phi \tag{1.66}
\end{gather*}
$$

is already the key to the proof of invertibility:

### 2.14

Theorem 1.3. The Fourier transform gives an isomorphism $\mathcal{F}: \mathcal{S}\left(\mathbb{R}^{n}\right) \longleftrightarrow$ $\mathcal{S}\left(\mathbb{R}^{n}\right)$ with inverse $\mathcal{G}$.

Proof. We shall use the idea of the Schwartz kernel theorem. It is important not to use this theorem itself, since the Fourier transform is a key tool in the (simplest) proof of the kernel theorem. Thus we consider the composite map

$$
\begin{equation*}
\mathcal{G} \circ \mathcal{F}: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{1.67}
\end{equation*}
$$

and write down its kernel. Namely

$$
\begin{gather*}
K(\phi)=(2 \pi)^{-n} \iiint e^{i y \cdot \xi-i x \cdot \xi} \phi(y, x) d x d \xi d y  \tag{1.68}\\
\forall \phi \in \mathcal{S}\left(\mathbb{R}_{y}^{n} \times \mathbb{R}_{x}^{n}\right) \Longrightarrow K \in \mathcal{S}^{\prime}\left(\mathbb{R}^{2 n}\right)
\end{gather*}
$$

The integrals in $\left(\frac{2.15}{1.68}\right)$ are iterated, i.e. should be performed in the order indicated. Notice that if $\psi, \zeta \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ then indeed

$$
\begin{align*}
(\mathcal{G} \cdot \mathcal{F}(\psi))(\zeta)=\int \zeta(y)(2 \pi)^{-n}\left(\int e^{i y \cdot \xi} \int e^{-i x \cdot \xi} \psi(x) d x d \xi\right) & d y d \xi d y  \tag{1.69}\\
& =K(\zeta \boxtimes \psi)
\end{align*}
$$

so $K$ is the Schwartz kernel ${ }_{2}$ of $\mathcal{G} \cdot \mathcal{F}$.
The two identities ( $\left(\frac{2.13}{1.66}\right)$ translate (with essentially the same proofs) to the conditions on $K$ :
2.17

$$
\left\{\begin{array}{l}
\left(D_{x_{j}}+D_{y_{j}}\right) K(x, y)=0  \tag{1.70}\\
\left(x_{j}-y_{j}\right) K(x, y)=0
\end{array} \quad j=1, \ldots, n\right.
$$

Next we use the freedom to make linear changes of variables, setting

$$
\begin{align*}
& K_{L}(x, z)=K(x, x-z), K_{L} \in \mathcal{S}^{\prime}\left(\mathbb{R}^{2 n}\right)  \tag{1.71}\\
& \text { i.e. } K_{L}(\phi)=K(\psi), \psi(x, y)=\phi(x, x-y)
\end{align*}
$$

where the notation will be explained later. Then ( $\left.\frac{2.17}{1.70}\right)$ becomes
2.19

$$
\begin{equation*}
D_{x_{j}} K_{L}(x, z)=0 \text { and } z_{j} K_{L}(x, z)=0 \text { for } j=1, \ldots n \tag{1.72}
\end{equation*}
$$

This puts us in a position to apply the two little lemmas. The first says $K_{L}(x, z)=$ $f(z)$ for some $f \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ and then the second says $f(z)=c \delta(z)$. Thus

$$
\begin{equation*}
K(x, y)=c \delta(x-y) \Longrightarrow \mathcal{G} \cdot \mathcal{F}=c \mathrm{Id} \tag{1.73}
\end{equation*}
$$

It remains only to show that $c=1$. That $c \neq 0$ is obvious (since $\mathcal{F}(\delta)=1$ ). The easiest way to compute the constant is to use the integral identity

$$
\begin{equation*}
\int_{-\infty}^{\infty} e^{-x^{2}} d x=\pi^{\frac{1}{2}} \tag{1.74}
\end{equation*}
$$

to show that ${ }^{4}$

[^2]\[

$$
\begin{gather*}
\mathcal{F}\left(e^{-|x|^{2}}\right)=\pi^{\frac{n}{2}} e^{-|\xi|^{2} / 4} \\
\Longrightarrow \mathcal{G}\left(e^{-|\xi|^{2} / 4}\right)=\pi^{-\frac{n}{2}} e^{-|x|^{2}}  \tag{1.75}\\
\Longrightarrow \mathcal{G} \cdot \mathcal{F}=\operatorname{Id} .
\end{gather*}
$$
\]

Now $(2 \pi)^{n} \mathcal{G}$ is actually the adjoint of $\mathcal{F}$ :

$$
\begin{equation*}
\int \phi(\zeta) \overline{\mathcal{F} \psi}(\zeta) d \zeta=(2 \pi)^{n} \int(\mathcal{G} \phi) \cdot \bar{\psi} d x \forall \phi, \psi \in \mathcal{S}\left(\mathbb{R}^{n}\right) . \tag{1.76}
\end{equation*}
$$

It follows that we can extend $\mathcal{F}$ to a map on tempered distributions

$$
\begin{gather*}
\mathcal{F}: \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \\
\mathcal{F} u(\bar{\phi})=u\left((2 \pi)^{n} \overline{\mathcal{G} \phi}\right) \forall \phi \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{1.77}
\end{gather*}
$$

Then we conclude
2.20 Corollary 1.1. The Fourier transform extends by continuity to an isomorphism

$$
\begin{equation*}
\mathcal{F}: \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \longleftrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \tag{1.78}
\end{equation*}
$$

with inverse $\mathcal{G}$, satisfying the identities ( $\frac{2.13}{1.66)}$.
Although I have not discussed Lebesgue integrability I assume familiarity with the basic Hilbert space

$$
\begin{aligned}
& L^{2}\left(\mathbb{R}^{n}\right)= \\
& \qquad\left\{u: \mathbb{R}^{n} \longrightarrow \mathbb{C} ; f \text { is measurable and } \int_{\mathbb{R}^{n}}|f(x)|^{2} d x<\infty\right\} / \sim \\
& \\
& \quad f \sim g \Longleftrightarrow f=g \text { almost everywhere. }
\end{aligned}
$$

This also injects by the same integration map $\left(\frac{1.2 .2000 .274}{1.102) \text { with }} \mathcal{S}\left(\mathbb{R}^{n}\right)\right.$ as a dense subset

$$
\mathcal{S}\left(\mathbb{R}^{n}\right) \hookrightarrow L^{2}\left(\mathbb{R}^{p}\right) \hookrightarrow \mathcal{S}\left(\mathbb{R}^{n}\right)
$$

Proposition 1.3. The Fourier transform extends by continuity from the dense subspace $\mathcal{S}\left(\mathbb{R}^{n}\right) \subset L^{2}\left(\mathbb{R}^{n}\right)$, to an isomorphism

$$
\mathcal{F}: L^{2}\left(\mathbb{R}^{n}\right) \longleftrightarrow L^{2}\left(\mathbb{R}^{n}\right)
$$

satisfying $\|\mathcal{F} u\|_{L^{2}}=(2 \pi)^{\frac{1}{2} n}\|u\|_{L^{2}}$.
${ }_{103}$ Roof. Given the density of $\mathcal{S}\left(\mathbb{R}^{n}\right)$ in $L^{2}\left(\mathbb{R}^{n}\right)$, this is also a consequence of $\left(\frac{1}{1.103}\right)$, since setting $\phi=\mathcal{F} u$, for $u \in \mathcal{S}\left(\mathbb{R}^{n}\right)$, gives Parseval's formula

$$
\int \mathcal{F} u(\zeta) \overline{\mathcal{F} v(\zeta)}=(2 \pi)^{n} \int u(x) \overline{v(x)} d x
$$

Setting $v=u$ gives norm equality (which is Plancherel's formula).
An outline of the proof of the density statement is given in the problems below.

### 1.8. Differential operators

The simplest examples of the Fourier transform of distributions are immediate consequences of the definition and (1.66). Thus

$$
\begin{equation*}
\mathcal{F}(\delta)=1 \tag{1.79}
\end{equation*}
$$

as already noted and hence, from (1.136),

$$
\begin{equation*}
\mathcal{F}\left(D^{\alpha} \delta(x)\right)=\xi^{\alpha} \quad \forall \alpha \in \mathbb{N}_{0}^{n} \tag{1.80}
\end{equation*}
$$

Now, recall that the space of distributions with support the point 0 is just:

$$
\begin{equation*}
\left\{u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) ; \sup (u) \subset\{0\}\right\}=\left\{u=\sum_{\text {finite }} c_{\alpha} D^{\alpha} \delta\right\} \tag{1.81}
\end{equation*}
$$

Thus we conclude that the Fourier transform gives an isomorphism

$$
\begin{equation*}
\mathcal{F}:\left\{u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) ; \operatorname{supp}(u) \subset\{0\}\right\} \longleftrightarrow \mathbb{C}[\xi]=\{\text { polynomials in } \xi\} . \tag{1.82}
\end{equation*}
$$

Another way of looking at this same isomorphism is to consider partial differential operators with constant coefficients:

$$
\begin{gather*}
P(D): \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}\left(\mathbb{R}^{n}\right) \\
P(D)=\sum c_{\alpha} D^{\alpha} . \tag{1.83}
\end{gather*}
$$

The identity becomes

$$
\begin{equation*}
\mathcal{F}(P(D) \phi)(\xi)=P(\xi) \mathcal{F}(\phi)(\xi) \forall \phi \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{1.84}
\end{equation*}
$$

and indeed the same formula holds for all $\phi \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$. Using the simpler notation $\hat{u}(\xi)=\mathcal{F} u(\xi)$ this can be written

$$
\begin{equation*}
P \widehat{(D) u}(\xi)=P(\xi) \hat{u}(\xi), P(\xi)=\sum c_{\alpha} \xi^{\alpha} \tag{1.85}
\end{equation*}
$$

The polynomial $P$ is called the (full) characteristic polynomial of $P(D)$; of course it determines $P(D)$ uniquely.

It is important for us to extend this formula to differential operators with variable coefficients. Using (1.59) and the inverse Fourier transform we get

$$
\begin{equation*}
P(D) u(x)=(2 \pi)^{-n} \iint e^{i(x-y) \cdot \xi} P(\xi) u(y) d y d \xi \tag{1.86}
\end{equation*}
$$

where again this is an iterated integral. In particular the inversion formula is just the case $P(\xi)=1$. Consider the space

$$
\begin{equation*}
\mathcal{C}_{\infty}^{\infty}\left(\mathbb{R}^{n}\right)=\left\{u: \mathbb{R}^{n} \longrightarrow \mathbb{C} ; \sup _{x}\left|D^{\alpha} u(x)\right|<\infty \forall \alpha\right\} \tag{1.87}
\end{equation*}
$$

the space of $\mathcal{C}^{\infty}$ function with all derivatives bounded on $\mathbb{R}^{n}$. Of course

$$
\begin{equation*}
\mathcal{S}\left(\mathbb{R}^{n}\right) \subset \mathcal{C}_{\infty}^{\infty}\left(\mathbb{R}^{n}\right) \tag{1.88}
\end{equation*}
$$

but $\mathcal{C}_{\infty}^{\infty}\left(\mathbb{R}^{n}\right)$ is much bigger, in particular $1 \in \mathcal{C}_{\infty}^{\infty}\left(\mathbb{R}^{n}\right)$. Now by Leibniz' formula

$$
\begin{equation*}
D^{\alpha}(u v)=\sum_{\beta \leq \alpha}\binom{\alpha}{\beta} D^{\beta} u \cdot D^{\alpha-\beta} v \tag{1.89}
\end{equation*}
$$

it follows that $\mathcal{S}\left(\mathbb{R}^{n}\right)$ is a module over $\mathcal{C}_{\infty}^{\infty}\left(\mathbb{R}^{n}\right)$. That is,

$$
\begin{equation*}
u \in \mathcal{C}_{\infty}^{\infty}\left(\mathbb{R}^{n}\right), \phi \in \mathcal{S}\left(\mathbb{R}^{n}\right) \Longrightarrow u \phi \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{1.90}
\end{equation*}
$$

From this it follows that if

$$
\begin{equation*}
P(x, D)=\sum_{|\alpha| \leq m} p_{\alpha}(x) D^{\alpha}, p_{\alpha} \in \mathcal{C}_{\infty}^{\infty}\left(\mathbb{R}^{n}\right) \tag{1.91}
\end{equation*}
$$

then $P(x, D): \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}\left(\mathbb{R}^{n}\right)$. The formula $\left(\frac{2.26}{1.86}\right)$ extends to

$$
\begin{equation*}
P(x, D) \phi=(2 \pi)^{-n} \int e^{i(x-y) \cdot \xi} P(x, \xi) \phi(y) d y d \xi \tag{1.92}
\end{equation*}
$$

where again this is an iterated integral. Here

$$
\begin{equation*}
P(x, \xi)=\sum_{|\alpha| \leq m} p_{\alpha}(x) \xi^{\alpha} \tag{1.93}
\end{equation*}
$$

is the (full) characteristic polynomial of $P$.

### 1.9. Radial compactification

For later purposes, and general propaganda, consider the quadratic radial compactification of $\mathbb{R}^{n}$. The smooth map

$$
\begin{equation*}
\mathrm{QRC}: \mathbb{R}^{n} \ni x \longmapsto \frac{x}{\left(1+|x|^{2}\right)^{\frac{1}{2}}} \in \mathbb{R}^{n} \tag{1.94}
\end{equation*}
$$

is 1-1 and maps onto the interior of the unit ball, $\mathbb{B}^{n}=\{|x| \leq 1\}$. Consider the subspace

$$
\begin{equation*}
\dot{\mathcal{C}}^{\infty}\left(\mathbb{B}^{n}\right)=\left\{u \in \mathcal{S}\left(\mathbb{R}^{n}\right) ; \operatorname{supp}(u) \subset \mathbb{B}^{n}\right\} \tag{1.95}
\end{equation*}
$$

This is just the set of smooth functions on $\mathbb{R}^{n}$ which vanish outside the unit ball. Then the composite ('pull-back') map

$$
\begin{equation*}
\mathrm{QRC}^{*}: \dot{\mathcal{C}}^{\infty}\left(\mathbb{B}^{n}\right) \ni u \longmapsto u \circ \mathrm{QRC} \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{1.96}
\end{equation*}
$$

is a topological isomorphism. A proof is indicated in the problems below.
The dual space of $\dot{\mathcal{C}}^{\infty}\left(\mathbb{B}^{n}\right)$ is generally called the space of 'extendible distributions' on $\mathbb{B}^{n}$ - because they are all given by restricting elements of $\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ to $\dot{\mathcal{C}}^{\infty}\left(\mathbb{B}^{n}\right)$. Thus QRC also identifies the tempered distributions on $\mathbb{R}^{n}$ with the extendible distributions on $\mathbb{B}^{n}$. We shall see below that various spaces of functions on $\mathbb{R}^{n}$ take interesting forms when pulled back to $\mathbb{B}^{n}$. I often find it useful to 'bring infinity in' in this way.

Why is this the 'quadratic' radial compactification, and not just the radial compactification? There is a good reason which is discussed in the problems below. The actual radial compactification is a closely related map which identifies Euclidean space, $\mathbb{R}^{n}$, with the interior of the upper half of the $n$-sphere in $\mathbb{R}^{n+1}$ :
1.2 .2000 .275
(1.97) $\quad \mathrm{RC}: \mathbb{R}^{n} \ni x \longmapsto\left(\frac{1}{\left(1+|x|^{2}\right)^{\frac{1}{2}}}, \frac{x}{\left(1+|x|^{2}\right)^{\frac{1}{2}}}\right)$

$$
\in \mathbb{S}^{n, 1}=\left\{X=\left(X_{0}, X^{\prime}\right) \in \mathbb{R}^{n+1} ; X_{0} \geq 0, X_{0}^{2}+\left|X^{\prime}\right|^{2}=1\right\}
$$

Since the half-sphere is diffeomorphic to the ball (as compact manifolds with boundary) these two maps can be compared - they are not the same. However it is true that RC also identifies $\mathcal{S}\left(\mathbb{R}^{n}\right)$ with $\dot{\mathcal{C}}^{\infty}\left(\mathbb{S}^{n, 1}\right)$.

### 1.10. Problems

Problem 1.1. Suppose $\phi: \mathbb{R}^{n} \longrightarrow \mathbb{C}$ is a function such that for each point $\bar{x} \in \mathbb{R}^{n}$ and each $k \in \mathbb{N}_{0}$ there exists a constant $\epsilon_{k}>0$ and a polynomial $p_{k}(x ; \bar{x})$ (in $x$ ) for which

$$
\begin{equation*}
\left|\phi(x)-p_{k}(x ; \bar{x})\right| \leq \frac{1}{\epsilon_{k}}|x-\bar{x}|^{k+1} \quad \forall|x-\bar{x}| \leq \epsilon_{k} \tag{1.98}
\end{equation*}
$$

Does it follow that $\phi$ is infinitely differentiable - either prove this or give a counterexample.

Problem 1.2. Show that the function $u(x)=\exp (x) \cos \left[e^{x}\right]$ 'is' a tempered distribution. Part of the question is making a precise statement as to what this means!

Problem 1.3. Write out a careful (but not necessarily long) proof of the 'easy' direction of the Schwartz kernel theorem, that any $K \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n+m}\right)$ defines a continuous linear operator

$$
\begin{equation*}
O_{K}: \mathcal{S}\left(\mathbb{R}^{m}\right) \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \tag{1.99}
\end{equation*}
$$

[with respect to the weak topology on $\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ and the metric topology on $\mathcal{S}\left(\mathbb{R}^{m}\right)$ ] by

$$
\begin{equation*}
O_{K} \phi(\psi)=K(\psi \boxtimes \phi) \tag{1.100}
\end{equation*}
$$

[Hint: Work out what the continuity estimate on the kernel, $K$, means when it is paired with an exterior product $\psi \boxtimes \phi$.]

Problem 1.4. Show that $d$ in $\left(\frac{1.7}{1.6}\right)$ is a metric on $\mathcal{S}\left(\mathbb{R}^{n}\right)$. [Hint: If $\|\cdot\|$ is a norm on a vector space show that

$$
\left.\frac{\|u+v\|}{1+\|u+v\|} \leq \frac{\|u\|}{1+\|u\|}+\frac{\|v\|}{1+\|v\|} .\right]
$$

Problem 1.5. Show that a sequence $\phi_{n}$ in $\mathcal{S 1} \mathcal{S}\left(\mathbb{R}_{2000}^{n}\right.$. is ${ }_{260}$ Cauchy, resp. converges to $\phi$, with respect to the metric $d$ in Problem 1.4 if and only if $\phi_{n}$ is Cauchy, resp. converges to $\phi$, with respect to each of the norms $\|\cdot\|_{k}$.
31.1 .2000 .262

Problem 1.6. Show that a linear map $F: \mathcal{S}\left(\mathbb{R}^{n}{ }_{31.2}{ }_{1.2000}\left(\mathbb{R}_{260}^{p}\right)\right.$ is continuous with respect to the metric topology given in Problem $\sqrt{1.4}$ if and only if for each $k$ there exists $N(k) \in \mathbb{N}$ a constant $C_{k}$ such that

$$
\|F \phi\|_{k} \leq C_{k}\|\phi\|_{N(k)} \forall \phi \in \mathcal{S}\left(\mathbb{R}^{n}\right)
$$

Give similar equivalent conditions for continuity of a linear map $f: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathbb{C}$ and for a bilinear map $\mathcal{S}\left(\mathbb{R}^{n}\right) \times \mathcal{S}\left(\mathbb{R}^{p}\right) \longrightarrow \mathbb{C}$.
31.1 .2000 .265

Problem 1.7. Check the continuity of ( $\frac{11.121,2000.264}{1}$
Problem 1.8. Prove Proposition $\frac{1.17}{1.2}$. [Hint: It is only necessary to show that if $u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ is fixed then for any of the open sets in (1.1), $B$, (with all the $\epsilon_{j}>0$ ) there is an element $\phi \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ such that $u-T_{\phi} \in B$. First show that if $\phi_{1}^{\prime}, \ldots \phi_{p}^{\prime}$ is a basis for $\Phi$ then the set

$$
\begin{equation*}
B^{\prime}=\left\{v \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) ;\left|\left\langle v, \phi_{j}^{\prime}\right\rangle\right|<\delta_{j}\right. \tag{1.101}
\end{equation*}
$$

is contained in $B$ if the $\delta_{j}>0$ are chosen small enough. Taking the basis to be orthonormal, show that $u-\psi \in B^{\prime}$ can be arranged for some $\psi \in \Phi$.]

Problem 1.9. Compute the Fourier transform of $\exp \left(-|x|^{2}\right) \in \mathcal{S}\left(\mathbb{R}^{n}\right)$. [Hint: The Fourier integral is a product of 1-dimensional integrals so it suffices to assume $x \in \mathbb{R}$. Then

$$
\int e^{-i \xi x} e^{-x^{2}} d x=e^{-\xi^{2} / 4} \int e^{-\left(x+\frac{i}{2} \xi\right)^{2}} d x .
$$

Interpret the integral as a contour integral and shift to the new contour where $x+\frac{i}{2} \xi$ is real.]

Problem 1.10. Show that $\left(\frac{1.13}{1.23}\right)$ makes sense for $\phi \in L^{2}\left(\mathbb{R}^{n}\right)$ (the space of (equivalence classes of) Lebesgue square-integrable functions and that the resulting $\operatorname{map} L^{2}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ is an injection.

Problem 1.11. Suppose $u \in L^{2}\left(\mathbb{R}^{n}\right)$ and that

$$
D_{1} D_{2} \cdots D_{n} u \in(1+|x|)^{-n-1} L^{2}\left(\mathbb{R}^{n}\right)
$$

where the derivatives are defined using Problem $\frac{\mathbb{P 1} 1.5}{1.10}$. Using repeated integration, show that $u$ is necessarily a bounded continuous function. Conclude further that for $u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$

$$
\begin{align*}
& D^{\alpha} u \in(1+|x|)^{-n-1} L^{2}\left(\mathbb{R}^{n}\right) \forall|\alpha| \leq k+n \\
\Longrightarrow & D^{\alpha} u \text { is bounded and continuous for }|\alpha| \leq k . \tag{1.102}
\end{align*}
$$

[This is a weak form of the Sobolev embedding theorem.]
P2.1
Problem 1.12. The support of a (tempered) distribution can be defined in terms of the support of a test function. For $\phi \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ the $\operatorname{support}, \operatorname{supp}(\phi)$, is the closure of the set of points at which it takes a non-zero value. For $u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ we define

$$
\begin{equation*}
\operatorname{supp}(u)=O^{\complement}, O=\bigcup\left\{O^{\prime} \subset \mathbb{R}^{n} \text { open; } \operatorname{supp}(\phi) \subset O^{\prime} \Longrightarrow u(\phi)=0\right\} \tag{1.103}
\end{equation*}
$$

Show that the definitions for $\mathcal{S}\left(\mathbb{R}^{n}\right)$ and $\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ are consistent with the inclusion $\mathcal{S}\left(\mathbb{R}^{n}\right) \subset \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$. Prove that $\operatorname{supp}(\delta)=\{0\}$.

P2.2 Problem 1.13. For simplicity in $\mathbb{R}$, i.e. with $n=1$, prove Schwartz theorem concerning distributions with support the origin. Show that with respect to the norm $\|\cdot\|_{k}$ the space

$$
\begin{equation*}
\{\phi \in \mathcal{S}(\mathbb{R}) ; \phi(x)=0 \text { in }|x|<\epsilon, \epsilon=\epsilon(\phi)>0\} \tag{1.104}
\end{equation*}
$$

is dense in

$$
\begin{equation*}
\left\{\phi \in \mathcal{S}(\mathbb{R}) ; \phi(x)=x^{k+1} \psi(x), \psi \in \mathcal{S}(\mathbb{R})\right\} \tag{1.105}
\end{equation*}
$$

Use this to show that

$$
\begin{equation*}
u \in \mathcal{S}^{\prime}(\mathbb{R}), \operatorname{supp}(u) \subset\{0\} \Longrightarrow u=\sum_{\ell, \text { finite }} c_{\ell} D_{x}^{\ell} \delta(x) \tag{1.106}
\end{equation*}
$$

P2.3 Problem 1.14. Show that if $P$ is a differential operator with coefficients in $\mathcal{C}_{\infty}^{\infty}\left(\mathbb{R}^{n}\right)$ then $P$ is local in the sense that

$$
\begin{equation*}
\operatorname{supp}(P u) \subset \operatorname{supp}(u) \quad \forall u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \tag{1.107}
\end{equation*}
$$

The converse of this, for an operator $P: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}\left(\mathbb{R}^{n}\right)$ where (for simplicity) we assume

$$
\begin{equation*}
\operatorname{supp}(P u) \subset K \subset \mathbb{R}^{n} \tag{1.108}
\end{equation*}
$$

for a fixed compact set $K$, is Peetre's theorem. How would you try to prove this? (No full proof required.)
P2.4 Problem 1.15. (Schwartz representation theorem) Show that, for any $p \in \mathbb{R}$ the map
2.28
(1.109) $\quad R_{p}: \mathcal{S}\left(\mathbb{R}^{n}\right) \ni \phi \longmapsto\left(1+|x|^{2}\right)^{-p / 2} \mathcal{F}^{-1}\left[\left(1+|\xi|^{2}\right)^{-p / 2} \mathcal{F} \phi\right] \in \mathcal{S}\left(\mathbb{R}^{n}\right)$
is an isomorphism and, using Problem $\frac{\mathbb{P 1} 1.6}{1.1}$ or otherwise,

$$
\begin{equation*}
p \geq n+1+k \Longrightarrow\left\|R_{p} \phi\right\|_{k} \leq C_{k}\|\phi\|_{L^{2}}, \forall \phi \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{1.110}
\end{equation*}
$$

Let $R_{p}^{t}: \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ be the dual map (defined by $\left.T_{p}^{t} u(\phi)=u\left(R_{p} \phi\right)\right)$. Show that $R_{p}^{t}$ is an isomorphism and that if $u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ satisfies

$$
\begin{equation*}
|u(\phi)| \leq C\|\phi\|_{k}, \forall \phi \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{1.111}
\end{equation*}
$$

then $R_{p}^{t} u \in L_{\mathbb{P} 1.5}^{2}\left(\mathbb{R}^{n}\right)$, if $p \geq n+1+k$, in ${ }^{2}$, in Problem $\frac{10}{1.10}$. Using Problem 1.11 show that $R_{n+1}\left(R_{n+1+k}^{t} u\right)$ is bounded and continuous and hence that

$$
\begin{equation*}
u=\sum_{|\alpha|+|\beta| \leq 2 n+2+k} x^{\beta} D^{\alpha} u_{\alpha, \beta} \tag{1.112}
\end{equation*}
$$

for some bounded continuous functions $u_{\alpha, \beta}$.
Problem 1.16. (Schwartz kernel theorem.) Show that any continuous linear operator

$$
T: \mathcal{S}\left(\mathbb{R}_{y}^{m}\right) \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}_{x}^{n}\right)
$$

extends to a continuous linear operator

$$
T:\left(1+|y|^{2}\right)^{-k / 2} H^{k}\left(\mathbb{R}_{y}^{m}\right) \longrightarrow\left(1+|x|^{2}\right)^{-q / 2} H^{q}\left(\mathbb{R}_{x}^{n}\right)
$$

for some $k$ and $q$. Deduce that the operator

$$
\begin{aligned}
\tilde{T}=\left(1+\left|D_{x}\right|^{2}\right)^{(-n-1-q) / 2}\left(1+|x|^{2}\right)^{q / 2} \circ T \circ\left(1+|y|^{2}\right)^{k / 2}\left(1+|D|^{2}\right)^{-k / 2}: \\
L^{2}\left(\mathbb{R}^{m}\right) \longrightarrow \mathcal{C}_{\infty}\left(\mathbb{R}^{n}\right)
\end{aligned}
$$

is continuous with values in the bounded continuous functions on $\mathbb{R}^{n}$. Deduce that $\tilde{T}$ has Schwartz kernel in $\mathcal{C}_{\infty}\left(\mathbb{R}^{n} ; L^{2}\left(\mathbb{R}^{m}\right)\right) \subset \mathcal{S}^{\prime}\left(\mathbb{R}^{n+m}\right)$ and hence that $T$ itself has a tempered Schwartz kernel.

Problem 1.17. Radial compactification and symbols.
Problem 1.18. Series of problems discussing double polyhomogeneous symbols.

## CHAPTER 2

## Pseudodifferential operators on Euclidean space

Formula $\left(\frac{2.27}{1.92}\right)$ for the action of a differential operator (with coefficients in $\mathcal{C}_{\infty}^{\infty}\left(\mathbb{R}^{n}\right)$ ) on $\mathcal{S}\left(\mathbb{R}^{n}\right)$ can be written

$$
\begin{align*}
P(x, D) u & =(2 \pi)^{-n} \int e^{i(x-y) \cdot \xi} P(x, \xi) u(y) d y d \xi  \tag{2.1}\\
& =(2 \pi)^{-n} \int e^{i x \cdot \xi} P(x, \xi) \hat{u}(\xi) d \xi
\end{align*}
$$

where $\hat{u}(\xi)=\mathcal{F} u(\xi)$ is the Fourier transform of $u$. We shall generalize this formula by generalizing $P(x, \xi)$ from a polynomial in $\xi$ to a symbol, which is to say a smooth function satisfying certain uniformity conditions at infinity. In fact we shall also allow the symbol, or rather the amplitude, in the integral ( $\left(\frac{3.1}{2.1}\right)$ to depend in addition on the 'incoming' variables, $y$ :

$$
\begin{equation*}
A(x, D) u=(2 \pi)^{-n} \int e^{i(x-y) \cdot \xi} a(x, y, \xi) u(y) d y d \xi, u \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{2.2}
\end{equation*}
$$

Of course it is not clear that this integral is well-defined.
To interpret $\left(\frac{3.2}{2.2}\right)$ we shall first look into the definition and properties of symbols. Then we show how this integral can be interpreted as an oscillatory integral and that it thereby defines an operator on $\mathcal{S}\left(\mathbb{R}^{n}\right)$. We then investigate the properties of these pseudodifferential operators at some length.

$$
\begin{equation*}
|p(\xi)| \leq C(1+|\xi|)^{m} \quad \forall \xi \in \mathbb{R}^{n} \tag{2.3}
\end{equation*}
$$

Since successive derivatives, $D_{\xi}^{\alpha} p(\xi)$, are polynomials of degree $m-|\alpha|$, for any multiindex $\alpha$, we get the family of estimates

$$
\begin{equation*}
\left|D_{\xi}^{\alpha} p(\xi)\right| \leq C_{\alpha}(1+|\xi|)^{m-|\alpha|} \quad \forall \xi \in \mathbb{R}^{n}, \alpha \in \mathbb{N}_{0}^{n} \tag{2.4}
\end{equation*}
$$

Of course if $|\alpha|>m$ then $D_{\xi}^{\alpha} p \equiv 0$, so we can even take the constant $C_{\alpha}$ to be independent of $\alpha$. If we consider the characteristic polynomial $P(x, \xi)$ of a differential operator of order $m_{13}$ with coefficients in $\mathcal{C}_{\infty}^{\infty}\left(\mathbb{R}^{n}\right)$ (i.e. all derivatives of the coefficients are bounded) ( $\overline{2.4}$ ) is replaced by

$$
\begin{equation*}
\left|D_{x}^{\alpha} D_{\xi}^{\beta} P(x, \xi)\right| \leq C_{\alpha, \beta}(1+|\xi|)^{m-|\beta|} \quad \forall(x, \xi) \in \mathbb{R}^{n} \times \mathbb{R}^{n}, \alpha, \beta \in \mathbb{N}_{0}^{n} \tag{2.5}
\end{equation*}
$$

There is no particular reason to have the same number of $x$ variables as of $\xi$ variables, so in general we define:

Definition 2.1. The space $S_{\infty}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$ of symbols of order $m$ consists of those functions $a \in \mathcal{C}^{\infty}\left(\mathbb{R}^{p} \times \mathbb{R}^{n}\right)$ satisfying all the estimates

$$
\begin{equation*}
\left|D_{z}^{\alpha} D_{\xi}^{\beta} a(z, \xi)\right| \leq C_{\alpha, \beta}(1+|\xi|)^{m-|\beta|} \text { on } \mathbb{R}^{p} \times \mathbb{R}^{n} \quad \forall \alpha \in \mathbb{N}_{0}^{p}, \beta \in \mathbb{N}_{0}^{n} \tag{2.6}
\end{equation*}
$$

For later reference we even define $S_{\infty}^{m}\left(\Omega ; \mathbb{R}^{n}\right)$ when $\Omega_{3} \subset \mathbb{R}^{p}$ and $\Omega \subset \operatorname{clos}(\operatorname{int}(\Omega))$ as consisting of those $a \in \mathcal{C}^{\infty}\left(\operatorname{int}(\Omega) \times \mathbb{R}^{n}\right)$ satisfying (LV.6) for $(z, \xi) \in \operatorname{int}(\Omega) \times \mathbb{R}^{n}$.

The estimates $\left(\frac{3.7}{2.6}\right)$ can be rewritten

$$
\begin{equation*}
\|a\|_{N, m}=\sup _{\substack{z \in \operatorname{int}(\Omega) \\ \xi \in \mathbb{R}^{n}}} \max _{|\alpha|+|\beta| \leq N}(1+|\xi|)^{-m+|\beta|}\left|D_{z}^{\alpha} D_{\xi}^{\beta} a(z, \xi)\right|<\infty \tag{2.7}
\end{equation*}
$$

With these norms $S_{\infty}^{m}\left(\Omega ; \mathbb{R}^{n}\right)$ is a Fréchet space, rather similar in structure to $\mathcal{C}_{\infty}^{\infty}\left(\mathbb{R}^{n}\right)$. Thus the topology is given by the metric

$$
\begin{equation*}
d(a, b)=\sum_{N \geq 0} 2^{-N} \frac{\|a-b\|_{N, m}}{1+\|a-b\|_{N, m}}, a, b \in S_{\infty}^{m}\left(\Omega ; \mathbb{R}^{n}\right) \tag{2.8}
\end{equation*}
$$

The subscript ' $\infty$ ' here is not standard notation. It refers to the assumption of uniform boundedness of the derivatives of the 'coefficients'. More standard notation would be just $S^{m}\left(\Omega \times \mathbb{R}^{n}\right)$, especially for $\Omega=\mathbb{R}^{p}$, but I think this is too confusing.

A more significant issue is: Why this class precisely? As we shall see below, there are other choices which are not only possible but even profitable to make. However, the present one has several virtues. It is large enough to cover most of the straightforward things we want to do (at least initially) and small enough to 'work' easily. It leads to what I shall refer to as the 'traditional' algebra of pseudodifferential operators.

Now to some basic properties. First notice that

$$
\begin{equation*}
(1+|\xi|)^{m} \leq C(1+|\xi|)^{m^{\prime}} \forall \xi \in \mathbb{R}^{n} \Longleftrightarrow m \leq m^{\prime} \tag{2.9}
\end{equation*}
$$

Thus we have an inclusion

$$
\begin{equation*}
S_{\infty}^{m}\left(\Omega ; \mathbb{R}^{n}\right) \hookrightarrow S_{\infty}^{m^{\prime}}\left(\Omega ; \mathbb{R}^{n}\right) \forall m^{\prime} \geq m \tag{2.10}
\end{equation*}
$$

Moreover this inclusion is continuous, since from ( $\left(\frac{3.8}{2.7}\right),\|a\|_{N, m^{\prime}} \leq\|a\|_{N, m}$ if $a \in$ $S^{m}\left(\Omega ; \mathbb{R}^{n}\right)$ and $m^{\prime} \geq m$. Since these spaces increase with $m$ we think of them as a filtration of the big space

$$
\begin{equation*}
S_{\infty}^{\infty}\left(\Omega ; \mathbb{R}^{n}\right)=\bigcup_{m} S_{\infty}^{m}\left(\Omega ; \mathbb{R}^{n}\right) \tag{2.11}
\end{equation*}
$$

Notice that the two ' $\infty$ s' here are quite different. The subscript refers to the fact that the 'coefficients' are bounded and stands for $L^{\infty}$ whereas the superscript $\infty$ stands really for $\mathbb{R}$. The residual space of this filtration is

$$
\begin{equation*}
S_{\infty}^{-\infty}\left(\Omega ; \mathbb{R}^{n}\right)=\bigcap_{m} S_{\infty}^{m}\left(\Omega ; \mathbb{R}^{n}\right) \tag{2.12}
\end{equation*}
$$

In fact the inclusion $\left(\frac{3.10}{2.10}\right)$ is never dense if $m^{\prime}>m$. Instead we have the following rather technical, but nevertheless very useful, result.

Lemma 2.1. For any $m \in \mathbb{R}$ and any $a \in S_{\infty}^{m}\left(\Omega ; \mathbb{R}^{n}\right)$ there is a sequence in $S_{\infty}^{-\infty}\left(\Omega ; \mathbb{R}^{n}\right)$ which is bounded in $S_{\infty}^{m}\left(\Omega ; \mathbb{R}^{n}\right)$ and converges to $a$ in the topology of $S_{\infty}^{m^{\prime}}\left(\Omega ; \mathbb{R}^{n}\right)$ for any $m^{\prime}>m$; in particular $S_{\infty}^{-\infty}\left(\Omega ; \mathbb{R}^{n}\right)$ is dense in the space $S_{\infty}^{m}\left(\Omega ; \mathbb{R}^{n}\right)$ in the topology of $S_{\infty}^{m^{\prime}}\left(\Omega ; \mathbb{R}^{n}\right)$ for $m^{\prime}>m$.

The reason one cannot take $m^{\prime}=m$ here is essentially the same reason that underlies the fact that $\mathcal{S}\left(\mathbb{R}^{n}\right)$ is not dense in $\mathcal{C}_{\infty}^{\infty}\left(\mathbb{R}^{n}\right)$. Namely any uniform limit obtained from a converging Schwartz sequence must vanish at infinity. In particular the constant function $1 \in S_{\infty}^{0}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$ cannot be in the closure in this space of $S_{\infty}^{-\infty}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$ if $n>0$.

Proof. Choose $\phi \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)$ with $0 \leq \phi(\xi) \leq 1, \phi(\xi)=1$ if $|\xi|<1, \phi(\xi)=0$ if $|\xi|>2$ and consider the sequence

$$
\begin{equation*}
a_{k}(z, \xi)=\phi(\xi / k) a(z, \xi), \quad a \in S_{\infty}^{m}\left(\Omega ; \mathbb{R}^{n}\right) \tag{2.13}
\end{equation*}
$$

We shall show that $a_{k} \in S_{\infty}^{-\infty}\left(\Omega, \mathbb{R}^{n}\right)$ is a bounded sequence in $S_{\infty}^{m}\left(\Omega ; \mathbb{R}^{n}\right)$ and that $a_{k} \longrightarrow a$ in $S_{\infty}^{m^{\prime}}\left(\Omega ; \mathbb{R}^{n}\right)$ for any $m^{\prime}>m$. Certainly for each $N$

$$
\begin{equation*}
\left|a_{k}(z, \xi)\right| \leq C_{N, k}(1+|\xi|)^{-N} \tag{2.14}
\end{equation*}
$$

since $\phi$ has compact support. Leibniz' formula gives

$$
\begin{equation*}
D_{z}^{\alpha} D_{\xi}^{\beta} a_{k}(z, \xi)=\sum_{\beta^{\prime} \leq \beta}\binom{\beta^{\prime}}{\beta} k^{-\left|\beta^{\prime}\right|}\left(D^{\beta^{\prime}} \phi\right)(\xi / k) D_{z}^{\alpha} D_{\xi}^{\beta-\beta^{\prime}} a(z, \xi) \tag{2.15}
\end{equation*}
$$

On the support of $\phi(\xi / k),|\xi| \leq k$ so, using the symbol estimates on $a$, it follows that $a_{k}$ is bounded in $S_{\infty}^{m}\left(\Omega ; \mathbb{R}^{n}\right)$. We easily conclude that

$$
\begin{equation*}
\left|D_{z}^{\alpha} D_{\xi}^{\beta} a_{k}(z, \xi)\right| \leq C_{N, \alpha, \beta, k}(1+|\xi|)^{-N} \quad \forall \alpha, \beta, N, k \tag{2.16}
\end{equation*}
$$

Thus $a_{k} \in S_{\infty}^{-\infty}\left(\Omega ; \mathbb{R}^{n}\right)$.
So consider the difference

$$
\begin{equation*}
\left(a-a_{k}\right)(z, \xi)=(1-\phi)(\xi / k) a(z, \xi) \tag{2.17}
\end{equation*}
$$

Now, $|(1-\phi)(\xi / k)|=0$ in $|\xi| \leq k$ so we only need estimate the difference in $|\xi| \geq k$ where this factor is bounded by 1 . In this region $1+|\xi| \geq 1+k$ so, since $-m^{\prime}+m<0$,

$$
\begin{align*}
& (1+|\xi|)^{-m^{\prime}}\left|\left(a-a_{k}\right)(z, \xi)\right| \leq  \tag{2.18}\\
& \quad(1+k)^{-m^{\prime}+m} \sup _{z, \xi}\left|(1+|\xi|)^{-m}\right| a(z, \xi) \mid \leq(1+k)^{-m^{\prime}+m}\|a\|_{0, m} \longrightarrow 0
\end{align*}
$$

This is convergence with respect to the first symbol norm.
Next consider the $\xi$ derivatives of ( $\overline{2} .147)$. Using Leibniz' formula

$$
\begin{gathered}
D_{\xi}^{\beta}\left(a-a_{k}\right)=\sum_{\gamma \leq \beta}\binom{\beta}{\gamma} D_{\xi}^{\beta-\gamma}(1-\phi)\left(\frac{\xi}{k}\right) \cdot D_{\xi}^{\gamma} a(z, \xi) \\
=(1-\phi)\left(\frac{\xi}{k}\right) \cdot D_{\xi}^{\beta} a(z, \xi)-\sum_{\gamma<\beta}\binom{\beta}{\gamma}\left(D^{\beta-\gamma} \phi\right)\left(\frac{\xi}{k}\right) \cdot k^{-|\beta-\gamma|} D_{\xi}^{\gamma} a(z, \xi) .
\end{gathered}
$$

In the first term, $D_{\xi}^{\beta} a(z, \xi)$ is a symbol of order $m-|\beta|$, so by the same argument as above

$$
\begin{equation*}
\sup _{\xi}(1+|\xi|)^{-m^{\prime}+|\beta|}(1-\phi)\left(\frac{\xi}{k}\right) D_{\xi}^{\beta} a(x, \xi) \longrightarrow 0 \tag{2.19}
\end{equation*}
$$

as $k \longrightarrow \infty$ if $m^{\prime}>m$. In all the other terms, $\left(D^{\beta-\gamma} \phi\right)(\zeta)$ has compact support, in fact $1 \leq|\zeta| \leq 2$ on the support. Thus for each term we get a bound

$$
\begin{equation*}
\sup _{k \leq|\xi| \leq 2 k}(1+|\xi|)^{-m^{\prime}+|\beta|} \cdot k^{-|\beta-\gamma|} C \cdot(1+|\xi|)^{m-|\gamma|} \leq C k^{-m^{\prime}+m} . \tag{2.20}
\end{equation*}
$$

The variables $z$ play the rôle of parameters so we have in fact shown that

$$
\begin{equation*}
\sup _{\substack{z \in \Omega \\ \xi \in \mathbb{R}^{n}}}(1+|\xi|)^{-m^{\prime}+|\beta|}\left|D_{z}^{\alpha} D_{\xi}^{\beta}\left(a-a_{k}\right)\right| \longrightarrow 0 \text { as } k \longrightarrow \infty \tag{2.21}
\end{equation*}
$$

This means $a_{k} \longrightarrow a$ in each of the symbol norms, and hence in the topology of $S_{\infty}^{m^{\prime}}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$ as desired.

In fact this proof suggests a couple of other 'obvious' results. Namely

$$
\begin{equation*}
S_{\infty}^{m}\left(\Omega ; \mathbb{R}^{n}\right) \cdot S_{\infty}^{m^{\prime}}\left(\Omega ; \mathbb{R}^{n}\right) \subset S_{\infty}^{m+m^{\prime}}\left(\Omega ; \mathbb{R}^{n}\right) \tag{2.22}
\end{equation*}
$$

This can be proved directly using Leibniz' formula:

$$
\begin{aligned}
& \sup _{\xi}(1+|\xi|)^{-m-m^{\prime}+|\beta|}\left|D_{z}^{\alpha} D_{\xi}^{\beta}(a(z, \xi) \cdot b(z, \xi))\right| \\
& \leq \sum_{\substack{\mu \leq \alpha \\
\gamma \leq \beta}}\binom{\alpha}{\mu}\binom{\beta}{\gamma} \sup _{\xi}(1+|\xi|)^{-m+|\gamma|}\left|D_{z}^{\mu} D_{\xi}^{\gamma} a(z, \xi)\right| \\
& \quad \times \sup _{\xi}(1+|\xi|)^{-m^{\prime}+|\beta-\gamma|}\left|D_{z}^{\alpha-\mu} D_{\xi}^{\beta-\gamma} b(z, \xi)\right|<\infty
\end{aligned}
$$

We also note the action of differentiation:

$$
\begin{align*}
D_{z}^{\alpha}: S_{\infty}^{m}\left(\Omega ; \mathbb{R}^{n}\right) \longrightarrow S_{\infty}^{m}\left(\Omega ; \mathbb{R}^{n}\right) \text { and } \\
D_{\xi}^{\beta}: S_{\infty}^{m}\left(\Omega ; \mathbb{R}^{n}\right) \longrightarrow S_{\infty}^{m-|\beta|}\left(\Omega ; \mathbb{R}^{n}\right) \tag{2.23}
\end{align*}
$$

In fact, while we are thinking about these things we might as well show the important consequence of ellipticity. A symbol $a \in S_{\infty}^{m}\left(\Omega ; \mathbb{R}^{n}\right)$ is said to be (globally) elliptic if

$$
\begin{equation*}
|a(z, \xi)| \geq \epsilon(1+|\xi|)^{m}-C(1+|\xi|)^{m-1}, \epsilon>0 \tag{2.24}
\end{equation*}
$$

or equivalently ${ }^{1}$

$$
\begin{equation*}
|a(z, \xi)| \geq \epsilon(1+|\xi|)^{m} \text { in }|\xi| \geq C_{\epsilon}, \epsilon>0 \tag{2.25}
\end{equation*}
$$

Lemma 2.2. If $a \in S_{\infty}^{m}\left(\Omega ; \mathbb{R}^{n}\right)$ is elliptic there exists $b \in S_{\infty}^{-m}\left(\Omega ; \mathbb{R}^{n}\right)$ such that

$$
\begin{equation*}
a \cdot b-1 \in S_{\infty}^{-\infty}\left(\Omega ; \mathbb{R}^{n}\right) \tag{2.26}
\end{equation*}
$$

Proof. Using $\left(\frac{3.20}{2.25}\right)$ choose $\phi$ as in the proof of Lemma $\frac{\frac{3}{2} .13}{2.1}$ and set

$$
b(z, \xi)= \begin{cases}\frac{1-\phi(\xi / 2 C)}{a(z, \xi)} & |\xi| \geq C  \tag{2.27}\\ 0 & |\xi| \leq C\end{cases}
$$

Then $b$ is $\mathcal{C}^{\infty}$ since $b=0$ in $C \leq|\xi| \leq C+\delta$ for some $\delta>0$. The symbol estimates follow by noting that, in $|\xi| \geq C$,

$$
\begin{equation*}
D_{z}^{\alpha} D_{\xi}^{\beta} b=a^{-1-|\alpha|-|\beta|} \cdot G_{\alpha \beta} \tag{2.28}
\end{equation*}
$$

where $G_{\alpha \beta}$ is a symbol of order $(|\alpha|+|\beta|) m-|\beta|$. This may be proved by induction. Indeed, it is trune when $\alpha=\beta=0$. Assuming $\left(\frac{1}{2.28}\right)$ for some $\alpha$ and $\beta$, differentiation of ( $\overline{2} .28$ ) gives

$$
\begin{aligned}
D_{z_{j}} D_{z}^{\alpha} D_{\xi}^{\beta} b=D_{z_{j}} a^{-1-|\alpha|-|\beta|} \cdot G_{\alpha \beta} & =a^{-2-|\alpha|-|\beta|} G^{\prime} \\
G^{\prime} & =(-1-|\alpha|-|\beta|)\left(D_{z_{j}} a\right) G_{\alpha \beta}+a D_{z_{j}} G_{\alpha \beta}
\end{aligned}
$$

[^3]By the inductive hypothesis, $G^{\prime}$ is a symbol of order $(|\alpha|+1+|\beta|) m-|\beta|$. A similar argument applies to derivatives with respect to the $\xi$ variables.

### 2.2. Pseudodifferential operators

Now we proceed to discuss the formula $\left(\frac{3.2}{2.2}\right)$ where we shall assume that, for some $w, m \in \mathbb{R}$,

$$
\begin{gather*}
a(x, y, \xi)=\left(1+|x-y|^{2}\right)^{w / 2} \tilde{a}(x, y, \xi)  \tag{2.29}\\
\tilde{a} \in S_{\infty}^{m}\left(\mathbb{R}_{(x, y)}^{2 n} ; \mathbb{R}_{\xi}^{n}\right)
\end{gather*}
$$

The extra 'weight' factor (which allows polynomial growth in the direction of $x-y$ ) turns out, somewhat enigmatically, to both make no difference and be very useful! Notice ${ }^{2}$ that if $a \in \mathcal{C}^{\infty}\left(\mathbb{R}^{2 n} \times \mathbb{R}^{n}\right)$ then $a \in\left(1+|x-y|^{2}\right)^{w / 2} S^{m}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right)$ if and only if

$$
\begin{equation*}
\left|D_{x}^{\alpha} D_{y}^{\beta} D_{\xi}^{\gamma} a(x, y, \xi)\right| \leq C_{\alpha, \beta, \gamma}(1+|x-y|)^{w}(1+|\xi|)^{m-|\gamma|} \forall \alpha, \beta, \gamma \in \mathbb{N}_{0}^{n} \tag{2.30}
\end{equation*}
$$

If $m<-n$ then, for each $u \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ the integral in $\left(\frac{\sqrt[3]{2.2}}{2.2}\right)$ is absolutely convergent, locally uniformly in $x$, since

$$
\begin{align*}
& |a(x, y, \xi) u(y)| \leq C(1+|x-y|)^{w}(1+|\xi|)^{m}(1+|y|)^{-N} \\
& \quad \leq C(1+|x|)^{w}(1+|\xi|)^{m}(1+|y|)^{m}, m<-n . \tag{2.31}
\end{align*}
$$

Here we have used the following simple consequence of the triangle inequality

$$
(1+|x-y|) \leq(1+|x|)(1+|y|)
$$

from which it follows that

$$
(1+|x-y|)^{w} \leq \begin{cases}(1+|x|)^{w}(1+|y|)^{w} & \text { if } w>0  \tag{2.32}\\ (1+|x|)^{w}(1+|y|)^{-w} & \text { if } w \leq 0\end{cases}
$$

Thus we conclude that, provided $m<-n$,

$$
\begin{equation*}
A: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow\left(1+|x|^{2}\right)^{w / 2} \mathcal{C}_{\infty}^{0}\left(\mathbb{R}^{n}\right) \tag{2.33}
\end{equation*}
$$

To show that, for general $m, A$ exists as an operator, we prove that its Schwartz kernel exists.

Proposition 2.1. The map, defined for $m<-n$ as a convergent integral, (2.34)

$$
\begin{aligned}
& \left(1+|x-y|^{2}\right)^{w / 2} S_{\infty}^{m}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right) \ni a \longmapsto I(a)= \\
& \quad(2 \pi)^{-n} \int e^{i(x-y) \cdot \xi} a(x, y, \xi) d \xi \in\left(1+|x|^{2}+|y|^{2}\right)^{w / 2} \mathcal{C}_{\infty}^{0}\left(\mathbb{R}^{2 n}\right)
\end{aligned}
$$

extends by continuity to
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$$
\begin{equation*}
I:\left(1+|x-y|^{2}\right)^{w / 2} S_{\infty}^{m}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right) \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{2 n}\right) \tag{2.35}
\end{equation*}
$$

for each $w, m \in \mathbb{R}$ in the topology of $S_{\infty}^{m^{\prime}}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right)$ for any $m^{\prime}>m$.

[^4]Proof. Since we already have the density of $S_{\infty}^{-\infty}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right)$ in $S_{\infty}^{m}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right)$ in the toplogy of $S_{3 \times 2}^{m^{\prime}}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right)$ for any $m^{\prime}>m$, we only need to show the continuity of the map ( $\overline{2.34}$ ) on this residual subspace with respect to the topology of $S_{\infty}^{m^{\prime}}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right)$ for any $m^{\prime}$, which we may as well write as $m$. What we shall show is that, for each $w, m_{1} \in \mathbb{R}$, there are integers $N, k \in \mathbb{N}$ such that, in terms of the norms in ( $\left(\frac{3.7}{2.7}\right)$ and (1.7. $)$

$$
\begin{align*}
|I(a)(\phi)| \leq C\|\tilde{a}\|_{N, m}\|\phi\|_{k} & \forall \phi \in \mathcal{S}\left(\mathbb{R}^{2 n}\right),  \tag{2.36}\\
& a=\left(1+|x-y|^{2}\right)^{w / 2} \tilde{a}, \tilde{a} \in S_{\infty}^{-\infty}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right)
\end{align*}
$$

To see this we just use integration by parts.
Set $\tilde{\phi}(x, y)=\left(1+|x-y|^{2}\right)^{w / 2} \phi(x, y)$. Observe that

$$
\begin{aligned}
& \left(1+\xi \cdot D_{x}\right) e^{i(x-y) \cdot \xi}=\left(1+|\xi|^{2}\right) e^{i(x-y) \cdot \xi} \\
& \left(1-\xi \cdot D_{y}\right) e^{i(x-y) \cdot \xi}=\left(1+|\xi|^{2}\right) e^{i(x-y) \cdot \xi}
\end{aligned}
$$

Thus we can write, for $\tilde{a} \in S_{\infty}^{-\infty}$, with $a=\left(1+|x-y|^{2}\right)^{w / 2} \tilde{a}$ and for any $q \in \mathbb{N}$

$$
\begin{align*}
I(a)(\phi)= & \iint(2 \pi)^{-n} \int e^{i(x-y) \cdot \xi}\left(1+|\xi|^{2}\right)^{-2 q} \\
& \left(1-\xi \cdot D_{x}\right)^{q}\left(1+\xi \cdot D_{y}\right)^{q}[\tilde{a}(x, y, \xi) \tilde{\phi}(x, y)] d \xi d x d y  \tag{2.37}\\
= & \sum_{|\gamma| \leq 2 q} \iint\left(\int e^{i(x-y) \cdot \xi} a_{\gamma}^{(q)}(x, y, \xi) d \xi\right) D_{(x, y)}^{\gamma} \tilde{\phi}(x, y) d x d y
\end{align*}
$$

Here the $a_{\gamma}^{(q)}$ arise by expanding the powers of the operator

$$
\left(1-\xi \cdot D_{x}\right)^{q}\left(1+\xi \cdot D_{x}\right)^{q}=\sum_{|\mu|,|\nu| \leq q} C_{\mu, \nu} \xi^{\mu+\nu} D_{x}^{\mu} D_{y}^{\nu}
$$

and applying Leibniz' formula. Thus $a_{\gamma}^{(q)}$ arises from terms in which $2 q-|\gamma|$ derivatives act on $\tilde{a}$ so it is of the form

$$
\begin{gathered}
a_{\gamma}=\left(1+|\xi|^{2}\right)^{-2 q} \sum_{|\mu| \leq|\gamma|,|\gamma| \leq 2 q} C_{\mu, \gamma} \xi^{\gamma} D_{(x, y)}^{\mu} \tilde{a} \\
\Longrightarrow\left\|a_{\gamma}\right\|_{N, m} \leq C_{m, q, N}\|\tilde{a}\|_{N+2 q, m-2 q} \forall m, N, q
\end{gathered}
$$

So (for given $m$ ) if we take $-2 q+m<-n$, e.g. $q>\max \left(\frac{n+m}{2}, 0\right)$ and use the integrability of $(1+|x|+|y|)^{-2 n-1}$ on $\mathbb{R}^{2 n}$, then

$$
\begin{equation*}
|I(a)(\phi)| \leq C\|\tilde{a}\|_{2 q, m}\|\tilde{\phi}\|_{2 q+2 n+1} \leq C\|\tilde{a}\|_{2 q, m}\|\phi\|_{2 q+w+2 n+1} \tag{2.38}
\end{equation*}
$$

This is the estimate $\left(\frac{3.27}{2.36}\right)$, which proves the desired continuity.
In showing the existence of the Schwartz' kernel in this proof we do not really need to integrate by parts in both $x$ and $y$; either separately will do the trick. We can use this observation to show that these pseudodifferential operator act on $\mathcal{S}\left(\mathbb{R}^{n}\right)$.

Lemma 2.3. If $a \in\left(1+|x-y|^{2}\right)^{w / 2} S_{\infty}^{m}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right)$ then the operator $A$, with Schwartz kernel $I(a)$, is a continuous linear map

$$
\begin{equation*}
A: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{2.39}
\end{equation*}
$$

We shall denote by $\Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ the linear space of operators $\left(\frac{3.30}{2.39}\right)$, corresponding to $\left(1+|x-y|^{2}\right)^{-w / 2} a \in S_{\infty}^{m}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right)$ for some $w$. I call them pseudodifferential operators 'af traditional type' - or type ' 1,0 ', the meaning of which is explained in Problem $\overline{2} .16$ below.

Proof. Proceeding as in $\left(\frac{\sqrt[3]{2} .28}{2.37}\right)$ but only integrating by parts in $y$ we deduce that, for $q$ large depending on $m$,

$$
\begin{aligned}
A u(\psi) & =\sum_{\gamma \leq 2 q}(2 \pi)^{-n} \iiint e^{i(x-y) \cdot \xi} a_{\gamma}(x, y, \xi) D_{y}^{\gamma} u(y) d \xi \psi(x) d y d x \\
a_{\gamma} & \in\left(1+|x-y|^{2}\right)^{w / 2} S^{m-q}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right) \text { if } a \in\left(1+|x-y|^{2}\right)^{w / 2} S^{m}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right) .
\end{aligned}
$$

The integration by parts is justified by continuity from $S_{\infty}^{-\infty}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right)$. Taking $-q+$ $m<-n-|w|$, this shows that $A u$ is given by the convergent integral

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| :--- |

$$
\begin{align*}
& A u(x)=\sum_{\gamma \leq 2 q}(2 \pi)^{-n} \iint e^{i(x-y) \cdot \xi} a_{\gamma}(x, y, \xi) D_{y}^{\gamma} u(y) d \xi d y  \tag{2.40}\\
& A: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow\left(1+|x|^{2}\right)^{\frac{|w|}{2}} \mathcal{C}_{\infty}^{0}\left(\mathbb{R}^{n}\right)
\end{align*}
$$

which is really just $\left(\frac{3.33}{2.33}\right)$ again. Here $\mathcal{C}_{\infty}^{0}\left(\mathbb{R}^{n}\right)$ is the Banach space of bounded continuous functions on $\mathbb{R}^{n}$, with the supremum norm. The important point is that the weight depends on $w$ but not on $m$. Notice that

$$
D_{x_{j}} A u(x)=(2 \pi)^{-n} \sum_{|\gamma| \leq 2 q} \iint e^{i(x-y) \cdot \xi}\left(\xi_{j}+D_{x_{j}}\right) a_{\gamma} \cdot D_{y}^{\gamma} u(y) d y d \xi
$$

and

$$
x_{j} A u(x)=(2 \pi)^{n} \sum_{\substack{|\gamma| \leq 2 q}} \iint e^{i(x-y) \cdot \xi}\left(-D_{\xi_{j}}+y_{j}\right) a_{\gamma} \cdot D_{y}^{\gamma} u(y) d y d \xi
$$

Proceeding inductively $\left(\frac{3.30}{2.39}\right)$ follows from $\left(\frac{3.33}{2.33}\right)$ or $\left(\frac{1.2 .2000 .282}{2.40}\right)$ since we conclude that

$$
x^{\alpha} D_{x}^{\beta} A u \in\left(1+|x|^{2}\right)^{\frac{|w|}{2}} \mathcal{C}_{\infty}^{0}\left(\mathbb{R}^{n}\right), \forall \alpha, \beta \in \mathbb{N}_{0}^{n}
$$

and this implies that $A u \in \mathcal{S}\left(\mathbb{R}^{n}\right)$.

### 2.3. Composition

There are two extreme cases of $I(a)$, namely where $a$ is independent of either $x$ or of $y$. Below we shall prove:

ThEOREM 2.1 (Reduction). Each $A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ can be written uniquely as $I\left(a^{\prime}\right)$ where $a^{\prime} \in S_{\infty}^{m}\left(\mathbb{R}_{x}^{n} ; \mathbb{R}_{\xi}^{n}\right)$.

This is the main step in proving the fundamental result of this Chapter, which is that two pseudodifferential operators can be composed to give a pseudodifferential operator and that the orders are additive. Thus our aim is to demonstrate the fundamental
4.1 Theorem 2.2. [Composition] The space $\Psi_{\infty}^{\infty}\left(\mathbb{R}^{n}\right)$ is an order-filtered $*$-algebra on $\mathcal{S}\left(\mathbb{R}^{n}\right)$.
${ }_{30}$ We have already shown that each $A \in \Psi_{\infty}^{\infty}\left(\mathbb{R}^{n}\right)$ defines a continuous linear map (2.39). We now want to show that
4.3

$$
\begin{gather*}
A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right) \Longrightarrow A^{*} \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)  \tag{2.41}\\
A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right), B \in \Psi_{\infty}^{m^{\prime}}\left(\mathbb{R}^{n}\right) \Longrightarrow A \circ B \in \Psi_{\infty}^{m+m^{\prime}}\left(\mathbb{R}^{n}\right), \tag{2.42}
\end{gather*}
$$

since this is what is menant by an order-filtered (the orders add on composition) *-algebra (meaning (ㄹ.4I) holds). In fact we will pick up some more information along the way.

### 2.4. Reduction

We proceed to prove Theorem $\frac{3.31}{2.1,}$ which we can restate as:
Proposition 2.2. The range of $\left(\frac{3.26}{2.34}\right)$ (for any $w$ ) is the same as the range of I restricted to the image of the inclusion map

$$
S_{\infty}^{m}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right) \ni a \longmapsto a(x, \xi) \in S_{\infty}^{m}\left(\mathbb{R}_{(x, y)}^{2 n} ; \mathbb{R}^{n}\right)
$$

Proof. Suppose $a \in\left(1+|x-y|^{2}\right)^{w / 2} S_{\infty}^{-\infty}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right)$ for some $w$, then

$$
\begin{equation*}
I\left(\left(x_{j}-y_{j}\right) a\right)=I\left(-D_{\xi_{j}} a\right) j=1, \ldots, n . \tag{2.43}
\end{equation*}
$$

Indeed this is just the result of inserting the identity

$$
D_{\xi_{j}} e^{i(x-y) \cdot \xi}=\left(x_{j}-y_{j}\right) e^{i(x-y) \cdot \xi}
$$

into $\left(\frac{3.26}{2.34}\right)$ and integrating by parts. Since both sides of $\left(\frac{4.5}{2.43)}\right.$ are continuous on $\left(1+|x-y|^{2}\right)^{w / 2} S_{\infty}^{\infty}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right)$ the identity $y_{5}$ holds in general. Notice that if $a$ is of order $m$ then $D_{\xi_{j}} a$ is of order $m-1$, so ( $\overline{L .43}$ ) shows that even though the operator with amplitude $\left(x_{j}-y_{j}\right) a(x, y, \xi)$ appears to have order $m$, it actually has order $m-1$.

To exploit ( ${ }^{4.4} \mathbf{2} .43$ ) consider the Taylor series (with Legendre's remainder) for $a(x, y, \xi)$ around $x=y$ :
(2.44) $a(x, y, \xi)=\sum_{|\alpha| \leq N-1} \frac{(-i)^{|\alpha|}}{\alpha!}(x-y)^{\alpha}\left(D_{y}^{\alpha} a\right)(x, x, \xi)$

$$
+\sum_{|\alpha|=N} \frac{(-i)^{|\alpha|}}{\alpha!}(x-y)^{\alpha} \cdot R_{N, \alpha}(x, y, \xi) .
$$

Here,

$$
\begin{equation*}
R_{N, \alpha}(x, y, \xi)=\int_{0}^{1}(1-t)^{N-1}\left(D_{y}^{\alpha} a\right)(x,(1-t) x+t y, \xi) d t \tag{2.45}
\end{equation*}
$$

Now,

$$
\begin{equation*}
(x-y)^{\alpha}\left(D_{y}^{\alpha} a\right)(x, y, \xi) \in\left(1+|x-y|^{2}\right)^{\frac{(w+|\alpha|)}{2}} S_{\infty}^{m}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right) \tag{2.46}
\end{equation*}
$$

Applying ( ${ }^{(2.5} .43$ ) repeatedly we see that if $A$ is the operator with kernel $I(a)$ then

$$
\begin{equation*}
A=\sum_{j=0}^{N-1} A_{j}+R_{N}, \quad A_{j} \in \Psi_{\infty}^{m-j}\left(\mathbb{R}^{n}\right), R_{N} \in \Psi_{\infty}^{m-N}\left(\mathbb{R}^{n}\right) \tag{2.47}
\end{equation*}
$$

where the $A_{j}$ have kernels

$$
\begin{equation*}
I\left(\sum_{|\alpha|=j} \frac{i^{|\alpha|}}{\alpha!}\left(D_{y}^{\alpha} D_{\xi}^{\alpha} a\right)(x, x, \xi)\right) \tag{2.48}
\end{equation*}
$$

To proceed further we need somehow to sum this series. Of course we cannot really do this, but we can come close!

### 2.5. Asymptotic summation

Suppose $a_{j} \in S_{\infty}^{m-j}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$. The fact that the orders are decreasing means that those symbols are getting very small, for $|\xi|$ large. The infinite series

$$
\begin{equation*}
\sum_{j} a_{j}(z, \xi) \tag{2.49}
\end{equation*}
$$

need not converge. However we shall say that it converges asymptotically, or since it is a series we say it is 'asymptotically summable,' if there exists $a \in S_{\infty}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$ such that,

$$
\begin{equation*}
\text { for every } N, a-\sum_{j=0}^{N-1} a_{j} \in S_{\infty}^{m-N}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right) \tag{2.50}
\end{equation*}
$$

We write this relation as
4.11

$$
\begin{equation*}
a \sim \sum_{j=0}^{\infty} a_{j} . \tag{2.51}
\end{equation*}
$$

Proposition 2.3. Any series $a_{j} \in S_{\infty}^{m-j}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$ is asymptotically summable, in the sense of $\left(\frac{4.50}{Z}\right)$, and the asymptotic sum is well defined up to an additive term in $S_{\infty}^{-\infty}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$.

Proof. The uniqueness part is easy. Suppose $a$ and $a^{\prime}$ both satisfy $\left(\frac{4.10}{2.50}\right)$. Taking the difference

$$
\begin{equation*}
a-a^{\prime}=\left(a-\sum_{j=0}^{N-1} a_{j}\right)-\left(a^{\prime}-\sum_{j=0}^{N-1} a_{j}\right) \in S_{\infty}^{m-N}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right) \tag{2.52}
\end{equation*}
$$

Since $S_{\infty}^{-\infty}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$ is just the intersection of the $S_{\infty}^{-N}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$ over $N$ it follows that $a-a^{\prime} \in S_{\infty}^{-\infty}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$, proving the uniqueness.

So to the existence of an asymptotic sum. To construct this (by Borel's $\operatorname{method}^{3}$ ) we cut off each term 'near infinity in $\xi$ '. Thus fix $\phi \in \mathcal{C} \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right)$ with $\phi(\xi)=0$ in $|\xi| \leq 1, \phi(\xi)=1$ in $|\xi| \geq 2,0 \leq \phi(\xi) \leq 1$. Consider a decreasing sequence

$$
\begin{equation*}
\epsilon_{0}>\epsilon_{1}>\cdots>\epsilon_{j} \downarrow 0 \tag{2.53}
\end{equation*}
$$

We shall set

$$
\begin{equation*}
a(z, \xi)=\sum_{j=0}^{\infty} \phi\left(\epsilon_{j} \xi\right) a_{j}(z, \xi) \tag{2.54}
\end{equation*}
$$

Since $\phi\left(\epsilon_{j} \xi\right)=0$ in $|\xi|<1 / \epsilon_{j} \rightarrow \infty$ as $j \rightarrow \infty$, only finitely many of these terms are non-zero in any ball $|\xi| \leq R$. Thus $a(z, \xi)$ is a well-defined $\mathcal{C}^{\infty}$ function. Of course we need to consider the seminorms, in $S_{\infty}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$, of each term.

[^5]The first of these is

$$
\begin{equation*}
\sup _{z} \sup _{\xi}(1+|\xi|)^{-m}\left|\phi\left(\epsilon_{j} \xi\right)\right|\left|a_{j}(z, \xi)\right| . \tag{2.55}
\end{equation*}
$$

Now $|\xi| \leq \frac{1}{\epsilon_{j}}$ on the support of $\phi\left(\epsilon_{j} \xi\right) a_{j}(z, \xi)$ and since $a_{j}$ is a symbol of order $m-j$ this allows us to estimate ( $\frac{4.16}{2.55}$ ) by

$$
\begin{gathered}
\sup _{z} \sup _{|\xi| \leq \frac{1}{\epsilon_{j}}}(1+|\xi|)^{-j} \cdot\left[(1+|\xi|)^{-m+j}\left|a_{j}(z, \xi)\right|\right] \\
\leq\left(1+\frac{1}{\epsilon_{j}}\right)^{-j} \cdot C_{j} \leq \epsilon_{j}^{j} \cdot C_{j}
\end{gathered}
$$

where the $C_{j}$ 's are fixed constants, independent of $\epsilon_{j}$.
Let us look at the higher symbol estimates. As usual we can apply Leibniz' formula:

$$
\begin{gathered}
\sup _{z} \sup _{\xi}(1+|\xi|)^{-m+|\beta|}\left|D_{z}^{\alpha} D_{\xi}^{\beta} \phi\left(\epsilon_{j} \xi\right) a_{j}(z, \xi)\right| \\
\leq \sum_{\mu \leq \beta} \sup _{z} \sup _{\xi}(1+|\xi|)^{|\beta|-|\mu|-j} \epsilon_{j}^{|\beta|-|\mu|}\left|\left(D^{\beta-\mu} \phi\right)\left(\epsilon_{j} \xi\right)\right| \\
\times(1+|\xi|)^{-m+j+|\mu|}\left|D_{z}^{\alpha} D_{\xi}^{\mu} a_{j}(z, \xi)\right|
\end{gathered}
$$

The term with $\mu=\beta$ we estimate as before and the others, with $\mu \neq \beta$ are supported in $\frac{1}{\epsilon_{j}} \leq|\xi| \leq \frac{2}{\epsilon_{j}}$. Then we find that for all $j$

$$
\begin{equation*}
\left\|\phi\left(\epsilon_{j} \xi\right) a_{j}(z, \xi)\right\|_{N, m} \leq C_{N, j} \epsilon_{j}^{j} \tag{2.56}
\end{equation*}
$$

where $C_{N, j}$ is independent of $\epsilon_{j}$.
So we see that for each given $N$ we can arrange that, for instance,

$$
\left\|\phi\left(\epsilon_{j} \xi\right) a_{j}(z, \xi)\right\|_{N, m} \leq C_{N} \frac{1}{j^{2}}
$$

by choosing the $\epsilon_{j}$ to satify

$$
C_{N, j} \epsilon_{j}^{j} \leq \frac{1}{j^{2}} \forall j \geq j(N)
$$

Notice the crucial point here, we can arrange that for each $N$ the sequence of norms in ( $\frac{4.17}{2.56}$ ) is dominated by $C_{N} j^{-2}$ by fixing $\epsilon_{j}<\epsilon_{j, N}$ for large $j$. Thus we can arrange convergence of all the sums

$$
\sum_{j}\left\|\phi\left(\epsilon_{j} \xi\right) a_{j}(z, \xi)\right\|_{N, m}
$$

by diagonalization, for example setting $\epsilon_{j}=\frac{1}{2} \epsilon_{j, j}$. Thus by choosing $\epsilon_{j} \downarrow 0$ rapidly enough we ensure that the series ( $\overline{2.54}$ ) converges. In fact the same argument allows us to ensure that for every $N$

$$
\begin{equation*}
\sum_{j \geq N} \phi\left(\epsilon_{j} \xi\right) a_{j}(z, \xi) \text { converges in } S_{\infty}^{m-N}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right) \tag{2.57}
\end{equation*}
$$

This certainly gives $\left(\frac{4.10}{2.50}\right)$ with $a$ defined by $\left(\frac{4.15}{2.54}\right)$.

### 2.6. Residual terms

Now we can apply Proposition $\frac{4.12}{2.3}$ to the series in ( $\frac{4.9}{2.48)}$, that is we can find $b \in S_{\infty}^{m}\left(\mathbb{R}_{x}^{n} ; \mathbb{R}_{\xi}^{n}\right)$ satisfying

$$
\begin{equation*}
b(x, \xi) \sim \sum_{\alpha} \frac{i^{|\alpha|}}{\alpha!}\left(D_{y}^{\alpha} a\right)(x, x, \xi) \tag{2.58}
\end{equation*}
$$

Let $B=f_{4}\left(\frac{8}{8}\right)$ be the operator defined by this amplitude (which is independent of $y)$. Now (2.47) says that

$$
A-B=\sum_{j=0}^{N-1} A_{j}+R_{N}-B
$$

and from $\left(\frac{4.10}{2.50}\right)$ applied to $\left(\frac{4.19}{2.58}\right)$

$$
B=\sum_{j=0}^{N-1} A_{j}+R_{N}^{\prime}, R_{N}^{\prime} \in \Psi_{\infty}^{m-N}\left(\mathbb{R}^{n}\right)
$$

Thus

$$
\begin{equation*}
A-B \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)=\bigcap_{N} \Psi_{\infty}^{N}\left(\mathbb{R}^{n}\right) \tag{2.59}
\end{equation*}
$$

Notice that, at this stage, we do not know that $A-B$ has kernel $I(c)$ with $c \in S_{\infty}^{-\infty}\left(\mathbb{R}^{2 n}, \mathbb{R}^{n}\right)$, just that it has kernel $I\left(c_{N}\right)$ with $c_{N} \in S_{\infty}^{N}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right)$ for each $N$.

However:
Proposition 2.4. An operator $A: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ is an element of the space $\Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$ if and only if its Schwartz kernel is $\mathcal{C}^{\infty}$ and satisfies the estimates

$$
\begin{equation*}
\left|D_{x}^{\alpha} D_{y}^{\beta} K(x, y)\right| \leq C_{N, \alpha, \beta}(1+|x-y|)^{-N} \forall \alpha, \beta, N \tag{2.60}
\end{equation*}
$$

Proof. Suppose first that $A \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$, which means that ${ }_{26} A \in \Psi_{\infty}^{N}\left(\mathbb{R}^{n}\right)$ for every $N$. The Schwartz kernel, $K_{A}$, of $A$ is therefore given by ( $\overline{2.34}$ ) with the amplitude $a_{N} \in S_{\infty}^{N}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right)$. For $N \ll-n-1-p$ the integral converges absolutely and we can integrate by parts to show that

$$
\begin{gathered}
(x-y)^{\alpha} D_{x}^{\beta} D_{y}^{\gamma} K_{A}(x, y) \\
=(2 \pi)^{-N} \int e^{i(x-y) \cdot \xi}\left(-D_{\xi}\right)^{\alpha}\left(D_{x}+i \xi\right)^{\beta}\left(D_{y}-i \xi\right)^{\gamma} a_{N}(x, y, \xi) d \xi
\end{gathered}
$$

which converges absolutely, and uniformly in $x$, $y$, provided $|\beta|+|\gamma|+N-|\alpha|<-n$. Thus

$$
\sup \left|(x-y)_{\mid 4.22}^{\alpha} D_{x}^{\beta} D_{y}^{\gamma} K\right|<\infty \forall \alpha, \beta, \gamma
$$

which is another way of writing $\left(\frac{4.22}{2.60}\right)$ i.e.

$$
\sup \left(1+|x-y|^{2}\right)^{N}\left|D_{x}^{\beta} D_{y}^{\gamma} K\right|<\infty \forall \beta, \gamma, N
$$

Conversely suppose that $\left(\frac{4.22}{2.60}\right)$ holds. Define

$$
\begin{equation*}
g(x, z)=K(x, x-z) \tag{2.61}
\end{equation*}
$$

The estimates $\left(\frac{4.22}{2.60}\right)$ become

$$
\begin{equation*}
\sup \left|D_{x}^{\alpha} z^{\gamma} D_{z}^{\beta} g(x, z)\right|<\infty \forall \alpha, \beta, \gamma \tag{2.62}
\end{equation*}
$$

That is, $g$ is rapidly decreasing with all its derivatives in $z$. Taking the Fourier transform,

$$
\begin{equation*}
b(x, \xi)=\int e^{-i z \cdot \xi} g(x, z) d z \tag{2.63}
\end{equation*}
$$

the estimate $\left(\frac{4.24}{2.62}\right)$ translates to

$$
\begin{gather*}
\sup _{x, \xi}\left|D_{x}^{\alpha} \xi^{\beta} D_{\xi}^{\gamma} b(x, \xi)\right|<\infty \forall \alpha, \beta, \gamma  \tag{2.64}\\
\Longleftrightarrow b \in S_{\infty}^{-\infty}\left(\mathbb{R}_{x}^{n} ; \mathbb{R}_{\xi}^{n}\right)
\end{gather*}
$$

Now the inverse Fourier transform in $\left(\frac{4.25}{2.63}\right)$, combined with $\left(\frac{4.23}{2.61}\right)$ gives

$$
4.27
$$

$$
\begin{equation*}
K(x, y)=g(x, x-y)=(2 \pi)^{-n} \int e^{i(x-y) \cdot \xi} b(x, \xi) d \xi \tag{2.65}
\end{equation*}
$$

i.e. $K=I(b)$. This certainly proves the proposition and actually gives the stronger result.

$$
\begin{equation*}
A \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right) \Longleftrightarrow A=I(c), c \in S_{\infty}^{-\infty}\left(\mathbb{R}_{x}^{n} ; \mathbb{R}_{\xi}^{n}\right) \tag{2.66}
\end{equation*}
$$

This also finishes the proof of Proposition $\frac{4.4}{2.2}$ since in $\left(\frac{4.19}{2.58}\right),\left(\frac{4.20}{2.59}\right)$ we have shown that

$$
\begin{equation*}
A=B+R, B=I(b), \quad R \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right) \tag{2.67}
\end{equation*}
$$

so in fact

$$
4.29
$$

$$
\begin{equation*}
A=I(e), e \in S_{\infty}^{m}\left(\mathbb{R}_{x}^{n} ; \mathbb{R}_{\xi}^{n}\right), e \sim \sum_{\alpha} \frac{i^{|\alpha|}}{\alpha!}\left(D_{y}^{\alpha} D_{\xi}^{\alpha} a\right)(x, x, \xi) \tag{2.68}
\end{equation*}
$$

### 2.7. Proof of Composition Theorem

First consider the adjoint formula. If

$$
A: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}\left(\mathbb{R}^{n}\right)
$$

the adjoint is the operator

$$
A^{*}: \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)
$$

defined by duality:

$$
\begin{equation*}
A^{*} u(\bar{\phi})=u(\overline{A \phi}) \forall \phi \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{2.69}
\end{equation*}
$$

Certainly $A^{*} u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ if $u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ since

$$
\begin{equation*}
A^{*} u(\psi)=u(\overline{A \bar{\psi}}) \text { and } \mathcal{S}\left(\mathbb{R}^{n}\right) \ni \psi \longmapsto \overline{A \bar{\psi}} \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{2.70}
\end{equation*}
$$

is clearly continuous. In terms of Schwartz kernels,
4.32

$$
\begin{align*}
A \phi(x) & =\int K_{A}(x, y) \phi(y) d y, \phi \in \mathcal{S}\left(\mathbb{R}^{n}\right)  \tag{2.71}\\
A^{*} u(x) & =\int K_{A^{*}}(x, y) u(y) d y, u \in \mathcal{S}\left(\mathbb{R}^{n}\right)
\end{align*}
$$

We then see that

$$
4.33
$$

$$
\begin{gather*}
\int K_{A^{*}}(x, y) u(y) \overline{\phi(x)} d y d x=\int \overline{K_{A}(x, y) \phi(y)} d y u(x) d x  \tag{2.72}\\
\Longrightarrow K_{A^{*}}(x, y)=\overline{K_{A}(y, x)}
\end{gather*}
$$

where we are using the uniqueness of Schwartz' kernels.
This proves ( $\overline{2} .41$ ) since

$$
\begin{align*}
\overline{K_{A}(y, x)} & =\overline{\left[\frac{1}{(2 \pi)^{n}} \int e^{i(y-x) \cdot \xi} a(y, x, \xi) d \xi\right]}  \tag{2.73}\\
& =\frac{1}{(2 \pi)^{n}} \int e^{i(x-y) \cdot \xi} \bar{a}(y, x, \xi) d \xi
\end{align*}
$$

i.e. $A^{*}=I(\bar{a}(y, x, \xi))$. Thus one advantage of allowing general operators $\left(\frac{3.26}{2.34}\right)$ is that closure under the passage to adjoint is immediate.

For the composition formula we need to apply Proposition $\frac{4.4}{2.2}$ twice. First to $A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$, to write it with symbol $a(x, \xi)$

$$
\begin{aligned}
A \phi(x) & =(2 \pi)^{-n} \int e^{i(x-y) \cdot \xi} a(x, \xi) \phi(y) d y d \xi \\
& =(2 \pi)^{-n} \int e^{i x \cdot \xi} a(x, \xi) \hat{\phi}(\xi) d \xi
\end{aligned}
$$

Then we also apply Proposition $\frac{4.4}{\frac{4.2}{2}}$ to $B^{*}$,

$$
B^{*} u(x)=(2 \pi)^{-n} \int e^{i x \cdot \xi} \bar{b}(x, \xi) \hat{u}(\xi) d \xi
$$

Integrating this against a test function $\phi \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ gives

$$
\begin{gather*}
\langle B \phi, u\rangle=\left\langle\phi, B^{*} u\right\rangle=(2 \pi)^{-n} \iint e^{-i x \cdot \xi} \phi(x) b(x, \xi) \overline{\hat{u}(\xi)} d \xi d x \\
\Longrightarrow \widehat{B \phi}(\xi)=\int e^{-i y \cdot \xi} b(y, \xi) \phi(y) d y \tag{2.74}
\end{gather*}
$$

Inserting this into the formula for $A \phi$ shows that

$$
\Longrightarrow A B(u)=(2 \pi)^{-n} \int e^{i(x-y) \cdot \xi} a(x, \xi) b(y, \xi) u(y) d y d \xi
$$

Since $a(x, \xi) b(y, \xi) \in S_{\infty}^{m+m^{\prime}}\left(\mathbb{R}_{(x, y)}^{2 n} ; \mathbb{R}_{\xi}^{n}\right)$ this shows that $A B \in \Psi_{\infty}^{m+m^{\prime}}\left(\mathbb{R}^{n}\right)$ as claimed.

### 2.8. Quantization and symbols

So, we have now shown that there is an 'oscillatory integral' interpretation of

$$
\begin{equation*}
K(x, y)=(2 \pi)^{-n} \int e^{i(x-y) \cdot \xi} a(x, y, \xi) d \xi=I(a) \tag{2.75}
\end{equation*}
$$

which defines, for any $w \in \mathbb{R}$, a continuous linear map

$$
I:\left(1+|x-y|^{2}\right)^{\frac{w}{2}} S_{\infty}^{\infty}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right) \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{2 n}\right)
$$

the range of which is the space of pseudodifferential operators on $\mathbb{R}^{n}$;

$$
\begin{gather*}
A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right) \Longleftrightarrow A: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \text { and } \\
\exists w \text { s.t. } K_{A}(x, y)=I(a), a \in\left(1+|x-y|^{2}\right)^{\frac{w}{2}} S_{\infty}^{m}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right) \tag{2.76}
\end{gather*}
$$

Furthermore, we have shown in Proposition $\frac{4.4}{2.2}$ that the special case, $w=0$ and $\partial_{y} a \equiv 0$, gives an isomorphism

$$
\begin{equation*}
\Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right) \underset{q_{L}}{\stackrel{\sigma_{L}}{\rightleftarrows}} S_{\infty}^{m}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right) \tag{2.77}
\end{equation*}
$$

The map here, $q_{L}=I$ on symbols independent of $y$, is the left quantization map and its inverse $\sigma_{L}$ is the left full symbol map. Next we consider some more consequences of this reduction theorem.

As well as the left quantization map leading to the isomorphism $\left(\frac{5.3}{2.3 .7}\right)$ there is a right quantization map, similarly derived from ( $\left.\bar{L} . \frac{1}{2} 5\right)$ :

$$
\begin{equation*}
q_{R}(a)=(2 \pi)^{-n} \int e^{i(x-y) \cdot \xi} a(y, \xi) d \xi, \quad a \in S_{\infty}^{m}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right) \tag{2.78}
\end{equation*}
$$

In fact using the adjoint operator, $*$, on operators and writing as well $*$ for complex conjugation of symbols shows that

$$
\begin{equation*}
q_{R}=* \cdot q_{L} \cdot * \tag{2.79}
\end{equation*}
$$

is also an isomorphism, with inverse $\sigma_{R}$

$$
\begin{equation*}
\Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right) \underset{q_{R}}{\stackrel{\sigma_{R}}{\rightleftarrows}} S_{\infty}^{m}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right) \tag{2.80}
\end{equation*}
$$

These are the two 'extreme' quantization procedures, see Problem ${ }^{\mathbb{P 5} .1}$. for another (more centrist) approach. Using the proof of the reduction theorem we find:
5.7 Lemma 2.4. For any $a \in S_{\infty}^{m}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right)$,

$$
\begin{equation*}
\sigma_{L}\left(q_{R}(a)\right)(x, \xi) \sim \sum_{\alpha} \frac{i^{|\alpha|}}{\alpha!} D_{x}^{\alpha} D_{\xi}^{\alpha} a(x, \xi) \sim e^{i<D_{x}, D_{\xi}>} a . \tag{2.81}
\end{equation*}
$$

For the moment the last asymptotic equality is just to help in remembering the formula, which is the same as given by the formal Taylor series expansion at the origin of the exponential.

Proof. This follows from the general formula ( $\left.\frac{(4.29}{2} .68\right)$.

### 2.9. Principal symbol

One important thing to note from $\left(\frac{5.8}{2.81}\right)$ is that

$$
\begin{equation*}
D_{x}^{\alpha} D_{\xi}^{\alpha} a(x, \xi) \in S_{\infty}^{m-|\alpha|}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right) \tag{2.82}
\end{equation*}
$$

so that for any pseudodifferential operator

$$
\begin{equation*}
A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right) \Longrightarrow \sigma_{L}(A)-\sigma_{R}(A) \in S_{\infty}^{m-1}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right) \tag{2.83}
\end{equation*}
$$

For this reason we consider the general quotient spaces

$$
\begin{equation*}
S_{\infty}^{m-[1]}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)=S_{\infty}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right) / S_{\infty}^{m-1}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right) \tag{2.84}
\end{equation*}
$$

and, for $a \in S_{\infty}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$, write $[a]$ for its image, i.e. equivalence class, in the quotient space $S_{\infty}^{m-[1]}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$. The 'principal symbol map'

$$
\begin{equation*}
\sigma_{m}: \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right) \longrightarrow S_{\infty}^{m-[1]}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right) \tag{2.85}
\end{equation*}
$$

$$
\text { is defined by } \sigma_{m}(A)=\left[\sigma_{L}(A)\right]=\left[\sigma_{R}(A)\right] \text {. }
$$

As distinct from $\sigma_{L}$ or $\sigma_{R}, \sigma_{m}$ depends on $m$, i.e. one needs to know that the order is at most $m$ before it is defined.

The isomorphism $\left(\frac{5.3}{2.7} 7\right)$ is replaced by a weaker (but very useful) exact sequence.

Lemma 2.5. For every $m \in \mathbb{R}$

$$
0 \hookrightarrow \Psi_{\infty}^{m-1}\left(\mathbb{R}^{n}\right) \hookrightarrow \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right) \xrightarrow{\sigma_{m}} S_{\infty}^{m-[1]}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right) \longrightarrow 0
$$

is a short exact sequence (the 'principal symbol sequence' or simply the 'symbol sequence').

Proof. This is just the statement that the range of each ${ }_{3}$ map is the null space of the next i.e. that $\sigma_{m}$ is surjective, which follpws from ( $\frac{5.3}{2.77}$ ), and that the null space of $\sigma_{m}$ is just $\Psi_{\infty}^{m-1}\left(\mathbb{R}^{n}\right)$ and this is again $\left(\frac{5.7}{2.7}\right)$ and the definition of $\sigma_{m}$.

The fundamental result proved above is that

$$
\begin{equation*}
\Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right) \cdot \Psi_{\infty}^{m^{\prime}}\left(\mathbb{R}^{n}\right) \subset \Psi_{\infty}^{m+m^{\prime}}\left(\mathbb{R}^{n}\right) \tag{2.86}
\end{equation*}
$$

In fact we showed that if $A=q_{L}(a), a \in S_{\infty}^{m}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right)$ and $B=q_{R}(b), b \in$ $S_{\infty}^{m^{\prime}}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right)$ then the composite operator has Schwartz kernel

$$
K_{A \cdot B}(x, y)=I(a(x, \xi) b(y, \xi))
$$

Using the formula $\left(\frac{4.29}{2.68}\right)$ again we see that

$$
\begin{equation*}
\sigma_{L}(A \cdot B) \sim \sum_{\alpha} \frac{i^{|\alpha|}}{\alpha!} D_{\xi}^{\alpha}\left[a(x, \xi) D_{x}^{\alpha} b(x, \xi)\right] \tag{2.87}
\end{equation*}
$$

Of course $b=\sigma_{R}(B)$ so we really want to rewrite $\left(\frac{5.12}{2.87}\right)$ in terms of $\sigma_{L}(B)$.

$$
\begin{gather*}
\sigma_{m+m^{\prime}}(A \circ B)=\sigma_{m}(A) \cdot \sigma_{m^{\prime}}(B)  \tag{2.88}\\
\sigma_{L}(A \circ B) \sim \sum_{\alpha} \frac{i^{|\alpha|}}{\alpha!} D_{\xi}^{\alpha} \sigma_{L}(A) \cdot D_{x}^{\alpha} \sigma_{L}(B) \tag{2.89}
\end{gather*}
$$

Proof. The simple formula $\left(\frac{5.14}{2.88}\right)$ is already immediate from $\left(\frac{5.12}{2.87}\right)$ since all terms with $|\alpha| \geq 1$ are of grder $m+m^{\prime}-|\alpha| \leq m .{ }^{5} m^{\prime}-1$. To get the 'full' formula ( $\overline{2} .89$ ) we can insert into ( $\overline{2.87}$ ) the inverse of ( $\overline{2} .81$ ), namely

$$
\sigma_{R}(x, \xi) \sim \sum_{\alpha} \frac{(-i)^{|\alpha|}}{\alpha!} D_{\xi}^{\alpha} D_{x}^{\alpha} \sigma_{L}(x, \xi) \sim e^{-i<D_{x}, D_{\xi}>} \sigma_{L}(x, \xi)
$$

This gives the double sum (still asymptotically convergent)

$$
\sigma_{L}(A \circ B) \sim \sum_{\beta} \sum_{\alpha} \frac{i^{|\alpha|}}{\alpha!} D_{\xi}^{\alpha}\left[\sigma_{L}(A) D_{x}^{\alpha} \frac{i^{\beta \mid}}{\beta!} D_{x}^{\beta} D_{\xi}^{\beta} \sigma_{L}(B)\right]
$$

Setting $\gamma=\alpha+\beta$ this becomes

$$
\sigma_{L}(A \circ B) \sim \sum_{\gamma} \frac{i^{|\gamma|}}{\gamma!} \sum_{0 \leq \alpha \leq \gamma} \frac{\gamma!(-1)^{|\gamma-\alpha|}}{\alpha!(\gamma-\alpha)!} D_{\xi}^{\alpha}\left[\sigma_{L}(A) \times D_{\xi}^{\gamma-\alpha} D_{x}^{\gamma} \sigma_{L}(B)\right]
$$

Then Leibniz' formula shows that this sum over $\alpha$ can be rewritten as

$$
\begin{aligned}
& \sigma_{L}(A \circ B) \sim \sum_{\gamma} \frac{i^{|\gamma|}}{\gamma!} D_{\xi}^{\gamma} \sigma_{L}(A) \cdot D_{x}^{\gamma} \sigma_{L}(B) \\
&\left.\sim e^{i<D_{y}, D_{\xi}>} \sigma_{L}(A)(x, \xi) \sigma_{L}(B)(y, \eta)\right|_{y=x, \eta=\xi}
\end{aligned}
$$

This is just $\left(\frac{5.15}{2.89}\right)$.
The simplicity of $\left(\frac{5.14}{2.88}\right)$ over $\left(\frac{5.15}{2.89}\right)$ is achieved at the expense of enormous loss of information. Still, many problems can be solved using ( 2.88 ) which we can think of as saying that the principal symbol maps give a homomorphism, for instance from the filtered algebra $\Psi_{\infty}^{0}\left(\mathbb{R}^{n}\right)$ to the commutative algebra $S_{\infty}^{0-[1]}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right)$.

### 2.10. Ellipticity

We say that an element of $\Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ is elliptic if it is invertible modulo an error in $\Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$ with the approximate inverse of order $-m$ i.e.

$$
A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right) \text { is elliptic }
$$

$$
\begin{equation*}
\Longleftrightarrow \exists \quad B \in \Psi_{\infty}^{-m}\left(\mathbb{R}^{n}\right) \text { s.t. } \quad A \circ B-\operatorname{Id} \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right) \tag{2.90}
\end{equation*}
$$

Thus ellipticity, here by definition, is invertibility in $\Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right) / \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$, so the inverse lies in $\Psi_{\infty}^{-m}\left(\mathbb{R}^{n}\right) / \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$. The point about ellipticity is that it is a phenomenon of the principal symbol.

### 5.17

5.18
5.19
5.20
5.21
5.22

Theorem 2.3. The following conditions on $A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ are equivalent

$$
\begin{align*}
& \exists[b] \in S_{\infty}^{-m-[1]}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right) \text { s.t. } \sigma_{m}(A) \cdot[b] \equiv 1 \text { in } S_{\infty}^{0-[1]}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right)  \tag{2.92}\\
& \exists b \in S_{\infty}^{-m}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right) \text { s.t. } \sigma_{L}(A) \cdot b-1 \in S_{\infty}^{-\infty}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right)  \tag{2.93}\\
& \quad \exists \epsilon>0 \text { s.t. }\left|\sigma_{L}(A)(x, \xi)\right| \geq \epsilon(1+|\xi|)^{m} \text { in }|\xi|>\frac{1}{\epsilon}
\end{align*}
$$

Proof. We shall show

$$
\begin{equation*}
\left(\frac{5.18}{2.91}\right) \Longrightarrow\left(\frac{15.19}{2.92}\right) \Longrightarrow\left(\frac{5.20}{2.93}\right) \Longleftrightarrow\left(\frac{5.21}{2.94}\right) \Longrightarrow\left(\frac{5.18}{2.91}\right) \tag{2.95}
\end{equation*}
$$

In fact Lemma $\frac{3.21}{2.2}$ shows the equivalence of $\left(\frac{5.20}{2.93}\right)$ and $\left(\frac{5.21}{(2.94)}\right.$. Since we know that $\sigma_{0}(\mathrm{Id})=1$ applying the identity $\left(\frac{5.14}{2.88}\right)$ to the definition of ellipticity in $\left(\frac{5.16}{2.90}\right)$ gives

$$
\begin{equation*}
\sigma_{m}(A) \cdot \sigma_{-m}(B) \equiv 1 \text { in } S_{\infty}^{0-[1]}\left(\mathbb{R}^{n}, \mathbb{R}^{n}\right) \tag{2.96}
\end{equation*}
$$

i.e. that $\left(\frac{5.18}{2.91}\right) \Longrightarrow\left(\frac{5.19}{2.92}\right)$.

Now assuming $\binom{5.44}{2.96}$ (i.e. $\left(\frac{5.19}{2.92}\right)$ ), and recalling that $\sigma_{m}(A)=\left[\sigma_{L}(A)\right]$ we find that a representative $b_{1}$ of the class $[b]$ must satisfy

$$
\begin{equation*}
\sigma_{L}(A) \cdot b_{1}=1+e_{1}, \quad e_{1} \in S_{\infty}^{-1}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right) \tag{2.97}
\end{equation*}
$$

this being the meaning of the equality of residue classes. Now for the remainder, $e_{1} \in S_{\infty}^{-1}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right)$, the Neumann series

$$
\begin{equation*}
f \sim \sum_{j \geq 1}(-1)^{j} e_{1}^{j} \tag{2.98}
\end{equation*}
$$

is asymptotically convergent, so $f \in S_{\infty}^{-1}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right)$ exists, and
(2.99)

$$
(1+f) \cdot\left(1+e_{1}\right)=1+e_{\infty}, \quad e_{\infty} \in S_{\infty}^{-\infty}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right)
$$

Then multiplying $\left(\frac{5.23}{2.97}\right)$ by $(1+f)$ gives

$$
\begin{equation*}
\sigma_{L}(A) \cdot\left\{b_{1}(1+f)\right\}=1+e_{\infty} \tag{2.100}
\end{equation*}
$$

which proves $\left(\frac{5.20}{2.93}\right)$, since $b=b_{1}(1+f) \in S_{\infty}^{-m}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right)$. Of course

$$
\begin{equation*}
\sup (1+|\xi|)^{N}\left|e_{\infty}\right|<\infty \quad \forall N \tag{2.101}
\end{equation*}
$$

SO

$$
\begin{equation*}
\exists C \text { s.t. }\left|e_{\infty}(x, \xi)\right|<\frac{1}{2} \text { in }|\xi|>C . \tag{2.102}
\end{equation*}
$$

From $\left(\frac{5.25}{2.100}\right)$ this means

$$
\begin{equation*}
\left|\sigma_{L}(A)(x, \xi)\right| \cdot|b(x, \xi)| \geq \frac{1}{2},|\xi|>C \tag{2.103}
\end{equation*}
$$

Since $|b(x, \xi)| \leq C(1+|\xi|)^{-m}$ (being a symbol of order $-m$ ), ( $\frac{5.26}{2.10} 3$ ) implies

$$
\begin{equation*}
\inf _{|\xi| \geq C}\left|\sigma_{L}(A)(x, \xi)\right|(1+|\xi|)^{-m} \geq C>0 \tag{2.104}
\end{equation*}
$$

which shows that $\left(\frac{5.20}{2.93}\right)$ implies $\left(\frac{5.21}{2.94}\right)$
Conversely, as already remarked, $\left(\frac{5.21}{2.94}\right)$ implies $\left(\frac{5.20}{2.93}\right)$
Now suppose ( $\left(\frac{5.20}{2.93}\right)$ holds. Set $B_{1}=q_{L}(b)$ then from $\left(\frac{5.14}{Z .88}\right)$ again

$$
\begin{equation*}
\sigma_{0}\left(A \circ B_{1}\right)=\left[q_{m}(A)\right] \cdot[b] \equiv 1 \tag{2.105}
\end{equation*}
$$

That is,

$$
\begin{equation*}
A \circ B_{1}-\operatorname{Id}=E_{1} \in \Psi_{\infty}^{-1}\left(\mathbb{R}^{n}\right) \tag{2.106}
\end{equation*}
$$

Consider the Neumann series of operators

$$
\begin{equation*}
\sum_{j \geq 1}(-1)^{j} E_{1}^{j} \tag{2.107}
\end{equation*}
$$

The corresponding series of (left-reduced) symbols is asymptotically summable so we can choose $F \in \Psi_{\infty}^{-1}\left(\mathbb{R}^{n}\right)$ with

$$
\begin{equation*}
\sigma_{L}(F) \sim \sum_{j \geq 1}(-1)^{j} \sigma_{L}\left(E_{1}^{j}\right) \tag{2.108}
\end{equation*}
$$

Then
(2.109)

$$
\left(\operatorname{Id}+E_{1}\right)(\operatorname{Id}+F)=\operatorname{Id}+E_{\infty}, E_{\infty} \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)
$$

Thus $B=B_{1}(\operatorname{Id}+F) \in \Psi_{\infty}^{-m}\left(\mathbb{R}^{n}\right)$ satisfies $\left(\frac{5.16}{2.90}\right)$ and it follows that $A$ is elliptic.
In the definition of ellipticity in $\left(\frac{5.16}{2.90}\right)$ we have taken $B$ to be a 'right parametrix', i.e. a right inverse modulo $\Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$. We can just as well take it to be a left parametrix.

Lemma 2.7. $A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ is elliptic if and only if there exists $B^{\prime} \in \Psi_{\infty}^{-m}\left(\mathbb{R}^{n}\right)$ such that
5.32

$$
\begin{equation*}
B^{\prime} \circ A=\operatorname{Id}+E^{\prime}, E^{\prime} \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right) \tag{2.110}
\end{equation*}
$$

and then if $B$ satisfies $\left(\frac{5.16}{2.90}\right), B-B^{\prime} \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$.
Proof. Certainly $\left(\frac{5.32}{5.119} 0\right)$ implies $\sigma_{-m}\left(B^{\prime}\right) \cdot \sigma_{m}(A) \equiv 1$, and the multiplication here is commutatiye so ( 2.92 ) holds and $A$ is elliptic. Conversely if $A$ is elliptic we get in place of ( $\left(\frac{5.29}{2.106)}\right.$

$$
B_{1} \circ A-\mathrm{Id}=E_{1}^{\prime} \in \Psi_{\infty}^{-1}\left(\mathbb{R}^{n}\right)
$$

Then defining $F^{\prime}$ as in $\left(\frac{5.45}{2.10} 8\right)$ with $E_{1}^{\prime}$ in place ${ }_{15} \mathrm{of}_{32} E_{1}$ we get $\left(\operatorname{Id}+F^{\prime}\right)\left(\operatorname{Id}+E_{1}^{\prime}\right)=$ $\operatorname{Id} .+E_{\infty}^{\prime}$ and then $B^{\prime}=\left(\operatorname{Id}+F^{\prime}\right) \circ B_{1}$ satisfies (L.110). Thus 'left' ellipticity as in ( $\frac{5.32}{2.110)}$ is equivalent to right ellipticity. Applying $B$ to ( ( 2.110$)$ gives

$$
\begin{equation*}
B^{\prime} \circ(\mathrm{Id}+E)=B^{\prime} \circ(A \circ B)=\left(\operatorname{Id}+E^{\prime}\right) \circ B \tag{2.111}
\end{equation*}
$$

which shows that $B-B^{\prime} \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$.
Thus a left parametrix of an elliptic element of $\Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ is always a right, hence two-sided, parametrix and such a parametrix is unique up to an additive term in $\Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$.

### 2.11. Elliptic regularity

One of the main reasons that the 'residual' terms are residual is that they are smoothing operators.

Lemma 2.8. If $E \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$ then

$$
\begin{equation*}
E: \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \cap \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right) \tag{2.112}
\end{equation*}
$$

Proof. This follows from Proposition $\frac{\frac{4.21}{2.4} \text { since we can regard the kernel as a }}{\text { a }}$ $\mathcal{C}^{\infty}$ function of $x$ taking values in $\mathcal{S}\left(\mathbb{R}_{y}^{n}\right)$.

Directly from the existence of parametrices for elliptic operators we can deduce the regularity of solutions to elliptic (pseudodifferential) equations.
20.2.1998.103

$$
\begin{equation*}
u=(B A-E) u=-E u \tag{2.113}
\end{equation*}
$$

and the conclusion follows from Lemma $\frac{5.42}{2.8 .}$

### 2.12. The Laplacian

Suppose that $g_{i j}(x)$ are the components of an ' $\infty$-metric' on $\mathbb{R}^{n}$, i.e.

$$
g_{i j}(x) \in \mathcal{C}_{\infty}^{\infty}\left(\mathbb{R}^{n}\right), i, j=1, \ldots, n
$$

$$
\begin{equation*}
\left|\sum_{i, j=1}^{n} g_{i j}(x) \xi_{i} \xi_{j}\right| \geq \epsilon|\xi|^{2} \quad \forall x \in \mathbb{R}^{n}, \xi \in \mathbb{R}^{n}, \epsilon>0 \tag{2.114}
\end{equation*}
$$

The Laplacian of the metric is the second order differential operator

$$
\begin{equation*}
\Delta_{g}=\sum_{i, j=1}^{n} \frac{1}{\sqrt{g}} D_{x_{i}} g^{i j} \sqrt{g} D_{x_{j}} \tag{2.115}
\end{equation*}
$$

where

$$
g(x)=\operatorname{det} g^{i j}(x), g^{i j}(x)=\left(g_{i j}(x)\right)^{-1}
$$

The Laplacian is determined by the integration by parts formula

$$
\begin{equation*}
\int_{\mathbb{R}^{n}} \sum_{i, j} g^{i j}(x) D_{x_{i}} \phi \cdot \overline{D_{x_{j}} \psi} d g=\int \Delta_{g} \phi \cdot \bar{\psi} d g \forall \phi, \psi \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{2.116}
\end{equation*}
$$

where
5.38

$$
\begin{equation*}
d g=\sqrt{g} d x \tag{2.117}
\end{equation*}
$$

Our assumption in $\left(\frac{5.34}{2.11} 4\right)$ shows that $\Delta=\Delta_{g} \in \operatorname{Diff}_{\infty}^{2}\left(\mathbb{R}^{n}\right) \subset \Psi_{\infty}^{2}\left(\mathbb{R}^{n}\right)$ is in fact elliptic, since

$$
\begin{equation*}
\sigma_{2}(\Delta)=\sum_{i, j=1} g^{i j} \xi_{i} \xi_{j} \tag{2.118}
\end{equation*}
$$

Thus $\Delta$ has a two-sided parametrix $B \in \Psi_{\infty}^{-2}\left(\mathbb{R}^{n}\right)$

$$
\begin{equation*}
\Delta \circ B \equiv B \circ \Delta \equiv \operatorname{Id} \quad \bmod \quad \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right) \tag{2.119}
\end{equation*}
$$

In particular we see from Proposition $\frac{20.2 .1998 .103}{2.5 \text { that } \Delta u}=0, u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ implies $u \in$ $\mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right)$.
S.L2.boundedness

### 2.13. $L^{2}$ boundedness

So far we have thought of pseudodifferential operators, the elements of $\Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ for some $m$, as defining continuous linear operators on $\mathcal{S}\left(\mathbb{R}^{n}\right)$ and, by duality, on $\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$. Now that we have proved the composition formula we can use it to prove other 'finite order' regularity results. The basic one of these is $L^{2}$ boundedness:
$6.1 \quad$ Proposition 2.6. [Boundedness] If $A \in \Psi_{\infty}^{0}\left(\mathbb{R}^{n}\right)$ then, by continuity from $\mathcal{S}\left(\mathbb{R}^{n}\right)$, A defines a bounded linear operator

$$
\begin{equation*}
A: L^{2}\left(\mathbb{R}^{n}\right) \longrightarrow L^{2}\left(\mathbb{R}^{n}\right) \tag{2.120}
\end{equation*}
$$

Our proof will be in two stages, the first part is by direct estimation. Namely, Schur's lemma ${ }^{4}$ gives a useful criterion for an integral operator to be bounded on $L^{2}$ 。

$$
\begin{equation*}
\sup _{x \in \mathbb{R}^{n}} \int_{\mathbb{R}^{n}}|K(x, y)| d y, \sup _{y \in \mathbb{R}^{n}} \int_{\mathbb{R}^{n}}|K(x, y)| d x<\infty \tag{2.121}
\end{equation*}
$$

then the operator $K: \phi \longmapsto \int_{\mathbb{R}^{n}} K(x, y) \phi(y) d y$ is bounded on $L^{2}\left(\mathbb{R}^{n}\right)$.
Proof. Since $\mathcal{S}\left(\mathbb{R}^{n}\right)$ is dense in $L^{2}\left(\mathbb{R}^{n}\right)$ (see Problem $\frac{1.2 .2000 .280}{2.18)}$ we only need to show the existence of a constant, $C$, such that
6.4

$$
\begin{equation*}
\int|K \phi(x)|^{2} d x \leq C \int|\phi|^{2} \quad \forall \phi \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{2.122}
\end{equation*}
$$

Writing out the integral on the left
6.5

$$
\begin{gather*}
\int\left|\int K(x, y) \phi(y) d y\right|^{2} d x \\
=\iiint K(x, y) \overline{K(x, z)} \phi(y) \overline{\phi(z)} d y d z d x \tag{2.123}
\end{gather*}
$$

is certainly absolutely convergent and

$$
\begin{aligned}
& \int|K \phi(x)|^{2} d x \\
& \leq\left(\left.\iiint|K(x, y) K(x, z)| \phi(y)\right|^{2} d y d x d z\right)^{\frac{1}{2}} \\
& \times\left(\left.\iiint|K(x, y) K(x, z)| \phi(z)\right|^{2} d z d x d y\right)^{\frac{1}{2}}
\end{aligned}
$$

[^6]These two factors are the same. Since

$$
\int|K(x, y)||K(x, z)| d x d z \leq \sup _{x \in \mathbb{R}^{n}} \int|K(x, z)| d z \cdot \sup _{y \in \mathbb{R}^{n}} \int \mid K(x, y \mid d x
$$

$\left(\frac{6.4}{(2.1} 22\right)$ follows. Thus $\left(\frac{5.51}{2.121}\right)$ gives ( $\left(\frac{6.4}{2.1} 22\right)$.
This standard lemma immediately implies the $L^{2}$ boundedness of the 'residual terms.' Thus, if $K \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$ then its kernel satisfies (立.60). This in particular implies

$$
|K(x, y)| \leq C(1+|x-y|)^{-n-1}
$$

and hence that $K$ satisfies $\left(\frac{5.51}{2.12} 1\right)$. Thus each $K \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$ is bounded on $L^{2}\left(\mathbb{R}^{n}\right)$.

### 2.14. Square root

To prove the general result, ( $\frac{16.2}{2.2} 20$ ), we shall use the clever idea, due to Hörmander, of using the (approximate) square root of an operator. We shall say that an element $[a] \in S_{\infty}^{m-[1]}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right)$ is positive if there is some $0<a \in S^{m}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right)$ in the equivalence class.

Proposition 2.7. Suppose $A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$, $m>0$, is self-adjoint, $A=A^{*}$, and elliptic with a positive principal symbol, then there exists $B \in \Psi_{\infty}^{m / 2}\left(\mathbb{R}^{n}\right), B=B^{*}$, such that

$$
\begin{equation*}
A=B^{2}+G, \quad G \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right) \tag{tabular}
\end{equation*}
$$

Proof. This is a good exercise in the use of the symbol calculus. Let $a \in$ $S_{\infty}^{m}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right), q_{11}>2.2000 .28$ positive representative of the principal symbol of $A$. Now (See Problem 2.19 for an outline of the proof)

$$
\begin{equation*}
b_{0}=a^{\frac{1}{2}} \in S_{\infty}^{m / 2}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right) \tag{2.126}
\end{equation*}
$$

Let $B_{0} \in \Psi_{\infty}^{m / 2}\left(\mathbb{R}^{n}\right)$ have principal symbol $b_{0}$. We can assume that $B_{0}=B_{0}^{*}$, since if not we just replace $B_{0}$ by $\frac{1}{2}\left(B_{0}+B_{0}^{*}\right)$ which has the same principal symbol.

The symbol calculus shows that $B_{0}^{2} \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ and

$$
\sigma_{m}\left(B_{0}^{2}\right)=\left(\sigma_{m / 2}\left(B_{0}\right)\right)^{2}=b_{0}^{2}=a_{0} \quad \bmod S_{\infty}^{m-1}
$$

Thus
6.9

$$
\begin{equation*}
A-B_{0}^{2}=E_{1} \in \Psi_{\infty}^{m-1}\left(\mathbb{R}^{n}\right) \tag{2.127}
\end{equation*}
$$

Then we proceed inductively. Suppose we have chosen $B_{j} \in \Psi_{\infty}^{m / 2-j}\left(\mathbb{R}^{n}\right)$, with $B_{j}^{*}=B_{j}$, for $j \leq N$ such that

$$
A-\left(\sum_{j=0}^{N} B_{j}\right)^{2}=E_{N+1} \in \Psi_{\infty}^{m-N-1}\left(\mathbb{R}^{n}\right)
$$

Of course we have done this for $N=0$. Then see the effect of adding $B_{N+1} \in$ $\Psi_{\infty}^{m / 2-N-1}\left(\mathbb{R}^{n}\right)$ :
6.11

$$
\begin{align*}
A-\left(\sum_{j=0}^{N+1} B_{j}\right)^{2}=E_{N+1}-\left(\sum_{j=0}^{N} B_{j}\right) & B_{N+1}  \tag{2.129}\\
& -B_{N+1}\left(\sum_{j=0}^{N} B_{j}\right)-B_{N+1}^{2}
\end{align*}
$$

On the right side all terms are of order $m-N-2$, except for
6.12

$$
\begin{equation*}
E_{N+1}-B_{0} B_{N+1}-B_{N+1} B_{0} \in \Psi_{\infty}^{m-N-1}\left(\mathbb{R}^{n}\right) \tag{2.130}
\end{equation*}
$$

The principal symbol, of order $m-N-1$, of this is just

$$
\begin{equation*}
\sigma_{m-N-1}\left(E_{N+1}\right)-2 b_{0} \cdot \sigma_{\frac{m}{2}-N-1}\left(B_{N+1}\right) \tag{2.131}
\end{equation*}
$$

Thus if we choose

$$
\sigma_{m / 2-N-1}\left(B_{N+1}\right)=\frac{1}{2} \frac{1}{b_{0}} \cdot \sigma_{m-N-1}\left(E_{N+1}\right)
$$

and replace $B_{N+1}$ by $\frac{1}{2}\left(B_{N_{1} 0^{1}}+B_{N+1}^{*}\right)$, we get the inductive hypothesis for $N+1$. Thus we have arranged ( $\overline{2} .128)$ for every $N$. Now define $B=\frac{1}{2}\left(B^{\prime}+\left(B^{\prime}\right)^{*}\right)$ where

$$
\begin{equation*}
\sigma_{L}\left(B^{\prime}\right) \sim \sum_{j=0}^{\infty} \sigma_{L}\left(B_{j}\right) \tag{2.132}
\end{equation*}
$$

Since all the $B_{j}$ are self-adjoint $B$ also satisfies $\left(\frac{6.14}{2.132}\right)$ and from $\left(\frac{6.10}{2.12} 8\right)$

$$
\begin{equation*}
A-B^{2}=A-\left(\sum_{j=0}^{N} B_{j}+B_{(N+1)}\right)^{2} \in \Psi_{\infty}^{m-N-1}\left(\mathbb{R}^{n}\right) \tag{2.133}
\end{equation*}
$$

for every $N$, since $B_{(N+1)}=B-\sum_{j=0}^{N} B_{j} \in \Psi_{\infty}^{m / 2-N-1}\left(\mathbb{R}^{n}\right)$. Thus $A-B^{2} \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$ and we have proved ( $\left.\frac{6.7}{\left(\frac{6}{2} .7\right.} 25\right)$, and so Proposition $\frac{6.6}{\frac{6.7}{2.7}}$.

### 2.15. Proof of Boundedness

Here is Hörmander's argument. We want to show that

$$
\begin{equation*}
\|A \phi\| \leq C\|\phi\| \quad \forall \phi \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{2.134}
\end{equation*}
$$

where $A \in \Psi_{\infty}^{0}\left(\mathbb{R}^{n}\right)$. The square of the left side can be written

$$
\int A \phi \cdot \overline{A \phi} d x=\int \phi \cdot \overline{\left(A^{*} A \phi\right)} d x
$$

So it suffices to show that
6.16

$$
\begin{equation*}
\left\langle\phi, A^{*} A \phi\right\rangle \leq C\|\phi\|^{2} \tag{2.135}
\end{equation*}
$$

Now $A^{*} A \in \Psi_{\infty}^{0}\left(\mathbb{R}^{n}\right)$ with $\sigma_{0}\left(A^{*} A\right)=\overline{\sigma_{0}(A)} \sigma_{0}(A) \in \mathbb{R}$. If $C>0$ is a large constant,

$$
C>\sup _{x, \xi}\left|\sigma_{L}\left(A^{*} A\right)(x, \xi)\right|
$$

then $C-A^{*} A$ has ${ }_{6}$ positive representative of its principal symbol. We can therefore apply Proposition $\frac{2.6}{2.7}$ to it:
6.17

$$
\begin{equation*}
C-A^{*} A=B^{*} B+G, \quad G \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right) \tag{2.136}
\end{equation*}
$$

This gives

$$
\begin{align*}
\left\langle\phi, A^{*} A \phi\right\rangle & =C\langle\phi, \phi\rangle-\left\langle\phi, B^{*} B \phi\right\rangle-\langle\phi, G \phi\rangle \\
& =C\|\phi\|^{2}-\|B \phi\|^{2}-\langle\phi, G \phi\rangle . \tag{2.137}
\end{align*}
$$

The second term on the right is negative and, since $G \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$, we can use the residual case discussed above to conclude that

$$
|\langle\phi, G \phi\rangle| \leq C^{\prime}\|\phi\|^{2} \Longrightarrow\|A \phi\|^{2} \leq C\|\phi\|^{2}+C^{\prime}\|\phi\|^{2}
$$

so ( $\left.\frac{6.2}{2.120}\right)$ holds and Proposition $\frac{6.1}{2.6}$ is proved.

### 2.16. Sobolev boundedness

Using the basic boundedness result, Proposition $\frac{16.1}{2.6}$, and the calculus of pseudodifferential operators we can prove more general results on the action of pseudodifferential operators on Sobolev spaces.

Recall that for any positive integer, $k$,

$$
\begin{equation*}
H^{k}\left(\mathbb{R}^{n}\right)=\left\{u \in L^{2}\left(\mathbb{R}^{n}\right) ; D^{\alpha} u \in L^{2}\left(\mathbb{R}^{n}\right) \forall|\alpha| \leq k\right\} \tag{2.138}
\end{equation*}
$$

Using the Fourier transform we find

$$
\begin{equation*}
u \in H^{k}\left(\mathbb{R}^{n}\right) \Longrightarrow \xi^{\alpha} \hat{u}(\xi) \in L^{2}\left(\mathbb{R}^{n}\right) \quad \forall|\alpha| \leq k \tag{2.139}
\end{equation*}
$$

Now these finitely many conditions can be written as just the one condition

$$
\begin{equation*}
\left(1+|\xi|^{2}\right)^{k / 2} \hat{u}(\xi) \in L^{2}\left(\mathbb{R}^{n}\right) \tag{2.140}
\end{equation*}
$$

Notice that $a(\xi)=\left(1+|\xi|^{2}\right)^{k / 2}=\langle\xi\rangle^{k} \in S_{\infty}^{k}\left(\mathbb{R}^{n}\right)$. Here we use the notation

$$
\begin{equation*}
\langle\xi\rangle=\left(1+|\xi|^{2}\right)^{\frac{1}{2}} \tag{2.141}
\end{equation*}
$$

for a smooth (symbol) of the size of $1+|\xi|$, thus $\left(\frac{6.21}{2.140}\right)$ just says

$$
\begin{equation*}
u \in H^{k}\left(\mathbb{R}^{n}\right) \Longleftrightarrow\langle D\rangle^{k} u \in L^{2}\left(\mathbb{R}^{n}\right) \tag{2.142}
\end{equation*}
$$

For negative integers

$$
6.24
$$

$$
\begin{equation*}
H^{k}\left(\mathbb{R}^{n}\right)=\left\{u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) ; u=\sum_{|\beta| \leq-k} D^{\beta} u_{\beta}, u_{\beta} \in L^{2}\left(\mathbb{R}^{n}\right)\right\},-k \in \mathbb{N} \tag{2.143}
\end{equation*}
$$

The same sort of discussion applies, showing that

$$
\begin{equation*}
u \in H^{k}\left(\mathbb{R}^{n}\right) \Longleftrightarrow\langle D\rangle^{k} u \in L^{2}\left(\mathbb{R}^{n}\right) \tag{2.144}
\end{equation*}
$$

In view of this we define the Sobolev space $H^{m}\left(\mathbb{R}^{n}\right)$, for any real order, by

$$
\begin{equation*}
u \in H^{m},\left(\mathbb{R}^{n}\right) \Longleftrightarrow\langle D\rangle^{m} u \in L^{2}\left(\mathbb{R}^{n}\right) \tag{2.145}
\end{equation*}
$$

It is a Hilbert space with

$$
\begin{equation*}
\|u\|_{m}^{2}=\left\|\langle D\rangle^{m} u\right\|_{L^{2}}^{2}=\int\left(1+|\xi|^{2}\right)^{m}|\hat{u}(\xi)|^{2} d \xi \tag{2.146}
\end{equation*}
$$

Clearly we have
6.28

$$
\begin{equation*}
H^{m}\left(\mathbb{R}^{n}\right) \supseteq H^{m^{\prime}}\left(\mathbb{R}^{n}\right) \text { if } m^{\prime} \geq m \tag{2.147}
\end{equation*}
$$

Notice that it is rather unfortunate that these spaces get smaller as $m$ gets bigger, as opposed to the spaces $\Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ which get bigger with $m$. Anyway that's life and we have to think of

$$
\begin{cases}H^{\infty}\left(\mathbb{R}^{n}\right)=\bigcap_{m} H^{m}\left(\mathbb{R}^{n}\right) & \text { as the residual space }  \tag{2.148}\\ H^{-\infty}\left(\mathbb{R}^{n}\right)=\bigcup_{m} H^{m}\left(\mathbb{R}^{n}\right) & \text { as the big space. }\end{cases}
$$

It is important to note that

$$
\begin{equation*}
\mathcal{S}\left(\mathbb{R}^{n}\right) \nsubseteq H^{\infty}\left(\mathbb{R}^{n}\right) \varsubsetneqq H^{-\infty}\left(\mathbb{R}^{n}\right) \varsubsetneqq \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \tag{2.149}
\end{equation*}
$$

In particular we do not capture all the tempered distributions in $H^{-\infty}\left(\mathbb{R}^{n}\right)$. We therefore consider weighted versions of these Sobolev spaces:

### 6.31

### 6.33

Theorem 2.4. For each $q, m, M \in \mathbb{R}$ each $A \in \Psi_{\infty}^{M}\left(\mathbb{R}^{n}\right)$ defines a continuous linear map

$$
\begin{equation*}
A:\langle x\rangle^{q} H^{m}\left(\mathbb{R}^{n}\right) \longrightarrow\langle x\rangle^{q} H^{m-M}\left(\mathbb{R}^{n}\right) \tag{2.151}
\end{equation*}
$$

Proof. Let us start off with $q=0$, so we want to show that

Now from $\left(\frac{6.26}{2.145)}\right.$ we see that
6.35 (2.153)

$$
\begin{aligned}
u & \in H^{m}\left(\mathbb{R}^{n}\right) \Longleftrightarrow\langle D\rangle^{m} u \in L^{2}\left(\mathbb{R}^{n}\right) \\
& \Longleftrightarrow\langle D\rangle^{m-M}\langle D\rangle^{M} u \in L^{2}\left(\mathbb{R}^{n}\right) \Longleftrightarrow\langle D\rangle^{M} u \in H^{m-M}\left(\mathbb{R}^{n}\right) \quad \forall m, M
\end{aligned}
$$

That is,
6.36 (2.154)

$$
\langle D\rangle^{M}: H^{m}\left(\mathbb{R}^{n}\right) \longleftrightarrow H^{m-M}\left(\mathbb{R}^{n}\right) \quad \forall m, M
$$

To prove ( $\frac{6.34}{2.152}$ ) it suffices to show that

$$
\begin{equation*}
B=\langle D\rangle^{-M+m} \cdot A \cdot\langle D\rangle^{-m}: L^{2}\left(\mathbb{R}^{n}\right) \longrightarrow L^{2}\left(\mathbb{R}^{n}\right) \tag{2.155}
\end{equation*}
$$

since then $A=\langle D\rangle^{-m+M} \cdot B \cdot\langle D\rangle^{m} \operatorname{maps} H^{m}\left(\mathbb{R}^{n}\right)$ to $H^{m-M}\left(\mathbb{R}^{n}\right)$ :


Since $B \in \Psi_{\infty}^{0}\left(\mathbb{R}^{n}\right)$, by the ${ }^{6034}$ position theorem, we already know ${ }^{6.37}\left(\frac{6.37}{2.155} 5\right)$.
Thus we have proved ( $\overline{2.15} 2)$. To prove the general case, ( $\overline{2.15} 1$ ), we proceed in the same spirit. Thus.33 $x\rangle^{q}$ is an isomorphism from $H^{m}\left(\mathbb{R}^{n}\right)$ to $\langle x\rangle^{q} H^{m}\left(\mathbb{R}^{n}\right)$, by definition. So to get ( $\overline{2} .151$ ) we need to show that

$$
\begin{equation*}
Q=\langle x\rangle^{-q} \cdot A \cdot\langle x\rangle^{q}: H^{m}\left(\mathbb{R}^{n}\right) \longrightarrow H^{m-M}\left(\mathbb{R}^{n}\right) \tag{2.157}
\end{equation*}
$$

i.e. satisfies $\left(\frac{6.34}{2.152)}\right.$. Consider the Schwartz kernel of $Q$. Writing $A$ in left-reduced form, with symbol $a$,

$$
\begin{equation*}
K_{Q}(x, y)=(2 \pi)^{-n} \int e^{i(x-y) \cdot \xi}\langle x\rangle^{-q} a(x, \xi) d \xi \cdot\langle y\rangle^{q} \tag{2.158}
\end{equation*}
$$

Now if we check that

$$
\begin{equation*}
\langle x\rangle^{-q}\langle y\rangle^{q} a(x, \xi) \in\left(1+|x-y|^{2}\right)^{\frac{|q|}{2}} S_{\infty}^{M}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right) \tag{2.159}
\end{equation*}
$$

then we know that $Q \in \Psi_{\infty}^{M}\left(\mathbb{R}^{n}\right)$ and we get $\left(\frac{6.39}{2.15} 7\right)$ from ( $\left.\frac{6.34}{2.15} 2\right)$. Thus we want to show that

$$
\begin{equation*}
\langle x-y\rangle^{-|q|} \frac{\langle y\rangle^{q}}{\langle x\rangle^{q}} a(x, \xi) \in S_{\infty}^{M}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right) \tag{2.160}
\end{equation*}
$$

assuming of course that $a(x, \xi) \in S_{\infty}^{M}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right)$. By interchanging the variables $x$ and $y$ if necessary we can assume that $q<0$. Consider separately the two regions

$$
\begin{align*}
& \left\{(x, y) ;|x-y|<\frac{1}{4}(|x|+|y|)\right\}=\Omega_{1}  \tag{2.161}\\
& \left\{(x, y) ;|x-y|>\frac{1}{8}(|x|+|y|)\right\}=\Omega_{2} .
\end{align*}
$$

In $\Omega_{1}, x$ is "close" to $y$, in the sense that

$$
\begin{equation*}
|x| \leq|x-y|+|y| \leq \frac{1}{4}(|x|+|y|)+|y| \Longrightarrow|x| \leq \frac{4}{3} \cdot \frac{5}{4}|y| \leq 2|y| \tag{2.162}
\end{equation*}
$$

Thus

$$
6.44
$$

$$
\begin{equation*}
\langle x-y\rangle^{-q} \cdot \frac{\langle x\rangle^{-q}}{\langle y\rangle^{-q}} \leq C \quad \text { in } \Omega_{1} \tag{2.163}
\end{equation*}
$$

On the other hand in $\Omega_{2}$,

$$
\begin{equation*}
|x|+|y|<8|x-y| \Longrightarrow|x|<8|x-y| \tag{2.164}
\end{equation*}
$$

so again

$$
\begin{equation*}
\langle x-y\rangle^{-q} \frac{\langle x\rangle^{-q}}{\langle y\rangle^{-q}} \leq C \tag{2.165}
\end{equation*}
$$

In fact we easily conclude that

$$
\begin{equation*}
\langle x-y\rangle^{-q} \frac{\langle y\rangle^{q}}{\langle x\rangle^{q}} \in \mathcal{C}_{\infty}^{\infty}\left(\mathbb{R}^{n}\right) \quad \forall q, \tag{2.166}
\end{equation*}
$$

since differentiation by $x$ or $y$ make all terms "smaller". This proves ( $\frac{6.42}{2.160 \text { ), hence }}$ $\left(\frac{2.41}{2.159)}\right.$ and $\left(\frac{0.39}{2.157)}\right.$ and therefore $\left(\frac{2.15}{2.15}\right)$, i.e. the theorem is proved.

### 2.17. Consequences

We can capture any tempered distribution in a weighted Sobolev space; this is really Schwartz' representation theorem which says that any $u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ is of the form

$$
\begin{equation*}
u=\sum_{\text {finite }} x^{\alpha} D_{x}^{\beta} u_{\alpha \beta}, u_{\alpha \beta} \text { bounded and continuous. } \tag{2.167}
\end{equation*}
$$

Clearly $\mathcal{C}_{\infty}^{0}\left(\mathbb{R}^{n}\right) \subset\langle x\rangle^{1+n} L^{2}\left(\mathbb{R}^{n}\right)$. Thus as a special case of Theorem $\frac{6.32}{2.4,}$

$$
D_{x}^{\beta}:\langle x\rangle^{1+n} L^{2}\left(\mathbb{R}^{n}\right) \longrightarrow\langle x\rangle^{1+n} H^{-|\beta|}\left(\mathbb{R}^{n}\right)
$$

so
Lemma 2.10.

$$
\begin{equation*}
\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)=\bigcup_{M}\langle x\rangle^{M} H^{-M}\left(\mathbb{R}^{n}\right) \tag{2.168}
\end{equation*}
$$

The elliptic regularity result we found before can now be refined:

$$
\begin{gather*}
A u \in\langle x\rangle^{p} H^{q}\left(\mathbb{R}^{n}\right), u \in\langle x\rangle^{p^{\prime}} H^{q^{\prime}}\left(\mathbb{R}^{n}\right) \\
\Longrightarrow u \in\langle x\rangle^{p^{\prime \prime}} H^{q^{\prime \prime}}\left(\mathbb{R}^{n}\right), p^{\prime \prime}=\max \left(p, p^{\prime}\right), q^{\prime \prime}=\max \left(q+m, q^{\prime}\right) . \tag{2.169}
\end{gather*}
$$

Proof. The existence of a left parametrix for $A, B \in \Psi_{\infty}^{-m}\left(\mathbb{R}^{n}\right)$,

$$
B \cdot A=\operatorname{Id}+G, G \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)
$$

means that

$$
\begin{equation*}
u=B(A u)+G u \in\langle x\rangle^{p} H^{q+m}\left(\mathbb{R}^{n}\right)+\langle x\rangle^{p^{\prime}} H^{\infty}\left(\mathbb{R}^{n}\right) \subset\langle x\rangle^{p^{\prime \prime}} H^{q+m}\left(\mathbb{R}^{n}\right) \tag{2.170}
\end{equation*}
$$

### 2.18. Polyhomogeneity

So.far we have been considering operators $A \in \Psi_{q \mathcal{S}}^{m}\left(\mathbb{R}^{n}\right)$ which correspond, via $\left(\frac{3.2}{2.2}\right)$, to amplitudes satisfying the symbol estimates ( $\left.\overline{2} .6\right)$, i.e., in $S_{\infty}^{m}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right)$. As already remarked, there are many variants of these estimates and corresponding spaces of pseudodifferential operators. Some weakening ofthe estimates is discussed in the problems below, starting with Problem $\frac{1.16 \text {. Here we consider a restriction }}{}$ of the spaces, in that we define

$$
\begin{equation*}
S_{\mathrm{ph}}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}\right) \subset S_{\infty}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right) \tag{2.171}
\end{equation*}
$$

The definition of the subspace ( (L.1.P.1 1 ) is straightforward. First we note that if $a \in \mathcal{C}^{\infty}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$ is homogeneous of degree $m \in \mathbb{R}$ in $|\xi| \geq 1$, then
eq:P.2 (2.172)

$$
a(z, t \xi)=t^{m} a(z, \xi),|t|,|\xi| \geq 1
$$

If it also satisfies the uniform regularity estimates

$$
\begin{equation*}
\sup _{z \in \mathbb{R}^{n},|\xi| \leq 2}\left|D_{z}^{\alpha} D_{\xi}^{\beta} a(z, \xi)\right|<\infty \forall \alpha, \beta \tag{2.173}
\end{equation*}
$$

then in fact

$$
\begin{equation*}
a \in S_{\infty}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right) \tag{2.174}
\end{equation*}
$$

Indeed, ( (Leq:P. 3.173 ) is exactly the restriction of the symbol estimates to $z \in \mathbb{R}^{p},|\xi| \leq 2$. On the other hand, in $|\xi| \geq 1, a(z, \xi)$ is homogeneous so

$$
\left|D_{z}^{\alpha} D_{\xi}^{\beta} a(z, \xi)\right|=|\xi|^{m-|\beta|}\left|D_{z}^{\alpha} D_{\xi}^{\beta} a(z, \hat{\xi})\right|, \quad \hat{\xi}=\frac{\xi}{|\xi|}
$$

from which the symbol estimates follow.
DEFINITION 2.2. For any $m_{\operatorname{l}} \in \mathbb{R}$, the subspace of (one-step) ${ }^{5}$ polyhomogeneous symbols is defined as a subset ( 2.17 I$)$ by the requirement that $a \in S_{\mathrm{ph}}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$ if and only if there exist elements $a_{m-j}(z, \xi) \in S_{\infty}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$ which are homogeneous of degree $m-j$ in $|\xi| \geq 1$, for $j \in \mathbb{N}_{0}$, such that

$$
\begin{equation*}
a \sim \sum_{j} a_{m-j} \tag{2.175}
\end{equation*}
$$

[^7]Clearly

$$
\begin{equation*}
S_{\mathrm{ph}}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right) \cdot S_{\mathrm{ph}}^{m^{\prime}}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right) \subset S_{\mathrm{ph}}^{m+m^{\prime}}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right) \tag{2.176}
\end{equation*}
$$

since the asymptotic expansion of the product is given by the formal product of the asymtotic expansion. In fact there is equality here, because

$$
\begin{equation*}
\left(1+|\xi|^{2}\right)^{m / 2} \in S_{\mathrm{ph}}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right) \tag{2.177}
\end{equation*}
$$

and multiplication by $\left(1+|\xi|^{2}\right)^{m / 2}$ is an isomorphism of the space $S_{\mathrm{ph}}^{0}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$ onto $S_{\mathrm{ph}}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$. Furthermore differentiation with respect to $z_{j}$ or $\xi_{l}$ preserves asymptotic homogeneity so

$$
\begin{aligned}
& D_{x_{j}}: S_{\mathrm{ph}}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right) \longrightarrow S_{\mathrm{ph}}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right) \\
& D_{\xi_{l}}: S_{\mathrm{ph}}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right) \longrightarrow S_{\mathrm{ph}}^{m-1}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right) \quad \forall j=1, \ldots, n .
\end{aligned}
$$

It is therefore no surprise that the polyhomogeneous operators form a subalgebra.
Proposition 2.9. The spaces $\Psi_{\mathrm{ph}}^{m}\left(\mathbb{R}^{n}\right) \subset \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ defined by the condition that the kernel of $A \in \Psi_{\mathrm{ph}}^{m}\left(\mathbb{R}^{n}\right)$ should be of the form $I(a)$ for some

$$
\begin{equation*}
a \in\left(1+|x-y|^{2}\right)^{w / 2} S_{\mathrm{ph}}^{m}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right) \tag{2.178}
\end{equation*}
$$

form an order-filtered $*$-algebra.
Proof. Since the definition shows that

$$
\Psi_{\mathrm{ph}}^{m}\left(\mathbb{R}^{n}\right) \subset \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)
$$

we know already that

$$
\Psi_{\mathrm{ph}}^{m}\left(\mathbb{R}^{n}\right) \cdot \Psi_{\mathrm{ph}}^{m^{\prime}}\left(\mathbb{R}^{n}\right) \subset \Psi_{\infty}^{m+m^{\prime}}\left(\mathbb{R}^{n}\right)
$$

To see that products are polyhomogeneous it suffices to use ( (kq:P. $\left.\frac{6}{2.16}\right)$ and (kq:P.8) which together show that the asymptotic formulæ describing the left symbols of $A \in \Psi_{\mathrm{ph}}^{m}\left(\mathbb{R}^{n}\right)$ and $B \in \Psi_{\mathrm{ph}}^{m^{\prime}}\left(\mathbb{R}^{m}\right)$, e.g.

$$
\left.\sigma_{L}(A) \sim \sum_{\alpha} \frac{i^{|\alpha|}}{\alpha!} D_{\xi}^{\alpha} D_{y}^{\alpha} a(x, y, \xi)\right|_{y=x}
$$

imply that $\sigma_{L}(A) \in S_{\mathrm{ph}}^{m}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right), \sigma_{L}(B) \in S_{\mathrm{ph}}^{m}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right)$. Then the asymptotic formula for the product shows that $\sigma_{L}(A \cdot B) \in S_{\mathrm{ph}}^{m+m^{\prime}}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right)$.

The proof of $*$-invariance is similarly elementary, since if $A=I(a)$ then $A^{*}=$ $I(b)$ with $b(x, y, z)=\overline{a(y, x, \xi)} \in S_{\mathrm{ph}}^{m}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right)$.

This subalgebra is usually denoted simply $\Psi^{m}\left(\mathbb{R}^{n}\right)$ and its elements are often said to be 'classical' pseudodifferential operators. As a small exercise in the use of the principal symbol map we shall show that
eq:P. 10

$$
\begin{gather*}
A \in \Psi_{\mathrm{ph}}^{m}\left(\mathbb{R}^{n}\right), A \text { elliptic } \Longrightarrow \exists \text { a parametrix } \\
B \in \Psi_{\mathrm{ph}}^{-m}\left(\mathbb{R}^{n}\right), A \cdot B-\mathrm{Id}, B \cdot A-\mathrm{Id} \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right) \tag{2.179}
\end{gather*}
$$

In fact we already know that $B \in \Psi_{\infty}^{-m}\left(\mathbb{R}^{n}\right)$ exists with these properties, and even that it is unique modulo $\Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$. To show that $B \in \Psi_{\mathrm{ph}}^{-m}\left(\mathbb{R}^{n}\right)$ we can use the principal symbol map.

For elements $A \in \Psi_{\mathrm{ph}}^{m}\left(\mathbb{R}^{n}\right)$ the principal symbol $\sigma_{m}(A) \in S_{\infty}^{m-[1]}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right)$ has a preferred class of representatives, namely the leading term in the expansion of $\sigma_{L}(A)$

$$
\sigma_{m}(A)=\sigma(\xi) a_{m}(x, \xi) \quad \bmod S_{\mathrm{ph}}^{m-1}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right)
$$

where $\sigma|\xi|=1$ in $|\xi| \geq 1, \sigma|\xi|=0$ in $|\xi| \leq 1 / 2$. It is even natural to identify the principal symbol with $a_{m}(x, \xi)$ as a homogeneous function. Then we can see that
(2.180) $A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right), \sigma_{m}(A)$ homogeneous of degree m

$$
\Longleftrightarrow \Psi_{\mathrm{ph}}^{m}\left(\mathbb{R}^{n}\right)+\Psi_{\infty}^{m-1}\left(\mathbb{R}^{n}\right)
$$

Indeed, we just subtract from $A$ an element $A_{1} \in \Psi_{\mathrm{ph}}^{m}\left(\mathbb{R}^{n}\right)$ with $\sigma_{m}\left(A_{1}\right)=\sigma_{m}(A)$, then $\sigma_{m}\left(A-A_{1}\right)=0$ so $A-A_{1} \in \Psi^{m-1}\left(\mathbb{R}^{n}\right)$.

So, returning to the proof of $\left(\frac{\operatorname{ea.P} 10}{2.179)}\right.$ note straight away that

$$
\sigma_{-m}(B)=\sigma_{m}(A)^{-1}
$$

has a homogeneous representative, namely $a_{m}(x, \xi)^{-1}$. Thus we have shown that for $j=1$
eq:P.12 (2.181)

$$
B \in \Psi_{\mathrm{ph}}^{-m}\left(\mathbb{R}^{n}\right)+\Psi_{\infty}^{-m-j}\left(\mathbb{R}^{n}\right)
$$

We take ( (eq:P.181) as an inductive hypthesis for general $j$. Writing this decomposition $B=B^{\prime}+B_{j}$ it follows from the identity $\left(\frac{\mathrm{eq}: \mathrm{P} .10}{2.19)}\right.$ that

$$
A \cdot B=A \cdot B^{\prime}+A B_{j}=\mathrm{Id} \quad \bmod \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)
$$

so

$$
A \cdot B_{j}=\operatorname{Id}-A B^{\prime} \in \Psi_{\mathrm{ph}}^{0}\left(\mathbb{R}^{n}\right) \cap \Psi_{\infty}^{-j}\left(\mathbb{R}^{n}\right)=\Psi_{\mathrm{ph}}^{-j}\left(\mathbb{R}^{n}\right)
$$

Now applying $B$ on the left, or using the principal symbol map, it follows that $B_{j} \in \Psi_{\mathrm{ph}}^{-m-j}\left(\mathbb{R}^{n}\right)+\Psi_{\infty}^{-m-j-1}\left(\mathbb{R}^{n}\right)$ which gives the inductive hypothesis ( (2.181) for $j+1$.

It is usually the case that a construction in $\Psi_{\infty}^{*}\left(\mathbb{R}^{n}\right)$, applied to an element of $\Psi_{\mathrm{ph}}^{*}\left(\mathbb{R}^{n}\right)$ will yield an element of $\Psi_{\mathrm{ph}}^{*}\left(\mathbb{R}^{n}\right)$ and when this is the case it can generally be confirmed by an inductive argument like that used above to check ( $\overline{2} .17 \overline{9}$ ).

As a subspace ${ }^{6}$

$$
S_{\mathrm{ph}}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right) \subset S_{\infty}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)
$$

is not closed. Indeed, since it contains $S_{\infty}^{-\infty}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$, its closure contains all of $S_{\infty}^{m^{\prime}}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$ for $m^{\prime}<m$. In fact it is a dense subspace. ${ }^{7}$ To capture its properties we can strengthen the topology $S_{\mathrm{ph}}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$ inherits from $S_{\infty}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$.

Then, as well as the symb:t norms $\|\cdot\|_{N, m}$ in $\left(\frac{3.8}{2.7}\right)$ we can add norms on the terms in the expansions in ( $\left(\frac{2.17}{2.1}\right)$

$$
\begin{equation*}
\left\|D_{x}^{\alpha} D_{\xi}^{\beta} a_{m-j}(x, \xi)\right\|_{L^{\infty}(G)}, G=\mathbb{R}^{p} \times\{1 \leq|\xi| \leq 2\} \tag{2.182}
\end{equation*}
$$

Then we can further add the symbol norms ensuring ( (2.17.P.5), i.e.,
eq:P. 15

$$
\begin{equation*}
\left\|a-\sum_{j=0}^{k} a_{m-j}\right\|_{m-k-1, N} \quad \forall k, N \tag{2.183}
\end{equation*}
$$

[^8]Together these give a countable number of norms on $S_{\mathrm{ph}}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$. With respect to the metric topology defined as in $\left(\frac{3.9}{2.8}\right)$ the spaces $S_{\mathrm{ph}}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$ are then complete. ${ }^{8}$. Later when we wish to topologize $\Psi_{\mathrm{ph}}^{m}\left(\mathbb{R}^{n}\right)$, or rather related algebras, it is this type of topology we will use. Namely we identify

$$
\begin{equation*}
\left(1+|\xi|^{2}\right)^{\frac{1}{2} m} \in S_{\infty}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right) \tag{2.186}
\end{equation*}
$$

for any $p$. Use this to show that

$$
S_{\infty}^{m}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right) \cdot S_{\infty}^{m^{\prime}}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)=S_{\infty}^{m+m^{\prime}}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)
$$

Problem 2.2. Consider $w=0$ and $n=2$ in the definition of symbols and show that if $a \in S_{\infty}^{1}\left(\mathbb{R}^{2}\right)$ is elliptic then for $r>0$ sufficiently large the integral

$$
\int_{0}^{2 \pi} \frac{1}{2 \pi} \frac{1}{a\left(r e^{i \theta}\right)} \frac{d}{d \theta} a\left(r e^{i \theta}\right) d \theta=\frac{1}{2 \pi} \int_{0}^{2 \pi} \frac{d}{d \theta} \log a\left(r e^{i \theta}\right) d \theta
$$

exists and is an integer independent of $r$, where $z=\xi_{1}+i \xi_{2}$ is the complex variable in $\mathbb{R}^{2}=\mathbb{C}$. Conclude that there is an elliptic symbol, $a$ on $\mathbb{R}^{2}$, such that there does not exist $b$, a symbol with

$$
\begin{equation*}
b \neq 0 \text { on } \mathbb{R}^{2} \text { and } a(\xi)=b(\xi) \text { for }|\xi|>r \tag{2.187}
\end{equation*}
$$

for any $r$.

[^9]\[

$$
\begin{equation*}
|a(z, \xi)| \leq C(1+|\xi|)^{m^{\prime}}, m^{\prime}<m \tag{2.188}
\end{equation*}
$$

\]

is necessarily in the space $S_{\infty}^{m^{\prime}+\epsilon}\left(\mathbb{R}_{z}^{p} ; \mathbb{R}_{\xi}^{n}\right)$ for all $\epsilon>0$.
22.2 .1998 .142
22.2 .1998 .143

$$
\begin{equation*}
c_{\phi}(z, \xi)=\phi\left(z, \frac{\xi}{|\xi|}\right)(1-\psi)(\xi) \in S^{0}\left(\mathbb{R}_{z}^{p} ; \mathbb{R}_{\xi}^{n}\right) \tag{2.189}
\end{equation*}
$$

If $a \in S_{\infty}^{m}\left(\mathbb{R}_{z}^{p} ; \mathbb{R}_{\xi}^{n}\right)$ define the cone support of $a$ in terms of its complement

$$
\begin{align*}
& \text { cone } \operatorname{supp}(a)^{\complement}=\left\{(\bar{z}, \bar{\xi}) \in \mathbb{R}_{z}^{p} \times\left(\mathbb{R}_{\xi}^{n} \backslash\{0\}\right) ; \exists\right.  \tag{2.190}\\
& \left.\qquad \phi \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}_{z}^{p} ; \mathbb{R}^{n}\right), \phi(\bar{z}, \bar{\xi}) \neq 0, \text { such that } c_{\phi} a \in S_{\infty}^{-\infty}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)\right\}
\end{align*}
$$

Show that if $a \in S_{\infty}^{m}\left(\mathbb{R}_{z}^{p} ; \mathbb{R}_{\xi}^{n}\right)$ and $b \in S_{\infty}^{m^{\prime}}\left(\mathbb{R}_{z}^{p} ; \mathbb{R}_{\xi}^{n}\right)$ then

$$
27.1 .2003 .30
$$

$$
\begin{equation*}
\text { cone } \operatorname{supp}(a b) \subset \text { cone } \operatorname{supp}(a) \cap \text { cone } \operatorname{supp}(b) . \tag{2.191}
\end{equation*}
$$

If $a \in S_{\infty}^{m}\left(\mathbb{R}_{z}^{p} ; \mathbb{R}_{\xi}^{n}\right)$ and cone $\operatorname{supp}(a) \emptyset$ does it follow that $a \in S_{\infty}^{-\infty}\left(\mathbb{R}_{z}^{p} ; \mathbb{R}_{\xi}^{n}\right)$ ?
Problem 2.5. Prove that $\left(\frac{3.32}{2.30}\right)$ is a characterization of functions $a \in(1+$ $\left.|x-y|^{2}\right)^{w / 2} S^{m}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right)$. [Hint: Use Liebniz' formula to show instead that the equivalent estimates

$$
\left|D_{x}^{\alpha} D_{y}^{\beta} D_{\xi}^{\gamma} a(x, y, \xi)\right| \leq C_{\alpha, \beta, \gamma}\left(1+|x-y|^{2}\right)^{w / 2}(1+|\xi|)^{m-|\gamma|} \forall \alpha, \beta, \gamma \in \mathbb{N}_{0}^{n}
$$

characterize this space.]
Problem 2.6. Show that $A \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$ if and only if its Schwartz kernel is $\mathcal{C}^{\infty}$ and satisfies all the estimates

$$
\begin{equation*}
\left|D_{x}^{\alpha} D_{y}^{\beta} a(x, y)\right|<C_{\alpha, \beta, N}(1+|x-y|)^{-N} \tag{2.192}
\end{equation*}
$$

for multiindices $\alpha, \beta \in \mathbb{N}_{0}^{n}$ and $N \in \mathbb{N}_{0}$.
prob:NN
prob:MM
prob:DD
prob:CC
prob:FF

Problem 2.7. Polyhomogeneous symbols as smooth functions.
Problem 2.8. General polyhomogeneous symbols and operators.
Problem 2.9. Density of polyhomogeneous symbols in $L^{\infty}$ symbols of the same order.

Problem 2.10. Completeness of the spaces of polyhomogeneous symbols.
Problem 2.11. Fourier transform??
Problem 2.12. Show that the kernel of any element of $\Psi_{\infty}^{\infty}\left(\mathbb{R}^{n}\right)$ is $\mathcal{C}^{\infty}$ away from the diagonal. Hint: Prove that $(x-y)^{\alpha} K(x, y)$ becomes increasingly smooth as $|\alpha|$ increases.
21.2 .1998 .117

Problem 2.13. Show that for any $m \geq 0$ the unit ball in $H^{m}\left(\mathbb{R}^{n}\right) \subset L^{2}\left(\mathbb{R}^{n}\right)$ is not precompact, i.e. there is a sequence $f_{j} \in H^{m}\left(\mathbb{R}^{n}\right)$ which has $\left\|f_{j}\right\|_{m} \leq 1$ and has no subsequence convergent in $L^{2}\left(\mathbb{R}^{n}\right)$.

Problem 2.14. Show that for any $R>0$ there exists $N>0$ such that the Hilbert subspace of $H^{N}\left(\mathbb{R}^{n}\right)$

$$
\begin{equation*}
\left\{u \in H^{N}\left(\mathbb{R}^{n}\right) ; u(x)=0 \text { in }|x|>R\right\} \tag{2.193}
\end{equation*}
$$

is compactly included in $H_{2}^{2}\left(\mathbb{R}^{n} n\right.$, i.e. the intersection of the unit ball in $H^{N}\left(\mathbb{R}^{n}\right)$ with the subspace $\left(\frac{2.193) \text { is precompact in } L^{2}\left(\mathbb{R}^{n}\right) \text {. Hint: This is true for any }}{2}\right.$ $N>0$, taking $N \gg 0$ will allow you to use the Sobolev embedding theorem and Arzela-Ascoli.

Problem 2.15. Using Problem $\frac{\text { 21.2.1998.118 }}{2.14 \text { (or otherwise) show that for any } \epsilon>0}$

$$
(1+|x|)^{\epsilon} H^{\epsilon}\left(\mathbb{R}^{n}\right) \hookrightarrow L^{2}\left(\mathbb{R}^{n}\right)
$$

is a compact inclusion, i.e. any infinite sequence $f_{n}$ such that $\left(1+|x|^{2}\right)^{-\epsilon}$ is bounded in $H^{\epsilon}\left(\mathbb{R}^{n}\right)$ has a subsequence convergent in $L^{2}\left(\mathbb{R}^{n}\right)$. Hint: Choose $\phi \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)$ with $\phi(x)=1$ in $|x|<1$ and, for each $k$, consider the sequence $\phi(x / k) f_{j}$. Show that the Fourier transform converts this into a sequence which is bounded in $(1+$ $\left.|\xi|^{2}\right)^{-\frac{1}{2} \epsilon} H^{N}\left(\mathbb{R}_{\xi}^{n}\right)$ for any $N$. Deduce that it has a convergent subsequence in $L^{2}\left(\mathbb{R}^{n}\right)$. By diagonalization (and using the rest of the assumption) show that $f_{j}$ itself has a convergent subsequence.

Problem 2.16. About $\rho$ and $\delta$.
Problem 2.17. Prove the formula ( $\frac{21.2 .1998 .104}{2.185)}$ for the left-reduced symbol of the operator $A_{T}$ obtained from the pseudodifferential operator $A$ by linear change of variables. How does the right-reduced symbol transform?

Problem 2.18. Density of $\mathcal{S}\left(\mathbb{R}^{n}\right)$ in $L^{2}\left(\mathbb{R}^{n}\right)$.
Problem 2.19. Square-root of a positive elliptic symbol is a symbol.
Problem 2.20. Write out a proof to Proposition $\frac{21.2 .1998 .107}{3.2 .}$. Hint (Just to do it el ${ }_{\text {el }}$ - 1998.112 egantly, it is straightforward enough): Write $A$ in right-reduced form as in ( $\frac{21.2}{2.74}$ ) and apply it to $\hat{u}$; this gives a formula for $\hat{A} u$.

Problem 2.21. Show that any continuous linear operator

$$
\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}\left(\mathbb{R}^{n}\right)
$$

has Schwartz kernel in $\mathcal{S}\left(\mathbb{R}^{2 n}\right)$.

## CHAPTER 3

## Isotropic and scattering calculi

In this chapter many of the general constructions with pseudodifferential operators are carried out in the context of two global calculi on $\mathbb{R}^{n}$. Partly this is done for the obvious reason, that these calculi and results have interesting applications, and partly it is preparatory to the discussion of the geometric algebras of pseudodifferential operators on a compact manifold without boundary and for the scattering algebra on a compact manifold with boundary. Thus, while this chapter is somewhat interstitial, it is designed to clarify the later discussions by separating the generalities of the construction from the particulars of the calculus involved. It should be noted that in this chapter it is generally the polyhomogeneous calculus which is under discussion unless it is explicitly stated to the contrary.

### 3.1. Isotropic operators

As noted in the discussion in Chapter $\frac{C D}{C, E u c l i d e a n}$, there are other sensible choices of the class of amplitudes which can be admitted in the definition of a space of pseudodifferential operators than the basic case of $S_{\infty}^{m}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right)$ discussed there. One of the smallest such choices is the class which is completely symmetric in the variables $x$ and $\xi$ and consists of the symbols on $\mathbb{R}^{2 n}$. Thus, $a \in S_{\infty}^{m}\left(\mathbb{R}_{x, \xi}^{2 n}\right)$ satisfies the estimates

$$
\begin{equation*}
\left|D_{x}^{\alpha} D_{\xi}^{\beta} a(x, \xi)\right| \leq C_{\alpha, \beta}(1+|x|+|\xi|)^{m-|\alpha|-|\beta|} \tag{3.1}
\end{equation*}
$$

for all multiindices $\alpha$ and $\beta$. If $m \leq 0$ this is in the space $S_{\infty}^{m}\left(\mathbb{R}_{x}^{n} ; \mathbb{R}_{\xi}^{n}\right)$; if $m>0$ it is not, however,
1.2 .2000 .303
8.2.1998.98

Lemma 3.1. For any $p$ and $n$

$$
S_{\infty}^{m}\left(\mathbb{R}^{p+n}\right) \subset \begin{cases}\bigcap_{0 \geq r \geq m}\left(1+|x|^{2}\right)^{r} S^{m-r}\left(\mathbb{R}_{x}^{p} ; \mathbb{R}_{\xi}^{n}\right) & m \leq 0  \tag{3.2}\\ \left(1+|x|^{2}\right)^{m / 2} S^{m}\left(\mathbb{R}_{x}^{p} ; \mathbb{R}_{\xi}^{n}\right), & m>0\end{cases}
$$

Proof. This follows from ( $\frac{1.2 .2000 .304}{3.1) \text { and the inequalities }}$

$$
\begin{gathered}
1+|x|+|\xi| \leq(1+|x|)(1+|\xi|) \\
1+|x|+|\xi| \geq(1+|x|)^{t}(1+|\xi|)^{1-t}, 0 \leq t \leq 1
\end{gathered}
$$

In view of these estimates the following definition makes sense.
1.2 .2000 .313
1.2.2000.314

Definition 3.1. For any $m \in \mathbb{R}$ we define

$$
\begin{equation*}
\Psi_{\mathrm{iso}}^{m}\left(\mathbb{R}^{n}\right) \subset \Psi_{\infty-\mathrm{iso}}^{m}\left(\mathbb{R}^{n}\right) \subset\langle x\rangle^{m_{+}} \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right) \tag{3.3}
\end{equation*}
$$

as the subspaces determined by
1.2 .2000 .315

$$
\begin{gather*}
A \in \Psi_{\mathrm{iso}}^{m}\left(\mathbb{R}^{n}\right) \Longleftrightarrow \sigma_{L}(A) \in S_{\mathrm{ph}}^{m}\left(\mathbb{R}^{2 n}\right) \\
A \in \Psi_{\infty-\text { iso }}^{m}\left(\mathbb{R}^{n}\right) \Longleftrightarrow \sigma_{L}(A) \in S_{\infty}^{m}\left(\mathbb{R}^{2 n}\right) \tag{3.4}
\end{gather*}
$$

As in the discussion of the traditional algebra in Chapter $\frac{\mathbb{C} \cdot \text { Euclidean }}{2}$ we show the $*-$ invariance and composition properties of these spaces of operators by proving an appropriate 'reduction' theorem. Note however that there is a small difficulty here. Namely it might be supposed that it is enough to analyse $I(a)$ for $a \in S_{\infty}^{m}\left(\mathbb{R}^{3 n}\right)$. This however is not the case. Indeed the definition above is in terms of left-reduced symbols. If $a \in S_{\infty}^{m}\left(\mathbb{R}^{2 n}\right)$ is regarded as a function on $\mathbb{R}^{3 n}$ which is independent of one of the variables then it is in general not an element of $S_{\infty}^{m}\left(\mathbb{R}^{3 n}\right)$. For this reason we consider some more 'hybrid' estimates.

Consider a subdivision of $\mathbb{R}^{3 n}$ into two closed regions:
1.2 .2000 .316

$$
\begin{align*}
R_{1}(\epsilon) & =\left\{(x, y, \xi) \in \mathbb{R}^{3 n} ;|x-y| \leq \epsilon\left(1+|x|^{2}+|y|^{2}+|\xi|^{2}\right)^{\frac{1}{2}}\right\} \\
R_{2}(\epsilon) & =\left\{(x, y, \xi) \in \mathbb{R}^{3 n} ;|x-y| \geq \epsilon\left(1+|x|^{2}+|y|^{2}+|\xi|^{2}\right)^{\frac{1}{2}}\right\} \tag{3.5}
\end{align*}
$$

If $a \in \mathcal{C}^{\infty}\left(\mathbb{R}^{3 n}\right)$ consider the estimates

$$
1.2 .2000 .317
$$

$$
\left|D_{x}^{\alpha} D_{y}^{\beta} D_{\xi}^{\gamma} a(x, y, \xi)\right| \leq C_{\alpha, \beta, \gamma} \begin{cases}\langle(x, y, \xi)\rangle^{m-|\alpha|-|\beta|-|\gamma|} & \text { in } R_{1}\left(\frac{1}{8}\right)  \tag{3.6}\\ \langle(x, y)\rangle^{m+}\langle\xi\rangle^{m-|\gamma|} & \text { in } R_{2}\left(\frac{1}{8}\right)\end{cases}
$$

The choice $\epsilon=\frac{1}{8}$ here is rather arbitrary. However if $\epsilon$ is decreased, but kept positive the same estimates continue to hold for the new subdivision, since the estimates in $R_{1}$ are stronger than those in $R_{2}$ (which is increasing at the expense of $R_{1}$ as $\epsilon$ decreases). Notice too that these estimates do in fact imply that $a \in$ $\langle(x, y)\rangle^{m+} S_{\infty}^{m}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right)$.
1.2 .2000 .318

Proposition 3.1. If $a_{4.19} \in \mathcal{C}^{\infty}\left(\mathbb{R}^{3 n}\right)$ satisfies the estimates $\left(\frac{1}{3.2 .2000 .317}\right.$ then $A=$ $I(a) \in \Psi_{\infty-\text { iso }}^{m}\left(\mathbb{R}^{n}\right)$ and (L2.58) holds for $\sigma_{L}(A)$.

Proof. We separate $a$ into two pieces. Choose $\chi \in \mathcal{C}_{c}^{\infty}(\mathbb{R})$ with $0 \leq \chi \leq 1$, with support in $\left[-\frac{1}{8}, \frac{1}{8}\right]$ and with $\chi \equiv 1$ on $\left[-\frac{1}{9}, \frac{1}{9}\right]$. Then consider the cutoff function on $\mathbb{R}^{3 n}$

$$
\begin{equation*}
\psi(x, y, \xi)=\chi\left(\frac{|x-y|^{2}}{1+|x|^{2}+|y|^{2}+|\xi|^{2}}\right) \tag{3.7}
\end{equation*}
$$

Clearly, $\psi$ has support in $R_{1}\left(\frac{1}{8}\right)$ and $\psi \in S_{\infty}^{0}\left(\mathbb{R}^{3 n}\right)$. It follows then that $a^{\prime}=\psi a \in$ $S_{\mathrm{iso}}^{m}\left(\mathbb{R}^{3 n}\right)$. On the other hand, $a^{\prime \prime}=(1-\psi) a$ has support in $R_{2}\left(\frac{1}{9}\right)$. In this region $|x-y|,\langle(x, y)\rangle$ and $\langle(x, y, \xi)\rangle$ are bounded by constant multiples of each other. Thus $a^{\prime \prime}$ satisfies the estimates
1.2 .2000 .320

$$
\begin{aligned}
\left|D_{x}^{\alpha} D_{y}^{\beta} D_{\xi}^{\gamma} a^{\prime \prime}(x, y, \xi)\right| \leq & C_{\alpha, \beta, \gamma}|x-y|^{m_{+}}\langle\xi\rangle^{m-|\gamma|} \\
& \leq C_{\alpha, \beta, \gamma}^{\prime}\langle(x, y, \xi)\rangle^{m_{+}}\langle\xi\rangle^{m-|\gamma|}, \operatorname{supp}\left(a^{\prime \prime}\right) \subset R_{2}\left(\frac{1}{9}\right)
\end{aligned}
$$

First we check that $I\left(a^{\prime \prime}\right) \in \mathcal{S}\left(\mathbb{R}^{2 n}\right)$. On $R_{2}\left(\frac{1}{9}\right)$ it is certainly the case that $|x-y| \geq \frac{1}{9}\langle(x, y)\rangle$ and by integration by parts

$$
|x-y|^{2 p} D_{x}^{\alpha} D_{y}^{\beta} I\left(a^{\prime \prime}\right)=I\left(\left|D_{\xi}\right|^{2 p} D_{x}^{\alpha} D_{y}^{\beta} a^{\prime \prime}\right)
$$

For all sufficiently large $p$ it follows from ( $\left(\frac{1}{3} .2 .2000 .320\right.$ that this is the product of $\langle(x, y)\rangle^{m_{+}}$ and a bounded continuous function. Thus, $I\left(a^{\prime \prime}\right) \in \mathcal{S}\left(\mathbb{R}^{2 n}\right)$ is the kernel of an operator in $\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$.

So it remains only to show that $A^{\prime}=I\left(a^{\prime}\right) \in \Psi_{\infty-\text { iso }}^{m}\left(\mathbb{R}^{n}\right)$. Certainly this is an element of $\langle x\rangle^{m+} \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$. The left-reduced symbol of $A^{\prime}$ has an asymptotic expansion, as $\xi \rightarrow \infty$, given by the usual formula, namely ( $\overline{2} .58$ ). Each of the terms in this expansion

$$
a_{L}\left(A^{\prime}\right) \sim \sum_{\alpha} \frac{i^{|\alpha|}}{\alpha!} D_{y}^{\alpha} D_{\xi}^{\alpha} a(x, x, \xi)
$$

is in the space $S_{\infty}^{m-2|\alpha|}\left(\mathbb{R}^{2 n}\right)$. Thus we can actually choose an asymptotic sum in the stronger sense that

$$
b^{\prime} \in S_{\infty}^{m}\left(\mathbb{R}^{2 n}\right), b_{N}=b^{\prime}-\sum_{|\alpha|<N} \frac{i^{|\alpha|}}{\alpha!} D_{x}^{\alpha} D_{\xi}^{\alpha} a(x, \xi) \in S_{\infty}^{m-2 N}\left(\mathbb{R}^{2 n}\right) \forall N .
$$

Consider the remainder term in ( $\left.\frac{4.8}{2.47}\right)$, given by $\left(\frac{4.6}{2.44}\right)$ and $\left(\frac{4.7}{2.45}\right)$. Integrating by parts in $\xi$ to remove the factors of $(x-y)^{\alpha}$ the remainder, $R_{N}$, can be written as a pseudodifferential operator with amplitude

$$
r_{N}(x, y, \xi)=\sum_{|\alpha|=N} \frac{i^{|\alpha|}}{\alpha!} \int_{0}^{1} d t(1-t)^{N}\left(D_{\xi}^{\alpha} D_{y}^{\alpha} a\right)((1-t) x+t y, \xi)
$$

This satisfies the estimates $\left(\frac{1}{3} \cdot 2 \cdot 2\right)$ with $m$ replaced by $m-2 N$. Indeed from the symbol estimates on $a^{\prime}$ the integrand satisfies the bounds

$$
\begin{aligned}
& \left|D_{x}^{\beta} D_{y}^{\gamma} D_{\xi}^{\delta} D_{\xi}^{\alpha} D_{y}^{\alpha} a^{\prime}((1-t) x+t y, \xi)\right| \\
& \quad \leq C\left(1+\mid\left(x+t(x-y)|+|\xi|)^{m-2 N-|\beta|-|\gamma|-|\delta|}\right.\right.
\end{aligned}
$$

In $R_{1}\left(\frac{1}{8}\right),|x-y| \leq \frac{1}{8}\langle(x, y, \xi)\rangle$ so $|x+t(x-y)|+|\xi| \geq \frac{1}{2}\langle(x, y, \xi)\rangle$ and these estimates imply the full. Symbol estimates there. On $R_{2}$ we immediately get the weaker estimates in (3.7).

Thus, for large $N$, the remainder term gives an operator in $\langle x\rangle^{\frac{m}{2}-N} \Psi_{\text {iso }}^{\frac{m}{2}-N}\left(\mathbb{R}^{n}\right)$. The difference between $A^{\prime}$ and the operator $B^{\prime} \in \Psi_{\infty-\text { iso }}^{m}\left(\mathbb{R}^{n}\right)$, which is $R_{N}$ plus an operator in $\Psi_{\infty-\text { iso }}^{m-2 N}\left(\mathbb{R}^{n}\right)$ for any $N$ is therefore in $\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$. Thus $A \in \Psi_{\infty-\text { iso }}^{m}\left(\mathbb{R}^{n}\right)$.

This is a perfectly adequate replacement in this context for our previous reduction theorem, so now we can show the basic result.

Theorem 3.1. The spaces $\Psi_{\infty-\mathrm{iso}}^{m}\left(\mathbb{R}^{n}\right)$ (resp. $\Psi_{\mathrm{iso}}^{m}\left(\mathbb{R}^{n}\right)$ ) of isotropic (resp: 2000.315 polyhomogeneous isotropic) pseudodifferential operators on $\mathbb{R}^{n}$, defined by ( $\mathbf{3 . 4}$ ) form an order-filtered $*$-algebra with residual space $\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)=\mathcal{S}\left(\mathbb{R}^{2 n}\right)$ (resp. the same) as spaces of kernels.

Proof. The condition that a continuous linear operator $A$ on $\mathcal{S}\left(\mathbb{R}^{n}\right)$ be an element of $\Psi_{\infty-\text { iso }}^{m}\left(\mathbb{R}^{n}\right)$ is that it be an element of $\left(1+|x|^{2}\right)^{m / 2} \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ if $m \geq 0$ or $\Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ if $m<0$ with left-reduced symbol an element of $S_{\infty}^{m}\left(\mathbb{R}_{x, \xi}^{2 n}\right)$ :

$$
\begin{equation*}
q_{l}: S_{\infty}^{m}\left(\mathbb{R}^{2 n}\right) \longleftrightarrow \Psi_{\infty-\mathrm{iso}}^{m}\left(\mathbb{R}^{n}\right) \tag{3.9}
\end{equation*}
$$

Thus $A^{*}$ has right-reduced symbol in $S_{\infty}^{m}\left(\mathbb{R}_{1.2 n}^{2 n}\right)$.2J0i.3 satisfies the estimates ( $\frac{1}{1} .2 .2000 .317$ a function of $x, y$ and $\xi$. Thus Propostion 3.1 shows that $A^{*} \in \Psi_{\infty-\text { iso }}^{m}\left(\mathbb{R}^{n}\right)$, since
its left-reduced symbol is in $S_{\infty}^{m}\left(\mathbb{R}^{2 n}\right)$. This shows the $*$-invariance. Moreover it also follows that any $B \in \Psi_{\infty-\text { iso }}^{m^{\prime}}\left(\mathbb{R}^{n}\right)$ has right-reduced symbol in $S_{\infty}^{m^{\prime}}\left(\mathbb{R}^{2 n}\right)$. Thus if $A \in \Psi_{\infty-\text { iso }}^{m}\left(\mathbb{R}^{n}\right)$ and $B \in \Psi_{\text {iso }}^{m^{\prime}}\left(\mathbb{R}^{n}\right)$ then using this result to right-reduce $B$ we see that the composite opertor has kernel $I\left(a_{L}(x, \xi) b_{R}(y, \xi)\right)$ where $a_{L} \in S_{\infty}^{m}\left(\mathbb{R}^{2 n}\right)$ and $b_{102} \in S_{31 \infty}^{m^{\prime}}\left(\mathbb{R}^{2 n}\right)$. Now it again follows that this product satisifes. the estimates (3.6) of order $m+m^{\prime}$. Hence, again applying Proposition 3.1, we conclude that $A \circ B \in \Psi_{\infty-\text { iso }}^{m+m^{\prime}}\left(\mathbb{R}^{n}\right)$. This proves the theorem for $\Psi_{\infty-\text { iso }}^{*}\left(\mathbb{R}^{n}\right)$.

The proof for the polyhomogeneous space $\Psi_{\text {iso }}^{m}\left(\mathbb{R}^{n}\right)$ follows immediately, since the symbol expansions all preserve polyhomogeneity.

One further property of the isotropic calculus that distinguishes it strongly from the traditional calculus is that it is invariant under Fourier transformation.
21.2 .1998 .107
S.Scattering.operators

Proposition 3.2. If $A \in \Psi_{\infty-\text { iso }}^{m}\left(\mathbb{R}^{n}\right)$ (resp. $\Psi_{\text {iso }}^{m}\left(\mathbb{R}^{n}\right)$ ) then $\hat{A} \in \Psi_{\infty-\text { iso }}^{m}\left(\mathbb{R}^{n}\right)$ (resp. $\Psi_{\text {iso }}^{m}\left(\mathbb{R}^{n}\right)$ ) where $\widehat{\hat{A} u}=A \hat{u}$ with $\hat{u}$ being the Fourier transform of $u \in \mathcal{S}\left(\mathbb{R}^{n}\right)$. 21.2.1998.108

The proof of this is outlined in Problem $\frac{21.2 .1}{2.20}$

### 3.2. Scattering operators

There is another calculus of pseudodifferential operators which is 'smaller' than the traditional calculus. It arises by taking amplitudes in $\left(\frac{3.2}{2.2}\right)$ which treat the base and fibre variables symmetrically, but not 'simultaneously.' Thus consider the spaces
21.2 .1998 .113

$$
\begin{align*}
& S_{\infty}^{l, m}\left(\mathbb{R}_{z}^{p}, \mathbb{R}_{\xi}^{n}\right)=\left\{a \in \mathcal{C}^{\infty}\left(\mathbb{R}^{p+n}\right)\right.  \tag{3.10}\\
&\left.\sup _{\mathbb{R}^{p+n}}(1+|z|)^{-l+|\alpha|}(1+|\xi|)^{-m+|\beta|}\left|D_{z}^{\alpha} D_{\xi}^{\beta} a(z, \xi)\right|<\infty, \forall \alpha, \beta\right\}
\end{align*}
$$

Observe that

$$
\begin{equation*}
S_{\infty}^{l, m}\left(\mathbb{R}_{z}^{p} ; \mathbb{R}_{\xi}^{n}\right) \subset\left(1+|z|^{2}\right)^{l / 2} S_{\infty}^{m}\left(\mathbb{R}_{z}^{p} ; \mathbb{R}_{\xi}^{n}\right) \tag{3.11}
\end{equation*}
$$

We can then define

$$
\begin{align*}
A \in \Psi_{\infty-\mathrm{sc}}^{l, m}\left(\mathbb{R}^{n}\right) \Longleftrightarrow A=\left(1+|x|^{2}\right)^{l / 2} B &  \tag{3.12}\\
& B \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right) \text { and } \sigma_{L}(B) \in S_{\infty}^{0, m}\left(\mathbb{R}_{x}^{n}, \mathbb{R}_{\xi}^{n}\right)
\end{align*}
$$

It follows directly from this definition and the properties of the 'traditional' operators that the left symbol map is an isomorphism

$$
1.2 .2000 .300
$$

$$
\begin{equation*}
\sigma_{L}: \Psi_{\infty-\mathrm{sc}}^{l, m}\left(\mathbb{R}^{n}\right) \longrightarrow S_{\infty}^{l, m}\left(\mathbb{R}_{x}^{n}, \mathbb{R}_{\xi}^{n}\right) \tag{3.13}
\end{equation*}
$$

To prove that this is an algebra, we need first the analogue of the asymptotic completeness, Proposition $\frac{4.3}{2.3}$, for symbols in $S_{\infty}^{*, *}\left(\mathbb{R}^{p} ; \mathbb{R}^{n}\right)$.
1.2 .2000 .292
1.2 .2000 .293

Lemma 3.2. If $a_{j} \in S_{\infty}^{l-j, m-j}\left(\mathbb{R}^{p}, \mathbb{R}^{n}\right)$ for $j \in \mathbb{N}_{0}$ then there exists

$$
\begin{equation*}
a \in S^{l, m}\left(\mathbb{R}^{p}, \mathbb{R}^{n}\right) \text { s.t. } a-\sum_{j=0}^{N} a_{j} \in S_{\infty}^{l-N, m-N}\left(\mathbb{R}^{p}, \mathbb{R}^{n}\right) \forall N \in \mathbb{N}_{0} \tag{3.14}
\end{equation*}
$$

Even though there is some potential for confusion we write $a \sim \sum_{j} a_{j}$ for a symbol $a$ satisfying (1.2.2000.

Proof. We use the same strategy as in the proof of Proposition $\frac{4.12}{2.3}$ with the major difference that there are essentially two different symbolic variables. Thus with the same notation as in $\left(\frac{1.54}{2.54}\right)$ we set

$$
\begin{equation*}
a=\sum_{j} \phi\left(\epsilon_{j} z\right) \phi\left(\epsilon_{j} \xi\right) a_{j}(z, \xi) \tag{3.15}
\end{equation*}
$$

and we proceed to check that if the $\epsilon_{j} \downarrow 0$ fast enough as. $j_{1.2} \rightarrow \infty$ then the series converges in $S_{\infty}^{l, m}\left(\mathbb{R}^{p}, \mathbb{R}^{n}\right)$ and the limit satisfies ( $(\mathbf{3} .14)$.

The first of the seminorms, for convergence, is

$$
A_{j}=\sup _{z} \sup _{\xi}(1+|z|)^{-l}(1+|\xi|)^{-m} \phi\left(\epsilon_{j} z\right) \phi\left(\epsilon_{j} \xi\right)\left|a_{j}(z, \xi)\right| .
$$

On the support of this function either $|z| \geq 1 / \epsilon_{j}$ or $|\xi| \geq 1 / \epsilon$. Thus

$$
\begin{aligned}
A_{j} \leq \sup _{z} \sup _{\xi}(1+|z|)^{-l+j}(1+ & |\xi|)^{-m+j}\left|a_{j}(z, \xi)\right| \\
\times \sup _{z} \sup _{\xi}(1+ & |z|)^{-j}(1+|\xi|)^{-j} \phi\left(\epsilon_{j} z\right) \phi\left(\epsilon_{j} \xi\right) \\
& \leq \epsilon_{j}^{j} \sup _{z} \sup _{\xi}(1+|z|)^{-l+j}(1+|\xi|)^{-m+j}\left|a_{j}(z, \xi)\right|
\end{aligned}
$$

The last term on the right is a seminorm on $S_{\infty}^{l-j, m-j}\left(\mathbb{R}^{p}, \mathbb{R}^{n}\right)$ so convergence follows by choosing the $\epsilon_{j}$ eventually smaller than a certain sequence of positive numbers. The same argument follows, as in the discussion leading to ( $\overline{2.56}$ ), for convergence of the series for the derivatives and also for the stronger convergence leading to (3.14). Since overall this is a countable collection of conditions, all can be arranged by diagonalization and the result follows.

With this result on asymptotic completeness the proof of Theorem $\frac{8.2 \cdot 1998.99}{3.1 \text { can be }}$ followed closely to yield the analogous result on products. In fact we can also define polyhomogeneous.2888. 275 rs. This requires a little work if we try to do it directly. However see ( $\overline{1.97}$ ) and Problem 1.17 which encourages us to identify

$$
\begin{gather*}
\mathrm{RC}_{p}^{*} \times \mathrm{RC}_{n}^{*}: S_{\mathrm{ph}}^{0,0}\left(\mathbb{R}^{p}, \mathbb{R}^{n}\right) \longleftrightarrow \mathcal{C}^{\infty}\left(\mathbb{S}^{p, 1} \times \mathbb{S}^{n, 1}\right)  \tag{3.16}\\
S_{\mathrm{ph}}^{l, m}\left(\mathbb{R}^{p}, \mathbb{R}^{n}\right)=\left(1+|z|^{2}\right)^{l / 2}\left(1+|\xi|^{2}\right)^{m / 2} S_{\mathrm{ph}}^{0,0}\left(\mathbb{R}^{p}, \mathbb{R}^{n}\right), l, m \in \mathbb{R}
\end{gather*}
$$

These definitions are discussed as problems starting at Problem PolyDouble 1.18 . Thus we simply define

$$
\begin{equation*}
\Psi_{\mathrm{sc}}^{l, m}\left(\mathbb{R}^{n}\right)=\left\{A \in \Psi_{\infty-\mathrm{sc}}^{l, m} ; \sigma_{L}(A) \in S_{\mathrm{ph}}^{l, m}\left(\mathbb{R}^{n}, \mathbb{R}^{n}\right)\right\} \tag{3.17}
\end{equation*}
$$

THEOREM 3.2. The spaces $\Psi_{\infty-\mathrm{sc}}^{l, m}\left(\mathbb{R}^{n}\right)$ (resp. $\Psi_{\mathrm{sc}}^{l, m}\left(\mathbb{R}^{n}\right)$ ) of scattering (resp. polyhomogeneous scattering) pseudodifferential operators on $\mathbb{R}^{n}$, form an orderbifiltered $*$-algebra

$$
\begin{equation*}
\Psi_{\infty-\mathrm{sc}}^{l, m}\left(\mathbb{R}^{n}\right) \circ \Psi_{\infty-\mathrm{sc}}^{l^{\prime}, m^{\prime}}\left(\mathbb{R}^{n}\right) \subset \Psi_{\infty-\mathrm{sc}}^{l+l^{\prime}, m+m^{\prime}}\left(\mathbb{R}^{n}\right) \tag{3.18}
\end{equation*}
$$

with residual spaces
1.2 .2000 .355

$$
\begin{equation*}
\bigcap_{l, m} \Psi_{\infty-\mathrm{sc}}^{l, m}\left(\mathbb{R}^{n}\right)=\bigcap_{l, m} \Psi_{\mathrm{sc}}^{l, m}\left(\mathbb{R}^{n}\right) \Psi_{\mathrm{iso}}^{-\infty}\left(\mathbb{R}^{n}\right)=\mathcal{S}\left(\mathbb{R}^{2 n}\right) \tag{3.19}
\end{equation*}
$$

### 3.3. The residual algebra isotropic algebra

The residual isotropic (and scattering) algebra has two important properties not shared by the residual algebra $\Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$, of which it is a subalgebra (and in fact in which it is an ideal). The first is that as operators on $L^{2}\left(\mathbb{R}^{n}\right)$ the residual isotropic operators are compact.

## 1.2 .2000 .309

1.2 .2000 .356
1.2 .2000 .321

Proposition 3.3. Elements of $\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ are characterized amongst continous operators on $\mathcal{S}\left(\mathbb{R}^{n}\right)$ by the fact that they extend by continuity to define continuous linear maps

$$
\begin{equation*}
A: \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{3.20}
\end{equation*}
$$

In particular the image of a bounded subset of $L^{2}\left(\mathbb{R}^{n}\right)$ under an element of $\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ is contained in a compact subset.

Proof . The kernels of elements of $\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ are in $\mathcal{S}\left(\mathbb{R}^{2 n}\right)$ so the mapping property $\left(\frac{1.2 .2000 .356}{3.20}\right)$ follows.

The norm $\sup _{|\alpha| \leq 1}\left|\langle x\rangle^{n+1} D^{\alpha} u(x)\right|$ is continuous on $\mathcal{S}\left(\mathbb{R}^{n}\right)$. Thus if $S \subset L^{2}\left(\mathbb{R}^{n}\right)$ is bounded and $A \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ the continuity of $A: L^{2}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}\left(\mathbb{R}^{n}\right)$ implies that $A(S)$ is bounded with respect to this norm. The theorem of Arzela-Ascoli shows that any sequence in $A(S)$ has a strongly convergent subsequence in $\langle x\rangle^{n} \mathcal{C}_{\infty}^{0}\left(\mathbb{R}^{n}\right)$ and such a sequence converges in $L^{2}\left(\mathbb{R}^{n}\right)$. Thus $A(S)$ has compact closure in $L^{2}\left(\mathbb{R}^{n}\right)$.

The second important property of the residual algebra is that it is 'bi-ideal' or a 'corner' in the bounded operators on $L^{2}\left(\mathbb{R}^{n}\right)$. Note that it is not an ideal.

Lemma 3.3. If $A_{1}, A_{2} \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ and $B$ is a bounded operator on $L^{2}\left(\mathbb{R}^{n}\right)$ then $A_{1} B A_{2} \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$.

Proof. The kernel of the composite $C=A_{1} B A_{2}$ can be written as a distributional pairing

$$
\begin{equation*}
C(x, y)=\int_{\mathbb{R}^{2 n}} B\left(x^{\prime}, y^{\prime}\right) A_{1}\left(x, x^{\prime}\right) A_{2}\left(y^{\prime}, y\right) d x^{\prime} d y^{\prime}=\left(B, A_{1}(x, \cdot) A_{2}(\cdot, y)\right) \in \mathcal{S}\left(\mathbb{R}^{2 n}\right) \tag{3.21}
\end{equation*}
$$

Thus the result follows from the continuity of the exterior product, $\mathcal{S}\left(\mathbb{R}^{2 n}\right) \times$ $\mathcal{S}\left(\mathbb{R}^{2 n}\right) \longrightarrow \mathcal{S}\left(\mathbb{R}^{4 n}\right)$.

In fact the same conclusion, with essentially the same proof, holds for any continuous linear operator $B$ from $\mathcal{S}\left(\mathbb{R}^{n}\right)$ to $\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$.

### 3.4. The residual isotropic ring

Recall that a bounded operator is said to have finite rank if its range is finite dimensional. If we consider a bounded operator $B$ on $L^{2}\left(\mathbb{R}^{n}\right)$ which is of finite rank then we may choose an orthonormal basis $f_{j}, j=1, \ldots, N$ of the range $B L^{2}\left(\mathbb{R}^{n}\right)$. The functionals $u \longmapsto\left\langle B u, f_{j}\right\rangle$ are continuous and so define non-vanishing elements $g_{j} \in L^{2}\left(\mathbb{R}^{n}\right)$. It follows that the Schwartz kernel of $B$ is

$$
\begin{equation*}
B=\sum_{j=1}^{N} f_{j}(x) \overline{g_{j}(y)} \tag{3.22}
\end{equation*}
$$

If $B \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ then the range must lie in $\mathcal{S}\left(\mathbb{R}^{n}\right)$ and similarly for the range of the adjoint, so the functions $f_{j}$ and $g_{j}$ are also in $\mathcal{S}\left(\mathbb{R}^{n}\right)$. Clearly the finite rank elements of $\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ form an ideal in $\Psi_{\infty-\text { iso }}^{\infty}\left(\mathbb{R}^{n}\right)$.

$$
\begin{equation*}
(\operatorname{Id}+A)(\operatorname{Id}+B)=\operatorname{Id}-\Pi_{1},(\operatorname{Id}+B)(\operatorname{Id}+A)=\operatorname{Id}-\Pi_{0} \tag{3.23}
\end{equation*}
$$

where $\Pi_{0}, \Pi_{1} \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ are the orthogonal projections onto the null spaces of $\operatorname{Id}+A$ and $\operatorname{Id}+A^{*}$ and furthermore, there is an element $A^{\prime} \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ of rank equal to the dimension of the null space such that $\mathrm{Id}+A+s A^{\prime}$ is an invertible operator on $L^{2}\left(\mathbb{R}^{n}\right)$ for all $s \neq 0$.

Proof. Most of these properties are a direct consequence of the fact that $A$ is compact as an operator on $L^{2}\left(\mathbb{R}^{n}\right)$ inevertheless we give brief proofs.

We have shown, in Proposition $\mathbf{3 . 3}$ that each $A \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ is compact. It follows that

$$
\begin{equation*}
N_{0}=\operatorname{Nul}(\operatorname{Id}+A) \subset L^{2}\left(\mathbb{R}^{n}\right) \tag{3.24}
\end{equation*}
$$

has compact unit ball. Indeed the unit ball, $B=\{u \operatorname{Nul}(\operatorname{Id}+A)\}$ satisfies $B=$ $A(B)$, since $u=-A u$ on $B$. Thus $B$ is closed and precompact. Any Hilbert space with a compact unit ball is finite dimensional, so $\operatorname{Nul}(\operatorname{Id}+A)$ is finite dimensional.

Now, let $R_{1}=\operatorname{Ran}(\operatorname{Id}+A)$ be the range of $\operatorname{Id}+A$; we wish to show that this is a closed subspace of $L^{2}\left(\mathbb{R}^{n}\right)$. Let $f_{k} \rightarrow f$ be a sequence in $R_{1}$, converging in $L^{2}\left(\mathbb{R}^{n}\right)$. For each $k$ there exists a unique $u_{k} \in L^{2}\left(\mathbb{R}^{n}\right)$ with $u_{k} \perp N_{0}$ and $(\operatorname{Id}+A) u_{k}=f_{k}$. We wish to show that $u_{k} \rightarrow u$. First we show that $\left\|u_{k}\right\|$ is bounded. If not, then along a subsequent $v_{j}=u_{k(j)},\left\|v_{j}\right\| \rightarrow \infty$. Set $w_{j}=v_{j} /\left\|v_{j}\right\|$. Using the compactness of $A, w_{j}=-A w_{j}+f_{k(j)} /\left\|v_{j}\right\|$ must have a convergent subsequence, $w_{j} \rightarrow w$. Then $(\operatorname{Id}+A) w=0$ but $w \perp N_{0}$ and $\|w\|=1$ which are contradictory. Thus the sequence $u_{k}$ is bounded in $L^{2}\left(\mathbb{R}^{n}\right)$. Then again $u_{k}=-A u_{k}+f_{k}$ has a convergent subsequence with limit $u$ which is a solution of $(\operatorname{Id}+A) u=f$; hence $R_{1}$ is closed. The orthocomplement of the range of a bounded operator is always the null space of its adjoint, so $R_{1}$ has a finite-dimensional complement $N_{1}=\operatorname{Nul}\left(\operatorname{Id}+A^{*}\right)$. The same argument applies to $\mathrm{Id}+A^{*}$ so gives the orthogonal decompositions

[^10]\[

$$
\begin{align*}
& L^{2}\left(\mathbb{R}^{n}\right)=N_{0} \oplus R_{0}, \quad N_{0}=\operatorname{Nul}(\operatorname{Id}+A), \quad R_{0}=\operatorname{Ran}\left(\operatorname{Id}+A^{*}\right) \\
& L^{2}\left(\mathbb{R}^{n}\right)=N_{1} \oplus R_{1}, \quad N_{1}=\operatorname{Nul}\left(\operatorname{Id}+A^{*}\right), R_{1}=\operatorname{Ran}(\operatorname{Id}+A) \tag{3.25}
\end{align*}
$$
\]

Thus we have shown that $\mathrm{Id}+A$ induces a continuous bijection $\tilde{A}: R_{0} \longrightarrow R_{1}$. From the closed graph theorem the inverse is a bounded operator $\tilde{B}: R_{1} \longrightarrow R_{0}$. In this case continuity also follows from the argument above. ${ }_{\tilde{B}}$ Thus $\tilde{B}_{0}$ is the generalized inverse of $\operatorname{Id}+A$ in the sense that $B=\tilde{B}-\operatorname{Id}$ atisfies ( 33.23$)$. It only remains to show that $B \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$. This follows from ( $\left(\frac{1}{3.23}\right)$, the identities in which show that
1.2 .2000 .329

$$
\begin{align*}
B=-A-A B-\Pi_{1},-B=A & +B A+\Pi_{0}  \tag{3.26}\\
& \Longrightarrow B=-A+A^{2}+A B A-\Pi_{1}+A \Pi_{0}
\end{align*}
$$

[^11]All terms here are in $\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$; for $A B A$ this follows from Proposition $\frac{1}{3} .2 .2$.
It remains to show the existence of the finite rank perturbation $A^{\prime}$. This is equivalent to the vanishing of the index, that is
1.2 .2000 .323

$$
\begin{equation*}
\operatorname{Ind}(\operatorname{Id}+A)=\operatorname{dim} \operatorname{Nul}(\operatorname{Id}+A)-\operatorname{dim} \operatorname{Nul}\left(\operatorname{Id}+A^{*}\right)=0 \tag{3.27}
\end{equation*}
$$

Indeed, let $f_{j}$ and $g_{j}, j=1, \ldots, N$, be respective bases of the two finite dimensional spaces $\operatorname{Nul}(\operatorname{Id}+A)$ and $\operatorname{Nul}\left(\operatorname{Id}+A^{*}\right)$. Then

$$
\begin{equation*}
A^{\prime}=\sum_{j=1}^{N} g_{j}(x) \overline{f_{j}(y)} \tag{3.28}
\end{equation*}
$$

is an isomorphism of $N_{0}$ onto $N_{1}$ which vanishes on $R_{0}$. Thus Id $+A+s A^{\prime}$ is the direct sum of $\operatorname{Id}+A$ as an operator from $R_{0}$ to $R_{1}$ and $s A^{\prime}$ as an operator from $N_{0}$ to $N_{1}$, invertible when $s \neq 0$.

There is a very simple proof $f^{2}$ of the equality $\frac{(1,2.2000 .323}{3.27) \text { it we use the trace func- }}$ tional discussed in Section $\frac{1}{3.14}$ below; this however is logically suspect as we use (although not crucially) approximation by finite rank operators in the discussion of the trace and this in turn might appear to use the present result via the discussion of ellipticity and the harmonic oscillator. Even though this is not really the case we give a clearly independent, but less elegant proof.

Consider the one-parameter family of operators $\operatorname{Id}+t A, A \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$. We shall see that the index, the difference in dimension between $\operatorname{Nul}(\operatorname{Id}+t A)$ and $\operatorname{Nul}\left(\operatorname{Id}+t A^{*}\right)$ is locally constant. To see this it is enough to consider a general $A$ near the point $t=1$. Consider the pieces of $A$ with respect to the decompositions $L^{2}\left(\mathbb{R}^{n}\right)=N_{i} \oplus R_{i}, i=0,1$, of domain and range. Thus $A$ is the sum of four terms which we write as a $2 \times 2$ matrix

$$
A=\left[\begin{array}{ll}
A_{00} & A_{01} \\
A_{10} & A_{11}
\end{array}\right]
$$

Since $\operatorname{Id}+A$ has only one term in such a decomposition, $\tilde{A}$ in the lower right, the solution of the equation $(\operatorname{Id}+t A) u=f$ can be written
(3.29) $(t-1) A_{00} u_{0}+(t-1) A_{01} u_{\perp}=f_{1},(t-1) A_{10} u_{0}+\left(A^{\prime}+(t-1) A_{11}\right) u_{\perp}=f_{\perp}$

Since $\tilde{A}$ is invertible, for $t-1$ small enough the second equation can be solved uniquely for $u_{\perp}$. Inserted into the first equation this gives
1.2 .2000 .326

$$
\begin{align*}
& G(t) u_{0}=f_{1}+H(t) f_{\perp}  \tag{3.30}\\
& \qquad \begin{aligned}
& G(t)=(t-1) A_{00}-(t-1)^{2} A_{01}\left(A^{\prime}+(t-1) A_{11}\right)^{-1} A_{10} \\
& H(t)=-(t-1) A_{01}\left(A^{\prime}+(t-1) A_{11}\right)^{-1}
\end{aligned}
\end{align*}
$$

The null space is therefore isomorphic to the null space of $G(t)$ and a complement to the range is isomorphic to a complement to the range of $G(t)$. Since $G(t)$ is a finite rank operator acting from $N_{0}$ to $N_{1}$ the difference of these dimension is constant in $t$, namely equal to $\operatorname{dim} N_{0}-\operatorname{dim} N_{1}$, near $t=1$ where it is defined.

[^12]This argument can be applied to $t A$ so the index is actually constant in $t \in[0,1]$ and since it certainly vanishes at $t=0$ it vanishes for all $t$. In fact, as we shall note below, $\operatorname{Id}+t A$ is invertible outside a discrete set of $t \in \mathbb{C}$.
1.2 .2000 .333

COROLLARY 3.1. If $\operatorname{Id}+A, A \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ is injective or surjective on $L^{2}\left(\mathbb{R}^{n}\right)$, ring $\operatorname{Id}+\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$.

Corollary 3.2. If $A \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ then as an operator on $\mathcal{S}\left(\mathbb{R}^{n}\right)$ or $\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$, $\mathrm{Id}+A$ is Fredholm in the sense that its null space is finite dimensional and its range is closed with a finite dimensional complement.

Proof. This follows from the existence of the generalized inverse of the form $\operatorname{Id}+B, B \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$.

### 3.5. Exponential and logarithm

1.2 .2000 .350
1.2 .2000 .351

$$
\begin{equation*}
\exp (A)=\sum_{j} \frac{1}{j!} A^{j}: \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right) \longrightarrow \operatorname{Id}+\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right) \tag{3.31}
\end{equation*}
$$

is a globally defined, entire, function with range containing a neighbourhood of the identity and with inverse on such a neighbourhood given by the analytic function

$$
1.2 .2000 .352
$$

$$
\begin{equation*}
\log (\operatorname{Id}+A)=\sum_{j} \frac{(-1)^{j}}{j} A^{j}, A \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right),\|A\|_{L^{2}}<1 \tag{3.32}
\end{equation*}
$$

### 3.6. Fredholm property

An element $A \in \Psi_{\infty-\text { iso }}^{m}\left(\mathbb{R}^{n}\right)$ is said to be elliptic (of order $m$ in the isotropic calculus) if its left-reduced symbol is elliptic in $S_{\infty}^{m}\left(\mathbb{R}^{2 n}\right)$.

Theorem 3.3. Each elliptic element $A \in \Psi_{\infty-\mathrm{iso}}^{m}\left(\mathbb{R}^{n}\right)$ has a two-sided parametrix $B \in \Psi_{\infty-\text { iso }}^{-m}\left(\mathbb{R}^{n}\right)$ in the sense that
1.2 .2000 .365

$$
\begin{equation*}
A \circ B-\mathrm{Id}, B \circ A-\mathrm{Id} \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right) \tag{3.33}
\end{equation*}
$$

and it follows that any $u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ satisfying $A u \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ is an element of $\mathcal{S}\left(\mathbb{R}^{n}\right)$; if $A \in \Psi_{\text {iso }}^{m}\left(\mathbb{R}^{n}\right)$ is elliptic then its parametrix is in $\Psi_{\text {iso }}^{-m}\left(\mathbb{R}^{n}\right)$.

Proof. This is just the inductive argument used to prove Lemma $\frac{5.31}{2.7 .}$ Nevertheless we repeat it here.

The existence of a right inverse for $\sigma_{k}(A)$ means that the equation $\sigma_{k}(A) c=d$ always has a solution $c \in \mathcal{A}_{j-k}$ for given $d \in \mathcal{A}_{j}$, namely $c=b d$. This in turn means that given $C_{j} \in \Psi_{\infty-\text { iso }}^{j}\left(\mathbb{R}^{n}\right)$ there always exists $B_{j} \in \Psi_{\infty-\text { iso }}^{j-k}\left(\mathbb{R}^{n}\right)$ such that $A B_{j}-D_{j} \in \Psi_{\infty-\text { iso }}^{j-1}\left(\mathbb{R}^{n}\right)$. Choosing $B_{0} \in \Psi_{\infty-\text { iso }}^{-k}\left(\mathbb{R}^{n}\right)$ to have $\sigma_{-k}\left(B_{0}\right)=b$ we can define $C_{1}=\operatorname{Id}-A B_{0} \in \Psi_{\infty-\text { iso }}^{-1}\left(\mathbb{R}^{n}\right)$. Then, proceeding inductively we may assume that $B_{j}$ for $j<l$ have been chosen such that $A\left(B_{0}+\cdots+B_{l-1}\right)-$ Id $=-C_{l} \in \Psi_{\infty-\text { iso }}^{-l}\left(\mathbb{R}^{n}\right)$. Then using the solvability we may choose $B_{l}$ so that $A B_{l}-C_{l}=-C_{l+1} \in \Psi_{\infty-\text { iso }}^{-l-1}\left(\mathbb{R}^{n}\right)$ which completes the induction, since $A\left(B_{0}+\right.$ $\left.\cdots+B_{l}\right)-\mathrm{Id}=A B_{l}-C_{l}=-C_{l+1}$. Finally by the asymptotic completeness we may choose $B \sim B_{0}+B_{1}+\ldots$ which is a right parametrix.

The argument showing the existence of a left parametrix for a left-elliptic operator is completely analogous.

Combining the earlier symbolic discussion and these analytic results we can see that elliptic operators in these calculi are Fredholm.

$$
\begin{equation*}
A B-\mathrm{Id}=\Pi_{1}, B A-\mathrm{Id}=\Pi_{0} \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right) \tag{3.34}
\end{equation*}
$$

where $\Pi_{1}$ and $\Pi_{0}$ are the finite rank orthogonal (in $L^{2}\left(\mathbb{R}^{n}\right)$ ) projections onto the null spaces of $A^{*}$ and $A$.

Proof. In the case of an elliptic isotropic operator or order $m$ we know that it has a parametrix $B^{\prime} \in \Psi_{\text {iso }}^{-m}\left(\mathbb{R}^{n}\right)$ modulo $\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$. Thus

$$
\begin{aligned}
& A B^{\prime}=\operatorname{Id}-E_{R}, E_{R} \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right) \\
& B^{\prime} A=\operatorname{Id}-E_{L}, E_{L} \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)
\end{aligned}
$$

Using Proposition $\frac{11.2 .2000 .312}{3.4 \text { it follows that the null space of } A \text { is contained in the null space }}$ of $B^{\prime} A=\operatorname{Id}-E_{L}$, hence is finite dimensional. Similarly, the range of $A$ contains the range of $A B^{\prime}=\mathrm{Id}-E_{R}$ so is closed with a finite codimensional complement. Defining $B$ as the linear map which vanishes $\left.{ }_{21} 3.1998 .1 A_{0}^{*}\right)$, and inverts $A$ on $\operatorname{Ran}(A)$ with values in $\operatorname{Ran}\left(A^{*}\right)=\operatorname{Nul}(A)^{\perp}$ gives (3.34). Furthermore these identities show that $B \in \Psi_{\infty-\text { iso }}^{-m}\left(\mathbb{R}^{n}\right)$ since applying $B^{\prime}$ gives

$$
\begin{array}{|l|}
\hline 1.2 .2000 .331  \tag{3.35}\\
\hline
\end{array}
$$

$$
\begin{array}{r}
B-E_{L} B=B^{\prime} A B=B^{\prime}-B^{\prime} \Pi_{1}, B-B E_{R}=B A B^{\prime}=B^{\prime}-\Pi_{0} B^{\prime} \Longrightarrow \\
B=B^{\prime}-B^{\prime} \Pi_{1}+E_{L} B^{\prime}+E_{L} B E_{R}-E_{L} \Pi_{0} B^{\prime} \in \Psi_{\infty-\text { iso }}^{-m}\left(\mathbb{R}^{n}\right) .
\end{array}
$$

The proof in the scattering case is essentially the same.
1.2 .2000 .332

Corollary 3.3. If $A \in \Psi_{\text {iso }}^{m}\left(\mathbb{R}^{n}\right)$ is elliptic then its generalized inverse lies in $\Psi_{\mathrm{iso}}^{-m}\left(\mathbb{R}^{n}\right)$ and similarly if $A \in \Psi_{\mathrm{sc}}^{l, m}\left(\mathbb{R}^{n}\right)$ is elliptic then its generalized inverse lies in $\Psi_{\mathrm{sc}}^{-l,-m}\left(\mathbb{R}^{n}\right)$.

### 3.7. The harmonic oscillator

The harmonic oscillator is the differential operator on $\mathbb{R}^{n}$

$$
H=\sum_{j=1}^{n}\left(D_{j}^{2}+x_{j}^{2}\right)=\Delta+|x|^{2}
$$

This is an elliptic element of $\Psi_{\text {iso }}^{2}\left(\mathbb{R}^{n}\right)$. The main immediate interest is in the spectral decomposition of $H$. The ellipticity of $H-\lambda, \lambda \in \mathbb{C}$, shows that

$$
\begin{equation*}
(H-\lambda) u=0, u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \Longrightarrow u \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{3.36}
\end{equation*}
$$

Since $H$ is (formally) self-adjoint, i.e., $H^{*}=H$, there are no non-trivial tempered solutions of $(H-\lambda) u=0, \lambda \notin \mathbb{R}$. Indeed if $(H-\lambda) u=0$,

$$
\begin{equation*}
0=\langle H u, u\rangle-\langle u, H u\rangle=(\lambda-\bar{\lambda})\langle u, u\rangle \Longrightarrow u=0 . \tag{3.37}
\end{equation*}
$$

As we shall see below in more generality, the spectrum of $H$ is a discrete subset of $\mathbb{R}$. In this case we can compute it explicitly.

The direct computation of eigenvalues and eigenfunctions is based on the properties of the creation and annihilation operators
eq:HO. 3

$$
\begin{equation*}
C_{j}=D_{j}+i x_{j}, \quad C_{j}^{*}=A_{j}=D_{j}-i x_{j}, j=1, \ldots, n \tag{3.38}
\end{equation*}
$$

These satisfy the elementary identities
eq:HO. 4

$$
\begin{align*}
& {\left[A_{j}, A_{k}\right]=\left[C_{j}, C_{k}\right]=0,\left[A_{j}, C_{k}\right]=2 \delta_{j k}, \quad j, k=1, \ldots, n} \\
& H=\sum_{j=1}^{n} C_{j} A_{j}+n, \quad\left[C_{j}, H\right]=-2 C_{j}, \quad\left[A_{j}, H\right]=2 A_{j} \tag{3.39}
\end{align*}
$$

Now, if $\lambda$ is an eigenvalue, $H u=\lambda u$, then
eq:HO. 5
prop:HO. 6
eq:HO. 7

$$
\begin{align*}
& H\left(C_{j} u\right)=C_{j}(H u+2 u)=(\lambda+2) C_{j} u \\
& H\left(A_{j} u\right)=A_{j}(H u-2 u)=(\lambda-2) A_{j} u \tag{3.40}
\end{align*}
$$

Proposition 3.7. The eigenvalues of $H$ are

$$
\begin{equation*}
\sigma(H)=\{n, n+2, n+4, \ldots\} \tag{3.41}
\end{equation*}
$$

Proof. We already know that eigenvalues must be real and from the decomposition of $H$ in (3.40) it follows that, for $u \in \mathcal{S}\left(\mathbb{R}^{n}\right)$,
eq:HO. 9

$$
\begin{equation*}
\langle H u, u\rangle=\sum_{j}\left\|A_{j} u\right\|^{2}+n\|u\|^{2} . \tag{3.42}
\end{equation*}
$$

Thus if $\lambda \in \sigma(H)$ is an eigenvalue then $\lambda \geq n$.
By direct computation we see that $n$ is an eigenvalue with a 1-dimensional eigenspace. Indeed, from (3.42), $H u=n u$ iff $A_{j} u=0$ for $j=1, \ldots, n$. In each variable separately

$$
A_{j} u\left(x_{j}\right)=0 \Leftrightarrow u\left(x_{j}\right)=c_{1} \exp \left(-\frac{x_{j}^{2}}{2}\right) .
$$

Thus the only tempered solutions of $A_{j} u=0, i=1, \ldots, n$ are the constant multiples of

$$
\begin{equation*}
u_{0}=\exp \left(-\frac{|x|^{2}}{2}\right) \tag{3.43}
\end{equation*}
$$

which is often called the ground state.
Now, if $\lambda$ is an eigenvalue with eigenfunction $u \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ it follows from ( k (3.40). 5 that $\lambda-2$ is an eigenvalue with eigenfunction $A_{j} u$. Since all the $A_{j} u$ cannot vanish unless $u$ is the ground state, it follows that the eigenvalues are contained in the set in (3.41). We can use the same argument to show that if $u$ is an eigenfunction with eigenvalue $\lambda$ then $C_{j} u$ is an eigenfunction with eigenvalue $\lambda+2$. Moreover, $C_{j} u \equiv 0$ would imply $u \equiv 0$ since $C_{j} v=0$ has no non-trivial tempered solutions, the solution in each variable being $\exp \left(x_{j}^{2} / 2\right)$.

Using the creation operators we can parameterize the eigenspaces quite explicitly.

Proposition 3.8. For each $k \in \mathbb{N}_{0}$ there is an isomorphism
(3.44) $\quad\left\{\right.$ Polynomials, homogeneous of degree $k$ on $\left.\mathbb{R}^{n}\right\} \ni p$

$$
\longmapsto p(C) \exp \left(-\frac{|x|^{2}}{2}\right) \in E_{k}
$$

where $E_{k}$ is the eigenspace of $H$ with eigenvalue $n+2 k$.

Proof. Notice that the $C_{j}, j=\overline{\bar{H}}, 1_{5} \ldots, n$ are commuting operators, so $p(C)$ is well-defined. By iteration from (3.40),

$$
\begin{equation*}
H C^{\alpha} u_{0}=C^{\alpha}(H+2|\alpha|) u_{0}=(n+2|\alpha|) C^{\alpha} u_{0} \tag{3.45}
\end{equation*}
$$

Thus ( $\mathrm{eq}: \mathrm{HO}$. 12 is a linear map into the eigenspace as indicated.
To see that (3.44) 1 s an isomorphism consider the action of the annihilation operators. Again from (3.40)

$$
|\beta|=|\alpha| \Longrightarrow A^{\beta} C^{\alpha} u_{0}= \begin{cases}0 & \beta \neq \alpha  \tag{3.46}\\ 2^{|\alpha|} \alpha!u_{0} & \beta=\alpha\end{cases}
$$

This allows us to recover the coefficients of $p$ from $p(C) u_{0}$, so ( $\frac{\text { (eq:HO. } 12}{}$. 12 injective. Conversely if $v \in E_{k} \subset \mathcal{S}\left(\mathbb{R}^{n}\right)$ is orthogonal to all the $C^{\alpha} u_{0}$ then

$$
\begin{equation*}
\left\langle A^{\alpha} v, u_{0}\right\rangle=\left\langle v, C^{\alpha} u_{0}\right\rangle=0 \forall|\alpha|=k . \tag{3.47}
\end{equation*}
$$

From (k. (3.40), the $A^{\alpha} v$ are all eigenfunctions of $H$ with eigenvalue $n$, so ( $\frac{\text { eq: } \mathrm{HO} .15}{3.47 \text { ) } 1 \mathrm{~m}}$ plies that $A^{\alpha} v=0$ for all $|\alpha|=k$. Proceeding inductively in $k$ we see that $A^{\alpha^{\prime}} A_{j} v=0$ for all $\left|\alpha^{\prime}\right|=k-1$ and $A_{j} v \in E_{k-1}$ implies $A_{\text {prop: }} \mathrm{Ap}_{\mathrm{p}} v_{11}=0, j=1, \ldots, n$. Since $v \in E_{k}$, $k>0$, this implies $v=0$ so Proposition 3.8 is proved.

Thus $H$ has eigenspaces as described in (3.44). The same argument shows that for any integer $p$, positive or negative, the eigenvalues of $H^{p}$ are precisely $(n+2 k)^{p}$ with the same eigenspaces $E_{k}$. For $p<0, H^{p}$ is a compact operator on $L^{2}\left(\mathbb{R}^{n}\right)$; this is obvious for large negative $p$. For example, if $p \leq-n-1$ then

$$
\begin{equation*}
x_{i}^{\beta} D_{j}^{\alpha} H \in \Psi_{\text {iso }}^{0}\left(\mathbb{R}^{n}\right),|\alpha| \leq n+1,|\beta| \leq n+1 \tag{3.48}
\end{equation*}
$$

are all bounded on $L^{2}$. If $S \subset L^{2}\left(\mathbb{R}^{n}\right)$ is bounded this implies that $H^{-n-1}(S)$ is bounded in $\langle x\rangle^{n+1} C_{\infty}^{1}\left(\mathbb{R}^{n}\right)$, so compact in $\langle x\rangle^{n} C_{\infty}^{0}\left(\mathbb{R}^{n}\right)$ and hence in $L^{2}\left(\mathbb{R}^{n}\right)$. It is a general fact that for compact self-adjoint operators, such as $H^{-n-2}$, the eigenfunctions span $L^{2}\left(\mathbb{R}^{n}\right)$. We give a brief proof of this for the sake of 'completeness'.
lem:HO.17
Lemma 3.4. The eigenfunction of $H, u_{\alpha}=\pi^{-\frac{n}{4}}\left(2^{|\alpha|} \alpha!\right)^{-1 / 2} C^{\alpha} u_{0}$ form an orthonormal basis of $L^{2}\left(\mathbb{R}^{n}\right)$.

Proof. Let $V \subset L^{2}\left(\mathbb{R}^{n}\right)$ be the closed subspace consisting of the orthocomplements of all the $u_{\alpha}$ 's. Certainly $H^{-n-2}$ acts on it as a compact self-adjoint operator. Since we have found all the eigenvalues of $H$, and hence of $H^{-n-1}$, it has no eigenvalue in $V$. We wish to conclude that $V=\{0\}$. Set

$$
\tau=\left\|H^{-n-1}\right\|_{V}=\sup \left\{\left\|H^{-n-1} \varphi\right\| ; \varphi \in V,\|\varphi\|=1\right\}
$$

Then there is a weakly convergent sequence $\varphi_{j} \rightharpoonup \varphi,\left\|\varphi_{j}\right\|=1$, so $\|\varphi\| \leq 1$, with $\left\|H^{-n-1} \varphi_{j}\right\| \rightarrow \tau$. The compactness of $H^{-n-2}$ allows a subsequence to be chosen such that $H^{-n-1} \varphi_{j} \rightarrow \psi$ in $L^{2}\left(\mathbb{R}^{n}\right)$. So, by the continuity of $H^{-n-1}, H^{-n-1} \varphi=\psi$ and $\left\|H^{-n-1} \varphi\right\|=\tau,\|\varphi\|=1$. If $\varphi^{\prime} \in V, \varphi^{\prime} \perp \varphi,\left\|\varphi^{\prime}\right\|=1$ then

$$
\begin{aligned}
& \tau^{2} \geq \| H^{-n-2}\left(\frac{\varphi+t \varphi^{\prime}}{\sqrt{1+t^{2}}}\right) \|^{2}=\tau^{2}+2 t\left\langle H^{-2 n-2} \varphi, \varphi^{\prime}\right\rangle+0\left(t^{2}\right) \\
& \Longrightarrow\left\langle H^{-2 n-2} \varphi, \varphi^{\prime}\right\rangle=0 \Longrightarrow H^{-2 n-2} \varphi=\tau^{2} \varphi
\end{aligned}
$$

This contradicts the fact that $H^{-2 n-2}$ has no eigenvalues in $V$, so $V=\{0\}$ and the eigenbasis is complete.

Thus, if $u \in L^{2}\left(\mathbb{R}^{n}\right)$ then

$$
\begin{equation*}
u=\sum_{\alpha} c_{\alpha} u_{\alpha}, \quad c_{\alpha}=\left\langle u, u_{\alpha}\right\rangle \tag{3.49}
\end{equation*}
$$

Lemma 3.5. If $u \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ the convergence in (3.49. 18 is rapid, i.e., $\left|C_{\alpha}\right| \leq$ $C_{N}(1+|\alpha|)^{-N}$ for all $N$ and the series converges in $\mathcal{S}\left(\mathbb{R}^{n}\right)$.

Proof. Since $u \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ implies $H^{N} u \in L^{2}\left(\mathbb{R}^{n}\right)$ we see that

$$
C_{N} \geq\left|\left\langle H^{N} u, u_{\alpha}\right\rangle\right|=\left|\left\langle u, H^{N} u_{\alpha}\right\rangle\right|=(n+2|\alpha|)^{N}\left|c_{\alpha}\right| \forall \alpha .
$$

Furthermore, $2 i x_{j}=C_{j}-A_{j}$ and $2 D_{j}=C_{j}+A_{j}$ so the polynomial derivatives of the $u_{\alpha}$ can be estimated (using the Sobolev embedding theorem) by polynomials in $\alpha$; this implies that the series converges in $\mathcal{S}\left(\mathbb{R}^{n}\right)$.

## 1.2 .2000 .407

then
1.2 .2000 .358

$$
\begin{equation*}
\inf _{E \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)}\|A+E\|_{\mathcal{B}\left(L^{2}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)\right.} \leq \sup _{p \in \mathbb{S}^{2 n-1}}\|a(p)\| \tag{3.50}
\end{equation*}
$$

Proof. It suffices to prove ( $\frac{1.2 .2000 .358}{3.50)}$ for all single operators $A \in \Psi_{\text {iso }}^{0}\left(\mathbb{R}^{n}\right)$. Indeed if $j_{v}(z)=z v$ is the linear map from $\mathbb{C}$ to $\mathbb{C}^{N}$ defined by $v \in \mathbb{C}^{N}$ then

$$
\begin{array}{|l|}
\hline 1.2 .2000 .363 \\
\hline
\end{array}
$$

$$
\begin{equation*}
\|A\|_{\mathcal{B}\left(L^{2}\left(\mathbb{R} ; \mathbb{C}^{N}\right)\right)}=\sup _{\left\{v, w \in \mathbb{C}^{N} ;\|v\|=\|w\|=1\right\}}\left\|j_{w}^{*} A j_{v}\right\|_{\mathcal{B}\left(L^{2}(\mathbb{R})\right)} \tag{3.51}
\end{equation*}
$$

Since the symbol of $j_{w}^{*} A j_{v}$ is just $j_{w}^{*} \sigma(A) j_{v}$, (1.2.2000.358 3.50$)$ follows from the corresponding equality for a single operator:
1.2 .2000 .364

$$
\begin{equation*}
\inf _{E \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)}\|A+E\|_{\mathcal{B}\left(L^{2}\left(\mathbb{R}^{n}\right)\right.} \leq \sup _{p \in \mathbb{S}^{2 n-1}}|a(p)|, a=\left.\sigma_{L}(A)\right|_{\mathbb{S}^{2 n-1}} \tag{3.52}
\end{equation*}
$$

The construction of the approximate square-root of $C-A^{*} A$ in Proposition $\frac{6.6}{2.7}$ only depends on the existence of a positive smooth square-root for $C-|a|^{2}$, so can be carried out for any

$$
\begin{equation*}
C>\sup _{p \in \mathbb{S}^{2 n-1}}|a(p)|^{2} \tag{3.53}
\end{equation*}
$$

Thus we conclude that with such a value of $C$

$$
\|A u\|^{2} \leq C\|u\|^{2}+\|\langle G u, u\rangle \mid \forall u \in L^{2}\left(\mathbb{R}^{n}\right)
$$

where $G \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$. Since $G$ is an isotropic smoothing operator, for any $\delta>0$ there is a finite dimensional subspace $W \subset \mathcal{S}\left(\mathbb{R}^{n}\right)$ such that

$$
\begin{equation*}
\|\langle G u, u\rangle\| \leq \delta\|u\|^{2} \forall u \in W^{\perp} \tag{3.54}
\end{equation*}
$$

Thus if we replace $A$ by $A\left(\operatorname{Id}-\Pi_{W}\right)=A+E$ where $E$ is a (finite rank) smoothing operator we see that

$$
\|(A+E) u\|^{2} \leq(C+\delta)\|G u\|^{2} \forall u \in \underset{\|^{1} .2 .2000 .363}{L^{2}\left(\mathbb{R}^{n}\right)} \Longrightarrow\|(A+E)\| \leq(C+\delta)^{\frac{1}{2}}
$$

This proves half of the desired estimate ( $\left(\frac{1}{3} \cdot 2.21\right)$, namely

$$
\begin{equation*}
\inf _{E \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)}\|A+E\|_{\mathcal{B}\left(L^{2}\left(\mathbb{R}^{n}\right)\right.} \leq \sup _{p \in \mathbb{S}^{2} n-1}|a(p)| \tag{3.55}
\end{equation*}
$$

To prove the opposite inequality, leading to $\left(\frac{1.2 .2000 .362}{3.55)}\right.$, it is enough to arrive at a contradiction by supposing to the contrary that there is some $A \in \Psi_{\text {iso }}^{0}\left(\mathbb{R}^{n}\right)$ satisfying the strict inequality

$$
\|A\|_{\mathcal{B}\left(L^{2}\left(\mathbb{R}^{n}\right)\right.}<\sup _{p \in \mathbb{S}^{2 n-1}}|a(p)| .
$$

From this it follows that we may choose $c>0$ such that $c=|a(p)|^{2}$ for some $p \in$ $\mathbb{S}^{2 n-1}$ and yet $A^{\prime}=A^{*} A-c$ has a bounded inverse, $B$. By making an arbitrariy small perturbation of the full symbol of $A^{\prime}$ we may assume that it vanishes identically near $p$. By ( $\mathbf{3} .5 \mathrm{~J}$ ) we may choose $G \in \Psi_{\text {iso }}^{0}\left(\mathbb{R}^{n}\right)$ with arbitrariy small $L^{2}$ such that $\tilde{A}=A^{\prime}+B$ has left symbol rapidly vanishing near $p$. When the norm of the perturbation is small enough, $\tilde{A}$ will still be invertible, with inverse $\tilde{B} \in \mathcal{B}\left(L^{2}\left(\mathbb{R}^{n}\right)\right.$. Now choose an element $G \in \Psi_{\text {iso }}^{0}\left(\mathbb{R}^{n}\right)$ with left symbol supported sufficiently near $p$, so that $G \circ \tilde{A} \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ but yet the principal symbol of $G$ should not vanish at $p$. Thus

$$
\begin{aligned}
G=G \circ \tilde{A} \circ \tilde{B}: L^{2}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}\left(\mathbb{R}^{n}\right) & \\
& G *=G=\tilde{B}^{*} \circ \tilde{A}^{*} \circ G^{*}: \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \longrightarrow L^{2}\left(\mathbb{R}^{n}\right) .
\end{aligned}
$$

It follows that $G^{*} G: \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}\left(\mathbb{R}^{n}\right)$ is an isotropic smoothing operator. This is the expected contradiction, since $G$, and hence $G^{*} G$, may . be ${ }_{2} 860.3602$ to have non-vanishing principal symbol at $p$. Thus we have proved ( 3.55 ) and hence the Proposition.

It is then easy to characterize the compact operators amongst the polyhomogeneous isotropic operators as those of negative.

Lemma 3.6. If $A \in \Psi_{\text {iso }}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$ then, as an operator on $L^{2}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$, $A$ is compact if and only if it has negative order.

Proof. The necessity of vanishing of the principal symbol for a compact operator follows from Proposition $\frac{1}{3.2}$ and the sufficiency follows from the density of $\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$ in $\Psi_{\text {iso }}^{-1}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$ in the topology of $\Psi_{\infty-\text { iso }}^{-\frac{1}{2}}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$ and hence in the topology of bounded operators. Thus, such an operator is the norm limit of compact operators so itself is compact.
 sumption of ellipticity in Proposition 3.6 .

Corollary 3.5. If $A \in \Psi_{\text {iso }}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$ then $A$ is Fredholm as an operator on $L^{2}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$ if and only if it is elliptic.

### 3.9. Sobolev spaces

The space of square-integrable functions plays a basic rôle in the theory of distributions; one reason for this is that it is associated with the embedding of $\mathcal{S}\left(\mathbb{R}^{n}\right)$ in $\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$. We know that pseudodifferential operators of order 0 are bounded on $L^{2}\left(\mathbb{R}^{n}\right)$. There is also a natural collection of Sobolev spaces associated to the isotropic calculus, and another associated to the scattering calculus. The isotropic Sobolev space of order $m$ may be defined as the collection of distributions mapped in $L^{2}\left(\mathbb{R}^{n}\right)$ by any one elliptic operator of order $-m$. Correspondingly the scattering Sobolev spaces have two orders.

Note that a differential operator $P\left(x, D_{x}\right)$ on $\mathbb{R}^{n}$ is an isotropic pseudodifferential operator if and only if its coefficients are polynomials. The fundamental symmetry between coefficients and differentiation suggest that the isotropic Sobolev spaces of non-negative integral order be defined by
1.2 .2000 .337

$$
\begin{equation*}
H_{\mathrm{iso}}^{k}\left(\mathbb{R}^{n}\right)=\left\{u \in L^{2}\left(\mathbb{R}^{n}\right) ; x^{\alpha} D_{x}^{\beta} u \in L^{2}\left(\mathbb{R}^{n}\right) \text { if }|\alpha|+|\beta| \leq k\right\}, k \in \mathbb{N} \tag{3.56}
\end{equation*}
$$

The norms

$$
\begin{equation*}
\|u\|_{k, \text { iso }}^{2}=\sum_{|\alpha|+|\beta| \leq k} \int_{\mathbb{R}^{n}}\left|x^{\alpha} D_{x}^{\beta} u\right|^{2} d x \tag{3.57}
\end{equation*}
$$

turn these into Hilbert spaces. For negative integral orders we identify the isotropic Sobolev spaces with the duals of these spaces

$$
\begin{equation*}
H_{\mathrm{iso}}^{k}\left(\mathbb{R}^{n}\right)=\left(H_{\mathrm{iso}}^{-k}\left(\mathbb{R}^{n}\right)\right)^{\prime} \hookrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right), k \in-\mathbb{N} \tag{3.58}
\end{equation*}
$$

The (continuous) injection into tempered distributions here arises from the density of the image of the inclusion $\mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow H_{\text {iso }}^{k}\left(\mathbb{R}^{n}\right)$.

## 1.2 .2000 .338

1.2 .2000 .341
1.2 .2000 .344

$$
\begin{equation*}
q_{k}(x, \xi)=\sum_{|\alpha|+|\beta|=k}\left(x^{\alpha} \xi^{\beta}\right)^{2} \geq c\left(|x|^{2}+|\xi|^{2}\right)^{k}, c>0 \tag{3.60}
\end{equation*}
$$

Thus $Q_{k}=\sum_{|\alpha|+|\beta| \leq k}\left(D^{\beta} x^{\alpha} x^{\alpha} D^{\beta}\right) \in \Psi_{\text {iso }}^{2 k}\left(\mathbb{R}^{n}\right)$, which has principal reduced symbol $q_{k}$, has a left parameterix $A_{k} \in \Psi_{\text {iso }}^{-2 k}\left(\mathbb{R}^{n}\right)$. This gives the identity

Lemma 3.7. For any $k \in \mathbb{Z}$,

$$
\begin{align*}
H_{\mathrm{iso}}^{k}\left(\mathbb{R}^{n}\right) & =\left\{u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) ; A u \in L^{2}\left(\mathbb{R}^{n}\right) \forall A \in \Psi_{\mathrm{iso}}^{-k}\right\}  \tag{3.59}\\
& =\left\{u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) ; \exists A \in \Psi_{\text {iso }}^{-k} \text { elliptic and such that } A u \in L^{2}\left(\mathbb{R}^{n}\right)\right\}
\end{align*}
$$

and $\mathcal{S}\left(\mathbb{R}^{n}\right) \hookrightarrow H_{\mathrm{iso}}^{k}\left(\mathbb{R}^{n}\right)$ is dense for each $k \in \mathbb{Z}$.
Proof. ${ }^{3}$ For $k \in \mathbb{N}$, the functions $x^{\alpha} \xi^{\beta}$ for $|\alpha|+|\beta|=k$ are 'collectively elliptic' in the sense that
(

$$
\begin{aligned}
& \sum_{|\alpha|+|\beta| \leq k} R_{\alpha, \beta} x^{\alpha} D^{\beta}=A_{k} Q_{k}=\mathrm{Id}+E, \text { where } \\
& R_{\alpha, \beta}=A_{k} D^{\beta} x^{\alpha} \in \Psi_{\text {iso }}^{-2 k+|\alpha|+|\beta|}\left(\mathbb{R}^{n}\right), E \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right) .
\end{aligned}
$$

[^13]Thus if $A \in \Psi_{\text {iso }}^{k}\left(\mathbb{R}^{n}\right)$

$$
A u=-A E u+\sum_{|\alpha|+|\beta| \leq k} A R_{\alpha, \beta} x^{\alpha} D^{\beta} u
$$

If $u \in H_{\text {iso }}^{k}\left(\mathbb{R}^{n}\right)$ then by definition $x^{\alpha} D^{\beta} u \in L^{2}\left(\mathbb{R}^{n}\right)$. By the boundedness of operators of order 0 on $L^{2}$, all terms on the right are in $L^{2}\left(\mathbb{R}^{n}\right)$ and we have shown the inclusion of $H_{\text {iso }}^{k}\left(\mathbb{R}^{n}\right)$ in the first space space on the right in $\left(\frac{1}{3.29} .20\right.$. . The converse is immediate, sp thisoproyes the first equality in ( 3.59 ) for $k>0$. Certainly the third space in (3.59) contains in the second. The existence of elliptic parametrix $B$ for the ellipic operator $A$ proves the converse since any isotropic pseudodifferential operator of order $A^{\prime}$ of order $k$ can be effectively factorized as

$$
A^{\prime}=A^{\prime}(B A+E)=B^{\prime} A+E^{\prime}, B^{\prime} \in \Psi_{\infty-\text { iso }}^{0}\left(\mathbb{R}^{n}\right), E^{\prime} \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)
$$

Thus, $A u \in L^{2}\left(\mathbb{R}^{n}\right)$ implies that $A^{\prime} u \in L^{2}\left(\mathbb{R}^{n}\right)$.
It also follows from second identification that $\mathcal{S}\left(\mathbb{R}^{n}\right)$ is dense in $H_{\text {loc }}^{k}\left(\mathbb{R}^{n}\right)$. Thus, if $A u \in L^{2}\left(\mathbb{R}^{n}\right)$ and we choose $f_{n} \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ with $f_{n} \rightarrow A u$ in $L^{2}\left(\mathbb{R}^{n}\right)$ then, with $B$ a parametrix for $A, u_{n}^{\prime}=B f_{n} \rightarrow B A u=u+E u$. Thus $u_{n}=u_{n}^{\prime}-E u \in \mathcal{S}\left(\mathbb{R}^{n}\right) \rightarrow u$ in $L^{2}\left(\mathbb{R}^{n}\right)$ and $A u_{n} \rightarrow u$ in $L^{2}\left(\mathbb{R}^{n}\right)$ proving the density.

The Riesz representation theorem shows that $v \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ is in the dual space, $H_{\mathrm{iso}}^{-k}\left(\mathbb{R}^{n}\right)$, if and only if there exists $v^{\prime} \in H_{\mathrm{iso}}^{k}\left(\mathbb{R}^{n}\right)$ such that

$$
\begin{array}{|l|}
\hline 1.2 .2000 .343 \\
\hline
\end{array}
$$

$$
\begin{align*}
& v(u)=\left\langle u, v^{\prime}\right\rangle_{k, \text { iso }}=\left\langle u, Q_{2 k} v^{\prime}\right\rangle_{L^{2}}, \forall u \in \mathcal{S}\left(\mathbb{R}^{n}\right) \hookrightarrow H_{\mathrm{iso}}^{k}\left(\mathbb{R}^{n}\right)  \tag{3.62}\\
& \text { with } Q_{2 k}=\sum_{|\alpha|+|\beta| \leq k} D^{\beta} x^{2 \alpha} D^{\beta}
\end{align*}
$$

This shows that $Q_{2 k}$ is an isomorphism of $H_{\mathrm{iso}}^{k}\left(\mathbb{R}^{n}\right)$ onto $H_{\mathrm{iso}}^{-k}\left(\mathbb{R}^{n}\right)$ as subspaces of $\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$. Notice that $Q_{2 k} \in \Psi_{\text {iso }}^{2 k}\left(\mathbb{R}^{n}\right)$ is elliptic, self-adjoint and_inyertible, since it is strictly positive. This now gives the same identification ( 3.59 ) for $k<0$.

The case $k=0$ follows directly from the $L^{2}$ boundedness of operators of order 0 so the proof is complete.

In view of this identification we define the isotropic Sobolev spaces or any real order the same way
1.2 .2000 .347

$$
\begin{equation*}
H_{\mathrm{iso}}^{s}\left(\mathbb{R}^{n}\right)=\left\{u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) ; A u \in L^{2}\left(\mathbb{R}^{n}\right) \forall A \in \Psi_{\mathrm{iso}}^{-s}\right\}, s \in \mathbb{R} \tag{3.63}
\end{equation*}
$$

These are Hilbertable spaces, with the Hilbert norm being given by $\|A u\|_{L^{2}\left(\mathbb{R}^{n}\right.}$ for any $A \in \Psi_{\text {iso }}^{s}\left(\mathbb{R}^{n}\right)$ which is elliptic and invertible.
1.2 .2000 .348

Proposition 3.10. Any element $A \in \Psi_{\infty-\text { iso }}^{m}\left(\mathbb{R}^{n}\right)$, $m \in \mathbb{R}$, defines a bounded linear operator

$$
\begin{equation*}
A: H_{\mathrm{iso}}^{s}\left(\mathbb{R}^{n}\right) \longrightarrow H_{\mathrm{iso}}^{s-m}\left(\mathbb{R}^{n}\right), \forall s \in \mathbb{R} \tag{3.64}
\end{equation*}
$$

This operator is Fredholm if and only if $A$ is elliptic. For any $s \in \mathbb{R}, \mathcal{S}\left(\mathbb{R}^{n}\right) \hookrightarrow$ $H_{\mathrm{iso}}^{s}\left(\mathbb{R}^{n}\right)$ is dense and $H_{\mathrm{iso}}^{-s}\left(\mathbb{R}^{n}\right)$ may be identified as the dual of $H_{\mathrm{iso}}^{s}\left(\mathbb{R}^{n}\right)$ with respect to the continuous extension of the $L^{2}$ pairing.

Proof. A straightforward application of the calculus, with the exception of the necessity of ellipticity for an isotropic pseudodifferential operator the to This is discussed in the problems beginning at Problem 3.10.

### 3.10. The residual group

By definition, $\mathcal{G}_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ is the set (if you want to be concrete you can think of them as operators on $L^{2}\left(\mathbb{R}^{n}\right)$ ) of invertible elements of the ring Id $+\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}_{1 \text { 1.2.2000.330 }}^{n}\right.$ If we identify this topologically with $\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ then, as follows from Corollary 3.1, $\mathcal{G}_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ is open. We will think of it as an infinite-dimensional manifold modeled, of course, on the linear space $\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right) \simeq \mathcal{S}\left(\mathbb{R}^{2 n}\right)$. Since I have no desire to get too deeply into the general theory of such Fréchet manifolds I will keep the discussion as elementary as possible.

The dual space of $\mathcal{S}\left(\mathbb{R}^{p}\right)$ is $\mathcal{S}^{\prime}\left(\mathbb{R}^{p}\right)$. If we want to think of $\mathcal{S}\left(\mathbb{R}^{p}\right)$ as a manifold we need to consider smooth functions and forms on it. In the finite-dimensional case, the exterior bundles are the antisymmetric parts of the tensor powers of the dual. Since we are in infinite dimensions the tensor power needs to be completed and the usual choice is the 'projective' tensor product. In our case this is something quite simple namely the $k$-fold completed tensor power of $\mathcal{S}^{\prime}\left(\mathbb{R}^{p}\right)$ is just $\mathcal{S}^{\prime}\left(\mathbb{R}^{k p}\right)$. Thus we set

$$
\begin{align*}
& \Lambda^{k} \mathcal{S}\left(\mathbb{R}^{p}\right)=\left\{u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{k p}\right) ;\right. \text { for any permutation }  \tag{3.65}\\
& \left.\qquad e, u\left(x_{e(1)}, \ldots x_{e(h)}\right)=\operatorname{sgn}(e) u\left(x_{1}, \ldots x_{k}\right)\right\}
\end{align*}
$$

In view of this it is enough for us to consider smooth functions on open sets $F \subset \mathcal{S}\left(\mathbb{R}^{p}\right)$ with values in $\mathcal{S}^{\prime}\left(\mathbb{R}^{p}\right)$ for general $p$. Thus

$$
\begin{equation*}
v: F \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{p}\right), F \subset \mathcal{S}\left(\mathbb{R}^{n}\right) \text { open } \tag{3.66}
\end{equation*}
$$

is continuously differentiable on $F$ if there exists a continuous map

$$
\begin{aligned}
& v^{\prime}: F \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n+p}\right) \text { and each } u \in F \text { has a neighbourhood } U \\
& \quad \text { such that for each } N \exists M \text { with } \\
& \qquad v v\left(u+u^{\prime}\right)-v(u)-v^{\prime}\left(u ; u^{\prime}\right)\left\|_{N} \leq C\right\| u^{\prime} \|_{M}^{2}, \forall u, u+u^{\prime} \in U .
\end{aligned}
$$

Then, as usual we define smoothness as infinite differentiability by iterating this definition. The smoothness of $v$ in this sense certainly implies that if $f: X \longrightarrow$ $\mathcal{S}\left(\mathbb{R}^{n}\right)$ is smooth then $v \circ F$ is smooth.

Thus we define the notion of a smooth form on $F \subset \mathcal{S}\left(\mathbb{R}^{n}\right)$, an open set, as a smooth map

$$
\begin{equation*}
\alpha: F \rightarrow \Lambda^{k} \mathcal{S}\left(\mathbb{R}^{p}\right) \subset \mathcal{S}^{\prime}\left(\mathbb{R}^{k p}\right) \tag{3.67}
\end{equation*}
$$

In particular we know what smooth forms are on $\mathcal{G}_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$.
The de Rham differential acts on forms as usual. If $v: F \rightarrow \mathbb{C}$ is a function then its differential at $f \in F$ is $d v: F \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)=\Lambda^{1} \mathcal{S}\left(\mathbb{R}^{n}\right)$, just the derivative. As in the finite-dimensional case $d$ extends to forms by enforcing the condition that $d v=0$ for constant forms and the identity distribution over exterior products

$$
\begin{equation*}
d(\alpha \wedge \beta)=(d \alpha) \wedge \beta+(-1)^{\operatorname{deg} \alpha} \alpha \wedge d \beta \tag{3.68}
\end{equation*}
$$

### 3.11. Representations

## Sect.radial.compactification

In $\S 1.9$ the compactification of Euclidean space to a ball, or half-sphere, is described. We make the following definition, recalling that $\rho \in \mathcal{C}^{\infty}\left(\mathbb{S}^{n,+}\right)$ is a boundary defining function.

$$
\begin{align*}
& \mathcal{L}\left(\mathbb{S}^{n,+}\right)=\bigcup_{k \in \mathbb{N}_{0}} \rho^{-k} \mathcal{C}^{\infty}\left(\mathbb{S}^{n,+}\right)  \tag{3.69}\\
& \rho^{-k} \mathcal{C}^{\infty}\left(\mathbb{S}^{n,+}\right)=\left\{u \in \mathcal{C}^{\infty}\left(\operatorname{int}\left(\mathbb{S}^{n,+}\right)\right) ; \rho^{k} u \in \mathcal{C}^{\infty}\left(\mathbb{S}^{n,+}\right)\right.
\end{align*}
$$

More generally if $m \in \mathbb{R}$ we denote by $\rho^{m} \mathcal{C}^{\infty}\left(\mathbb{S}^{n,+}\right)$ the space of functions which can be written as products $u=\rho^{m} v$, with $v \in \mathcal{C}^{\infty}\left(\mathbb{S}^{n,+}\right)$; again it can be identified with a subspace of the space of $\mathcal{C}^{\infty}$ functions on the open half-sphere.

Proposition 3.11. The compactification map ( $\left(\frac{1.104}{1.94)}\right.$ extends from $\left(\frac{11.106}{1.96)}\right.$ to give, for each $m \in \mathbb{R}$, an identification of $\rho^{-m} \mathcal{C}^{\infty}\left(\mathbb{S}^{n,+}\right)$ and $S_{\mathrm{cl}}^{m}\left(\mathbb{R}^{n}\right)$.

Thus, the fact that the $\Psi_{\text {iso,cl }}^{\mathbb{Z}}\left(\mathbb{R}^{n}\right)$ form an order-filtered $*$-algebra means that $\rho^{\mathbb{Z}} \mathcal{C}^{\infty}\left(\mathbb{S}^{2 n,+}\right)$ has a non-commutative product defined on it, with $\mathcal{C}^{\infty}\left(\mathbb{S}^{2 n,+}\right)$ a subalgebra, using the left symbol isomorphism, followed by compactification.

### 3.12. Symplectic invariance of the isotropic product

The composition law for the isotropic calculus, and in particular for it smoothing part, is derived from its identification as a subalgebra of the (weighted) spaces of pseudodifferential operator on $\mathbb{R}^{n}$. There is a much more invariant formulation of the product which puts into evidence more of the invariance properties.

Let $W$ be a real symplectic vector space. Thus, $W$ is a vector space equipped with a real, antisymmetic and non-degenerate bilinear form

$$
\begin{align*}
\omega: W \times W \longrightarrow \mathbb{R}, \omega\left(w_{1}, w_{2}\right)+\omega\left(w_{2}, w_{1}\right)=0 \forall w_{1}, w_{2} & \in W  \tag{3.70}\\
\omega\left(w_{1}, w\right) & =0 \forall w \in W \Longrightarrow w_{1}=0
\end{align*}
$$

A Lagrangian subspace of $W$ is a vector space $V \subset W$ such that $\omega$ vanishes when restricted to $V$ and such that $2 \operatorname{dim} V=\operatorname{dim} W$.

Lemma 3.8. Every symplectic vector space has a Lagrangian subspace and for any choice of Lagrangian subspace $U_{1}$ there is a second Lagrangian subspace $U_{2}$ such that $W=U_{1} \oplus U_{2}$ is a Lagrangian decomposition.

Proof. First we show that there is a Lagrangian subspace. If $\operatorname{dim} W>0$ then the antisymmetry of $\omega$ shows that any 1-dimensional vector subspace is isotropic, that is $\omega$ vanishes when restricted to it. Let $V$ be a maximal isotropic subspace, that is an isotropic subspace of maximal dimension amongst isotropic subspaces. Let $U$ be a complement to $V$ in $W$. Then

$$
\omega: V \times U \longrightarrow \mathbb{R}
$$

is a non-degenerate pairing. Indeed $u \in U$ and $\omega(v, u)=0$ for all $v \in V$ then $V+\mathbb{R}\{u\}$ is also isotropic, so $u=0$ by the assumed maximality. Similarly if $v \in V$ and $\omega(v, u)=0$ for all $u \in U$ then rechalling that $\omega$ vanishes on $V, \omega(v, w)=0$ for all $w \in W$ so $v=0$. The pairing (3.11) therefore identifies $U$ with $V^{\prime}$, the dual of $V$. In particular $\operatorname{dim} w=2 \operatorname{dim} V$.

Now, choose any Lagrangian subspace $U_{1}$. We proceed to show that there is a complementary Lagrangian subspace. Certainly there is a 1-dimensional subspace which does not meet $U_{1}$. Let $V$ be an isotropic subspace which does not meet $U_{1}$ and is of maximal dimension amongst such subspaces. Suppose that $\operatorname{dim} V<\operatorname{dim} U_{1}$. Choose $w \in W$ with $w \notin V \oplus U_{1}$. Then $V \ni v \longrightarrow \omega(w, v)$ is a linear functional
on $U_{1}$. Since $U_{1}$ can be completed to a complement, any such linear functional can be written $\omega\left(u_{1}, v\right)$ for some $u_{1} \in U_{1}$. It follows that $\omega\left(w-u_{1}, v\right)=0$ for all $v \in V$. Thus $V \oplus \mathbb{R}\left\{w-u_{1}\right\}$ a non-trivial isotropic extension of $V$, contradicting the assumed maximality. Thus $V=U_{2}$ is a complement of $U_{1}$.

Given such a Lagrangian decomposition of the symplectic vector space $W$, let $X_{1}, \ldots X_{n}$ be a basisfor the dual of $U_{1}$, and let $\Xi_{1}, \ldots, \Xi_{n}$ be the dual basis, of $U_{1}$ itself. The pairing ( 3.11 ) with $U=U_{1}$ and $V=U_{2}$ identifies $U_{2}=U_{1}^{\prime}$ so the $\Xi_{i}$ can also be regarded as a basis of the dual of $U_{2}$. Thus $X_{1} \ldots X_{n}, \Xi_{1}, \ldots, \Xi_{n}$ gives a basis of $W^{\prime}=U_{1}^{\prime} \oplus U_{2}^{\prime}$. The symplectic form can then be written

$$
\begin{equation*}
\omega\left(w_{1}, w_{2}\right)=\sum_{i=1}^{n}\left(\Xi_{i}\left(w_{1}\right) X_{i}\left(w_{2}\right)-\Xi_{i}\left(w_{2}\right) X_{i}\left(w_{1}\right)\right) \tag{3.72}
\end{equation*}
$$

This is the Darboux form of $\omega$. If the $X_{i}, \Xi_{i}$ are thought of as linear functions $x_{i}, \xi_{i}$ on, $V_{123}$ now fonsidered as a manifold then these are Darboux coordinates in which(3.72) becomes

$$
\begin{equation*}
\omega=\sum_{i=1}^{n} d \xi_{i} \wedge d x_{i} \tag{3.73}
\end{equation*}
$$

The symplectic form $\omega$ defines a volume form on $W$, namely the $n$-fold wedge product $\omega^{n}$. In Darboux coordinates this is just, up to sign, the Lebesgue form $d \xi d x$.

Proposition 3.12. On any symplectic vector space, $W$, the bilinear map on $\mathcal{S}(W)$,
$a \# b(w)=(2 \pi)^{-2 n} \int_{W^{2}} e^{i \omega\left(w_{1}, w_{2}\right)} a\left(w+w_{1}\right) b\left(w+w_{2}\right) \omega^{n}\left(w_{1}\right) \omega^{2}\left(w_{2}\right), \operatorname{dim} W=2 n$
defines an associative product isomorphic to the composition of $\Psi_{\text {iso }}^{-\infty}\left(U_{1}\right)$ for any Lagrangian decomposition $W=U_{1} \oplus U_{2}$.

Corollary 3.6. Extended by continuity in the symbol space (3.74) $\frac{1988.181}{\text { 3.7. }}$ a filtered product on $S^{\infty}(W)$ which is isomorphic to the isotropic algebra on $\mathbb{R}^{2 n}$ and is invariant under symplectic linear transformation of $W$.

Proof. Written in the form (3.74) the symplectic invariance is immediate. That is, if $F$ is a linear transformation of $W$ which preserves the symplectic form, $\omega\left(F w_{1}, F w_{2}\right)=\omega\left(w_{1}, w_{2}\right)$ then
23.3.1998.183

$$
\begin{equation*}
F^{*}(a \# b)=\left(F^{*} a\right) \#\left(F^{*} b\right) \forall a, b \in \mathcal{S}(W) \tag{3.75}
\end{equation*}
$$

The same result holds for general symbols once the continuity is established.
Let us start from the Weyl quatization of the isotropic algebra. As usual for computations we may assume that the amplitudes are of order $-\infty$. Thus, $A \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ may be written
23.3.1998.184

$$
\begin{equation*}
A u(x)=\int A(x, y) u(y)=(2 \pi)^{-n} \int e^{i(x-y) \cdot \xi} a\left(\frac{1}{2}(x+y), \xi\right) u(y) d y d \xi \tag{3.76}
\end{equation*}
$$

Both the kernel $A_{2}\left(x_{8} y_{8}\right.$ and the amplitude $a(x, \xi)$ are elements of $\mathcal{S}\left(\mathbb{R}^{2 n}\right)$. The relationship (3.76) and its inverse may be written
23.3 .1998 .185

$$
\begin{gather*}
A\left(s+\frac{t}{2}, s-\frac{t}{2}\right)=(2 \pi)^{-n} \int e^{i t \cdot \xi} a(s, \xi) d \xi \\
a(x, \xi)=\int e^{-i t \cdot \xi} A\left(x+\frac{t}{2}, x-\frac{t}{2}\right) d t \tag{3.77}
\end{gather*}
$$

If $A$ has Weyl symbol $a$ and $B$ has. Weyl symbol $b$ let $c$ be the Weyl symbol of the composite $A \circ B$. Using (3.77) and (3.66)

$$
\begin{gathered}
c(s, \zeta)=\int e^{-i t \cdot \zeta} A\left(s+\frac{t}{2}, z\right) B\left(z, s-\frac{t}{2}\right) d t \\
=(2 \pi)^{-2 n} \iiint d t d z d \xi d \eta e^{i \Phi} a\left(\frac{s}{2}+\frac{t}{4}+\frac{z}{2}, \xi\right) a\left(\frac{z}{2}+\frac{s}{2}-\frac{t}{4}, \eta\right) \\
\text { where } \Phi=-t \cdot \zeta+\left(s+\frac{t}{2}-z\right) \cdot \xi+\left(z-s+\frac{t}{2}\right) \cdot \eta .
\end{gathered}
$$

Changing variables of integration to $X=\frac{z}{2}+\frac{t}{4}-\frac{s}{2}, Y=\frac{z}{2}-\frac{t}{4}-\frac{s}{2}, \Xi=\xi-\zeta$ and $H=\eta-\zeta$ this becomes

$$
\begin{aligned}
& c(s, \zeta)=(2 \pi)^{-2 n} 4^{n} \iiint d Y d X d \Xi d H \\
& \quad e^{2 i(X \cdot H-Y \cdot \Xi)} a(X+s, \Xi+\zeta) a(Y+s, H+\zeta)
\end{aligned}
$$

23.3.1998.181

This reduces to (3.74), written out in Darboux coordinates, after the change of variable $H^{\prime}=2 H, \Xi^{\prime}=2 \Xi$ and $\zeta^{\prime}=2 \zeta$. Thus the precise isomorphism with the product in Weyl form is given by

$$
\begin{equation*}
A(x, y)=(2 \pi)^{-n} \int e^{i(x-y) \cdot \xi} a_{\omega}\left(\frac{1}{2}(x+y), 2 \xi\right) u(y) d y d \xi \tag{3.78}
\end{equation*}
$$

so that composition of kernels reduces to (23.3.1.
Discuss metaplectic group here.

### 3.13. Complex order

The identification of polyhomogeneous symbols of order zero on $\mathbb{R}^{2 n}$ with the smooth functions on the radial compactification allows us to define the isotropic operators of a given complex order $z \in \mathbb{C}$. Namely, we use the left quantization map to identify

$$
\begin{equation*}
\Psi_{\text {iso }}^{z}\left(\mathbb{R}^{n}\right)=\rho^{-z} \mathcal{C}^{\infty}\left(\mathbb{S}^{2 n, 1}\right) \subset \Psi_{\infty-\text { iso }}^{\Re z}\left(\mathbb{R}^{n}\right) \tag{3.79}
\end{equation*}
$$

Here, $\rho \in \mathcal{C}^{\infty}\left(\mathbb{S}^{2 n, 1}\right)$ is a boundary defining function. Any other boundary defining function is of the form $a \rho$ with $0<a \in \mathcal{C}^{\infty}\left(\mathbb{S}^{2 n, 1}\right)$. It follows that the definition is independent of the choice of $\rho$ since $a^{z} \in \mathcal{C}^{\infty}\left(\mathbb{S}^{2 n, 1}\right)$ for any $z \in \mathbb{Z}$.

In fact it is even more useful to consider holomorphic families. Thus if $\Omega \subset \mathbb{C}$ is an open set and $h: \Omega \longrightarrow \mathbb{C}$ is holomorphic then we may consider holomorphic families of order $h$ as elements of
1.2 .2000 .335

$$
\begin{align*}
\Psi_{\text {iso }}^{h(z)}\left(\mathbb{R}^{2 n}\right)=\{A: \Omega \longrightarrow & \Psi_{\infty-\text { iso }}^{\infty}\left(\mathbb{R}^{2 n}\right)  \tag{3.80}\\
& \left.\Omega \ni z \longmapsto \rho^{h(z)} A(z) \in \mathcal{C}^{\infty}\left(\mathbb{S}^{2 n, 1}\right) \text { is holomorphic. }\right\}
\end{align*}
$$

Note that a map from $\Omega \subset \mathbb{C}$ into $\mathcal{C}^{\infty}\left(\mathbb{S}^{2 n 1,}\right)$ is said to be holomorphic it is defines an element of $\mathcal{C}^{\infty}\left(\Omega \times \mathbb{S}^{2 n, 1}\right)$ which satisies the Cauchy-Riemann equation in the first variable.
1.2 .2000 .336
S.Traces.residual
eq:1

$$
\begin{equation*}
\operatorname{Tr}: M(N ; \mathbb{C}) \rightarrow \mathbb{C}, \quad \operatorname{Tr}(A)=\sum_{i=1}^{N} A_{i i} \tag{3.81}
\end{equation*}
$$

as the sum of the diagonal entries. The fundamental property of this functional is that

$$
\begin{equation*}
\operatorname{Tr}([A, B])=0 \forall A, B \in M(N ; \mathbb{C}) \tag{3.82}
\end{equation*}
$$

To check this it is only necessary to write down the definition of the composition in the algebra. Thus

$$
(A B)_{i j}=\sum_{k=1}^{N} A_{i k} B_{k j}
$$

It follows that

$$
\begin{aligned}
\operatorname{Tr}(A B) & =\sum_{i=1}^{N}(A B)_{i i}=\sum_{i, k=1}^{N} A_{i k} B_{k i} \\
& =\sum_{k=1}^{N} \sum_{i=1}^{N} B_{k i} A_{i k}=\sum_{k=1}^{N}(B A)_{k k}=\operatorname{Tr}(B A)
\end{aligned}
$$


Of course any multiple of $\operatorname{Tr}$ has the same property ( $\mathrm{eq} \mathrm{Ba}^{2} 2$ ) but the normalization condition

$$
\begin{equation*}
\operatorname{Tr}(\mathrm{Id})=N \tag{3.83}
\end{equation*}
$$

 $\operatorname{Tr} \in M(N ; \mathbb{C})^{\prime}$ as a point in the $N^{2}$ dimensional linear space which is the dual of $M(N ; \mathbb{C})$.

Proposition 3.13. If $h$ and $g$ are holomorphic functions on an open set $\Omega \subset \mathbb{C}$ and $A(z), B(z)$ are holomorphic familes of isotropic operators of orders $h(z)$ and $g(z)$ then the composite family $A(z) \circ B(z)$ is holomorphic of order $h(z)+g(z)$.

Proof. It suffices to consider an arbitrary open subset $\Omega^{\prime} \subset \Omega$ with compact closure inside $\Omega$. Then $h$ and $g$ have bounded real parts, so $A(z), B(z) \in$ $\Psi_{\infty-\text { iso }}^{M}\left(\mathbb{R}^{2 n}\right)$ for $z \in \Omega^{\prime}$ for some fixed $M$. It follows that the composite $A(z) \circ B(z) \in$ $\Psi_{\infty-\text { iso }}^{2 M}\left(\mathbb{R}^{2 n}\right)$. The symbol is given by the usual formula. Furthermore

The algebras we are studying are topological algebras, so it makes sense to consider continuous linear functionals on them. The most important of these is the trace. To remind you what it is we consider first its properties for matrix algebras.

Let $M(N ; \mathbb{C})$ denote the algebra of $N \times N$ complex matrices. We can simply define

Lemma 3.9. If $F: M(N ; \mathbb{C}) \rightarrow \mathbb{C}$ is a linear functional satisfying ( $\mathrm{B} \cdot \mathrm{eq}: 2 \mathrm{2}$ ) and $B \in M(N ; \mathbb{C})$ is any matrix such that $F(B) \neq 0$ then $F(A)=\frac{F(B)}{\operatorname{Tr}(B)} \operatorname{Tr}(A)$.

Proof. Consider the basis of $M(N ; \mathbb{C})$ given by the elementary matrices $E_{j k}$, where $E_{j k}$ has $j k$-th entry 1 and all others zero. Thus

$$
E_{j k} E_{p q}=\delta_{k p} E_{j q}
$$

If $j \neq k$ it follows that

$$
E_{j j} E_{j k}=E_{j k}, E_{j k} E_{j j}=0
$$

Thus

$$
F\left(\left[E_{j j}, E_{j k}\right]\right)=F\left(E_{j k}\right)=0 \text { if } j \neq k
$$

On the other hand, for any $i$ and $j$

$$
E_{j i} E_{i j}=E_{j j}, E_{i j} E_{j i}=E_{i i}
$$

so

$$
F\left(E_{j j}\right)=F\left(E_{11}\right) \forall j .
$$

Since the $E_{j k}$ are a basis,

$$
\begin{aligned}
F(A) & =F\left(\sum_{j, k=1}^{N} A_{i j} E_{i j}\right) \\
& =\sum_{j, l=1}^{N} A_{j j} F\left(E_{i j}\right) \\
& =F\left(E_{11}\right) \sum_{j=1}^{N} A_{j j}=F\left(E_{11}\right) \operatorname{Tr}(A)
\end{aligned}
$$

This proves the lemma.
For the isotropic smoothing algebra we have a similar result.
isotropic trace
Proposition 3.14. If $F: \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right) \simeq \mathcal{S}\left(\mathbb{R}^{2 n}\right) \longrightarrow \mathbb{C}$ is a continuous linear functional satisfying

$$
\begin{equation*}
F([A, B])=0 \forall A, B \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right) \tag{3.84}
\end{equation*}
$$

then $F([A, B])=0$ for all $A \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ and $B \in \Psi_{\text {iso }}^{\infty}\left(\mathbb{R}^{n}\right)$ and $F$ is a constant multiple of the functional

$$
\begin{equation*}
\operatorname{Tr}(A)=\int_{\mathbb{R}^{n}} A(x, x) d x \tag{3.85}
\end{equation*}
$$

Proof. Recall that $\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right) \subset \Psi_{\text {iso }}^{\infty}\left(\mathbb{R}^{n}\right)$ is an ideal so $A \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ and $B \in \Psi_{\text {iso }}^{\infty}\left(\mathbb{R}^{n}\right)$ implies that $A B, B A \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ and it follows that the equality $F(A B)=F(B A)$, or $F([A, B])=0$, is meaningful. To see that it holds we just use the continuity of $F$. We know that if $B \in \Psi_{\text {iso }}^{\infty}\left(\mathbb{R}^{n}\right)$ then there is a sequence $B_{n} \rightarrow B$ in the topology of $\Psi_{\text {iso }}^{m}\left(\mathbb{R}^{n}\right)$ for some $m$. Since this implies $A B_{n} \rightarrow A B$, $B_{n} A \rightarrow B A$ in $\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ we see that

$$
F([A, B])=\lim _{n \rightarrow \infty} F\left(\left[A, B_{n}\right]\right)=0
$$

We use this identity to prove ( $\frac{\text { 学: } 5.85}{}$ ). Take $B=x_{j}$ or $D_{j}, j=1, \ldots, n$. Thus for any $A \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$

$$
F\left(\left[A, x_{j}\right]\right)=F\left(\left[A, D_{j}\right]\right)=0
$$

Now consider $F$ as a distribution acting on the kernel $A \in \mathcal{S}\left(\mathbb{R}^{2 n}\right)$. Since the kernel of $\left[A, x_{j}\right]$ is $A(x, y)\left(y_{j}-x_{j}\right)$ and the kernel of $\left(A, D_{j}\right)$ is $-\left(D_{y_{j}}+D_{x_{j}}\right) A(x, y)$ we conclude that, as an element of $\mathcal{S}^{\prime}\left(\mathbb{R}^{2 n}\right), F$ satisfies

$$
\left(x_{j}-y_{j}\right) F(x, y)=0,\left(D_{x_{j}}+D_{y_{j}}\right) F(x, y)=0
$$

If we make the linear change of variables to $p_{i}=\frac{x_{i}+y_{i}}{2}, q_{i}=x_{i}-y_{i}$ and set $\tilde{F}(p, q)=F(x, y)$ these conditions become

$$
D_{q_{i}} \tilde{F}=0, p_{i} \tilde{F}=0, i=1, \ldots, N .
$$

As we know from Lemmas $\frac{2.3}{1.2}$ and $\frac{2.5}{1.3}$, this implies that $\tilde{F}=c \delta(p)$ so

$$
F(x, y)=c \delta(x-y)
$$

as a distribution. Clearly $\delta(x-y)$ gives the functional $\operatorname{Tr}$ defined by (3.85), so the proposition is proved.

We still need to justify the use of the same notation, Tr , for these two functionals. However, if $L \subset \mathcal{S}\left(\mathbb{R}^{n}\right)$ is any finite dimensional subspace we may choose an orthonal basis $\varphi_{i} \in L, i=1, \ldots, l$,

$$
\int_{\mathbb{R}^{n}}\left|\varphi_{i}(x)\right|^{2} d x=0, \int_{\mathbb{R}^{n}} \varphi_{i}(x) \overline{\varphi_{j}}(x) d x=0, i \neq j
$$

Then if $a_{i j}$ is an $l \times l$ matrix,

$$
A=\sum_{i, j=1}^{\ell} a_{i j} \varphi_{i}(x) \overline{\varphi_{j}(y)} \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)
$$

From (3.8.55) we see that

$$
\begin{aligned}
\operatorname{Tr}(A) & =\sum_{i j} a_{i j} \operatorname{Tr}\left(\varphi_{i} \bar{\varphi}_{j}\right) \\
& =\sum_{i j} a_{i j} \int_{\mathbb{R}^{n}} \varphi_{i}(x) \overline{\varphi_{j}}(x) d x \\
& =\sum_{i=1}^{n} a_{i i}=\operatorname{Tr}(a)
\end{aligned}
$$

Thus the two notions of trace coincide. In any case this already follows, up to a constant, from the uniqueness in Lemma 3.9.

### 3.15. Fredholm determinant

For $N \times N$ matrices, the linear space of which we denote $M(N ; \mathbb{C})$, the determinant is a multiplicative polynomial map

$$
\begin{equation*}
\operatorname{det}: M(N ; \mathbb{C}) \longrightarrow \mathbb{C}, \operatorname{det}(A B)=\operatorname{det}(A) \operatorname{det}(B), \operatorname{det}(\mathrm{Id})=1 \tag{3.86}
\end{equation*}
$$

It is not quite determined by these conditions, since $\operatorname{det}(A)^{k}$ also satisfies then. The fundamental property of the determinant is that it defines the group of invertible elements

$$
\begin{equation*}
\mathrm{GL}(N, \mathbb{C})=\{A \in M(N ; \mathbb{C}) ; \operatorname{det}(A) \neq 0\} \tag{3.87}
\end{equation*}
$$

A reminder of a direct definition is given in Problem $\frac{112}{3.2}$.
The Fredholm determinant is an extension of this definition to a function on the ring $\operatorname{Id}+\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$, or even further to $\operatorname{Id}+\Psi_{\infty-\text { iso }}^{-2 n-\epsilon}\left(\mathbb{R}^{n}\right)$ for $\epsilon>0$. This can
be done if seyeral ways using the density of finite rank operators, as shown in Corollary 3.4. We proceed by generalizing the formula relating the determinant to the trace. Thus, for any smooth curve with values in $\mathrm{GL}(N ; \mathbb{C})$ for any $N$,

$$
\begin{equation*}
\frac{d}{d s} \operatorname{det}\left(A_{s}\right)=\operatorname{det}\left(A_{s}\right) \operatorname{tr}\left(A_{s}^{-1} \frac{A_{s}}{d s}\right) \tag{3.88}
\end{equation*}
$$

In particular if $\left(\frac{1.2 .2000 .404}{3.86) ~ \text { is augmented by the normalization condition }}\right.$
iml. 1

$$
\begin{equation*}
\left.\frac{d}{d s} \operatorname{det}(\operatorname{Id}+s A)\right|_{s=0}=\operatorname{tr}(A) \forall A \in M(N ; \mathbb{C}) \tag{3.89}
\end{equation*}
$$

then it is determined.
A branch of the logarithm can be introduced along any curve, smoothly in the parameter, and then ( $\overline{3} .88$ ) can be rewritten

$$
\begin{equation*}
d \log \operatorname{det}(A)=\operatorname{tr}\left(A^{-1} d A\right) \tag{3.90}
\end{equation*}
$$

Here $\operatorname{GL}(N ; \mathbb{C})$ is regarded as a subset of the linear space $M(N ; \mathbb{C})$ and $d A$ is the canonical identification, at the point $A$, of the tangent space to $M(N, \mathbb{C})$ with $M(N, \mathbb{C})$ itself. This just arises from the fact that $M(N, \mathbb{C})$ is a linear space. 2000.409 Thus $d A\left(\left.\frac{d}{d s}(A+s B)\right|_{s=0}=B\right.$. This allows the expression on the right in ( $\left.\bar{B} \cdot \mathbf{9 0}\right)$ to be interpreted as a smooth 1-form on the manifold $\operatorname{GL}(N ; \mathbb{C})$. Note that it is independent of the local choice of logarithm.

To define the Fredholm determinant we shall extend the 1-form

$$
\begin{equation*}
\alpha=\operatorname{Tr}\left(A^{-1} d A\right) \tag{3.91}
\end{equation*}
$$

to the group $G_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right) \hookrightarrow \operatorname{Id}+\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$. Here $d A$ has essentially the same meaning as before, given that Id is fixed. Thus at any point $A=\operatorname{Id}+B \in \operatorname{Id}+\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ it is the identification of the tangent space with $\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ using the linear structure:

$$
d A\left(\left.\frac{d}{d s}(\operatorname{Id}+B+s E)\right|_{s=0}\right)=E, E \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)
$$

Since $d A$ takes values in $\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$, the trace functional in ( $\frac{1 \mathrm{iml} .2}{}{ }^{2}$ ) is well defined. 2000.409
The 1 -form $\alpha$ is closed. In the finite-dimensional case this follows from ( $\overline{3} .920)$. For (3.91) we can compute directly. Since $d(d A)=0$, essentially by definition, and

$$
\begin{equation*}
d A^{-1}=-A^{-1} d A A^{-1} \tag{3.92}
\end{equation*}
$$

we see that
iml. 5

$$
\begin{equation*}
d \alpha=-\operatorname{Tr}\left(A^{-1}(d A) A^{-1}(d A)\right)=0 \tag{3.93}
\end{equation*}
$$

Here we have used the trace identity, and the antisymmetry of the implicit wedge product in (3.93), to conlcude that $d \alpha=0$. For a more detailed discussion of this point see Problem 3.8.

From the fact that $d \alpha=0$ we can be confident that there is, locally near any point of $G_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$, a function $f$ such that $d f=\alpha$; then we will define the Fredholm determinant by $\operatorname{det}_{\mathrm{Fr}}(A)=\exp (f)$. To define $\operatorname{det}_{\text {Fr }}$ globally we need to see that this is well defined.
iml. 9
iml. 10

$$
\begin{equation*}
\int_{\gamma} \alpha=\int_{\mathbb{S}^{1}} \gamma^{*} \alpha \in 2 \pi i \mathbb{Z} \tag{3.94}
\end{equation*}
$$

That is, $\alpha$ defines an integral cohomology class, $\left[\frac{\alpha}{2 \pi i}\right] \in H^{1}\left(G_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right) ; \mathbb{Z}\right)$.

Proof. This is where we use the approximability by finite rank operators. If $\pi_{N}$ is the orthogonal projection onto the span of the eigenspaces of the smallIf $N$.rarmonic.oscillator est $N$ eigenvalues of the harmonic oscillator then we know from Section 3.7 that $\pi_{N} E \pi_{N} \rightarrow E$ in $\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ for any element. In fact it follows that for the smooth curve that $\gamma(s)=\operatorname{Id}+E(s)$ and $E_{N}(s)=\pi_{N} E(s) \pi_{N}$ converges uniformly with all $s$ derivatives. Thus, for some $N_{0}$ and all $N>N_{0}, \mathrm{Id}+E_{N}(s)$ is a smooth curve in $G_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ and hence $\gamma_{N}(s)=\operatorname{Id}_{N}+E_{N}(s)$ is a smooth curve in $\operatorname{GL}(N ; \mathbb{C})$. Clearly

$$
\begin{equation*}
\int_{\gamma_{N}} \alpha \longrightarrow \int_{\gamma} \alpha \text { as } N \rightarrow \infty \tag{3.95}
\end{equation*}
$$

and for finite $N$ it follows from the identity of the trace with the matrix trace (see Section $\frac{1.10}{1.10}$. gives (3.94).

Now, once we have $\left(\frac{(\mathrm{iml} .10}{3.94)}\right.$ and the connectedness of $G_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ we may define

$$
\begin{equation*}
\operatorname{det}_{\mathrm{Fr}}(A)=\exp \left(\int_{\gamma} \alpha\right), \gamma:[0,1] \longrightarrow G_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right), \gamma(0)=\mathrm{Id}, \gamma(1)=A \tag{3.96}
\end{equation*}
$$ Indeed, Lemma $\frac{\operatorname{iml} .9}{3.10}$ shows that this is independent of the path chosen from the

identity to $A$. Notice that the connectedness of $G_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ follows from the connectedness of the $\mathrm{GL}(N, \mathbb{C})$ and the density argument above.

The same arguments and results apply to $G_{\infty-\text { iso }}^{-2 n-\epsilon}\left(\mathbb{R}^{n}\right)$ using the fact that the trace functional extends continuously to $\Psi_{\infty-\text { iso }}^{-2 n-\epsilon}\left(\mathbb{R}^{n}\right)$ for any $\epsilon>0$.

Proposition 3.15. The Fredholm determinant, defined by ( $\frac{1 \mathrm{iml} .12}{3.96)}$ on $G_{\mathrm{iso}}^{-\infty}\left(\mathbb{R}^{n}\right)$ (or $G_{\mathrm{iso}}^{-2 n-\epsilon}\left(\mathbb{R}^{n}\right)$ for $\epsilon>0$ ) and to be zero on the complement in $\operatorname{Id}+\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ (or $\left.\operatorname{Id}+\Psi_{\text {iso }}^{-2 n-\epsilon}\left(\mathbb{R}^{n}\right)\right)$ is an entire function satisfying

$$
\left.\begin{array}{rl}
\operatorname{det}_{\mathrm{Fr}}(A B)=\operatorname{det}_{\mathrm{Fr}}(A) \operatorname{det}_{\mathrm{Fr}}(B), & A, B \tag{3.97}
\end{array}\right) \mathrm{Id}+\Psi_{\mathrm{iso}}^{-\infty}\left(\mathbb{R}^{n}\right), ~\left(o r \operatorname{Id}+\Psi_{\text {iso }}^{-2 n-\epsilon}\left(\mathbb{R}^{n}\right)\right), \operatorname{det}_{\mathrm{Fr}}(\mathrm{Id})=1 .
$$

Proof. We start with the multiplicative property of $\operatorname{det}_{\mathrm{Fr}}$ on $G_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$. Thus is $\gamma_{1}(s)$ is a smooth curve from Id to $A_{1}$ and $\gamma_{2}(s)$ is a smooth curve from Id to $A_{2}$ then $\gamma(s)=\gamma_{1}(s) \gamma_{2}(s)$ is a smooth curve from Id to $A_{1} A_{2}$. Consider the differential on this curve. Since

$$
\frac{d\left(A_{1}(s) A_{2}(s)\right)}{d s}=\frac{d A_{1}(s)}{d s} A_{2}(s)+A_{1}(s) \frac{d A_{2}(s)}{d s}
$$

the 1-form becomes

$$
\begin{equation*}
\gamma^{*}(s) \alpha(s)=\operatorname{Tr}\left(A_{2}(s)^{-1} \frac{d A_{2}(s)}{d s}\right)+\operatorname{Tr}\left(A_{2}(s)^{-1} A_{1}(s)^{-1} \frac{d A_{2}(s)}{d s} A_{2}(s)\right) \tag{3.98}
\end{equation*}
$$

In the second term on the right we can use the trace identity, since $\operatorname{Tr}(G A)=$ $\operatorname{Tr}(A G)$ if $G \in \Psi_{\text {iso }}^{\mathbb{Z}}\left(\mathbb{R}^{n}\right)$ and $A \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$. Thus ( $\left.\overline{3.98}\right)$ becomes

$$
\gamma^{*}(s) \alpha(s)=\gamma_{1}^{*} \alpha+\gamma_{2}^{*} \alpha
$$

Inserting this into the definition of $\operatorname{det}_{\mathrm{Fr}}$ gives $\left(\frac{\mathrm{iml} 1.14}{3.97}\right.$ when both factors are in $G_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$. Of course if either factor is not invertible, then so is the product and hence both $\operatorname{det}_{\mathrm{Fr}}(A B)$ and at least one of $\operatorname{det}_{\mathrm{Fr}}(A)$ and $\operatorname{det}_{\mathrm{Fr}}(B)$ vanishes. Thus $\left(\frac{1}{3.97}\right)$ ) $h o l d s$ in general when $\operatorname{det}_{\mathrm{Fr}}$ is extended to be zero on the non-invertible elements.

Thus it remains to establish the smoothness. That $\operatorname{det}_{\mathrm{Fr}}(A)$ is smooth in any real parameters in which $A \in G_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ depends, or indeed is holomorphic in holomorphic parameters, follows from the definition since $\alpha$ clearly depends smoothly, or holomorphically, on parameters. In fact the same follows if holomorphy is examined as a function of $E, A=\operatorname{Id}+E$, for $E \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$. Thus it is only smoothness across the non-invertibles that is at issue. To prove this we use the multiplicativity just established.

If $A=\operatorname{Id}+E$ is not invertible, $F \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ then it has a generalized inverse Id $+E^{\prime}$ as in Proposition 3.6 . Since $A$ has index zero, we may actually replace $E^{\prime}$ by $E^{\prime}+E^{\prime \prime}$, where $E^{\prime \prime}$ is an invertible linear map from the orthocomplement of the range of $A$ to its null space. Then $\operatorname{Id}+E^{\prime}+E^{\prime \prime} \in G_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ and $\left(\operatorname{Id}+E^{\prime}+E^{\prime \prime}\right) A=\operatorname{Id}-\Pi_{0}$. To prove the smoothness of $\operatorname{det}_{\mathrm{Fr}}$ on a neighbourhood of $A$ it is enough to prove the smoothness on a neighbourhood of $\mathrm{Id}-\Pi_{0}$ since $\mathrm{Id}+E^{\prime}+E^{\prime \prime}$ maps a neighbourhood of the first to a neighbourhood of the second and $\operatorname{det}_{\mathrm{Fr}}$ is multiplicative. Thus consider $\operatorname{det}_{\mathrm{Fr}}$ on a set $\mathrm{Id}-\Pi_{0}+E$ where $E$ is near 0 in $\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$, in particular we may assume that $\operatorname{Id}+E \in G_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$. Thus

$$
\operatorname{det}_{\mathrm{Fr}}\left(\operatorname{Id}+E-\Pi_{0}\right)=\operatorname{det}(\operatorname{Id}+E) \operatorname{det}\left(\operatorname{Id}-\Pi_{0}+\left(G_{E}-\operatorname{Id}\right) \Pi_{0}\right)
$$

were $G_{E}=(\operatorname{Id}+E)^{-1}$ depends holomorphically on $E$. Thus it suffices to prove the smoothness of $\operatorname{det}_{\mathrm{Fr}}\left(\operatorname{Id}-\Pi_{0}+H \Pi_{0}\right)$ where $H \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$

Consider the deformation $H_{s}=\Pi_{0} H \Pi_{0}+s\left(\operatorname{Id}-\Pi_{0}\right) H \Pi_{0}, s \in[0,1]$. If Id $-\Pi_{0}+$ $H_{s}$ is invertible for one value of $s$ it is invertible for all, since its range is always the range of $\mathrm{Id}-\Pi_{0}$ plus the range of $\Pi_{0} H \Pi_{0}$. It follows that $\operatorname{det}_{\mathrm{Fr}}\left(\mathrm{Id}-\Pi_{0}+H_{s}\right)$ is smooth in $s$; in fact it is constant. If the family is not invertible this follows immediately and if it is invertible then

$$
\begin{aligned}
& \frac{d \operatorname{det}_{\mathrm{Fr}}\left(\mathrm{Id}-\Pi_{0}+H_{s}\right)}{d s} \\
& \left.\quad=\operatorname{det}_{\mathrm{Fr}}\left(\mathrm{Id}-\Pi_{0}+H_{s}\right) \operatorname{Tr}\left(\left(\mathrm{Id}-\Pi_{0}+H_{s}\right)^{-1}\left(\operatorname{Id}-P i_{0}\right) H \Pi_{0}\right)\right)=0
\end{aligned}
$$

since the argument of the trace is finite rank and off-diagonal with respect to the decomposition by $\Pi_{0}$.

Thus finally it is enough to consider the smoothness of $\operatorname{det}_{\mathrm{Fr}}\left(\mathrm{Id}-\Pi_{0}+\Pi_{0} H \Pi_{0}\right)$ as a function of $H \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$. Since this is just $\operatorname{det}\left(\Pi_{0} H \Pi_{0}\right)$, interpreted as a finite rank map on the range of $\Pi_{0}$ the result follows from the finite dimensional case.
iml.18 Lemma 3.11. If $A \in G_{\text {iso }}^{\mathbb{Z}}\left(\mathbb{R}^{n}\right)$ and $B \in G_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ then $A B A^{-1} \in G_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ and

$$
\begin{equation*}
\operatorname{det}_{\mathrm{Fr}}\left(A B A^{-1}\right)=\operatorname{det}_{\mathrm{Fr}}(B) \tag{3.99}
\end{equation*}
$$

Proof. If $A B A^{-1}$ is not invertible then neither is $B$ so both sides of ( $\frac{1 \mathrm{iml} 1.19}{19}$ ) vanish. Thus we may assume that $B=\mathrm{Id}+E$ is invertible and let $\operatorname{Id}+E(s)$ be a smooth curve in $G_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ connecting it to the identity. Consider the function

$$
\begin{equation*}
\operatorname{det}_{\mathrm{Fr}}\left(A(\operatorname{Id}+E(s)) A^{-1}\right)=\operatorname{det}_{\mathrm{Fr}}\left(\operatorname{Id}+A E(s) A^{-1}\right) \tag{3.100}
\end{equation*}
$$

This is certainly smooth and non-vanishing and its logarithm has derivative

$$
\begin{aligned}
\operatorname{Tr}\left(A(\operatorname{Id}+E(s))^{-1} A^{-1}\right. & \left.\frac{d\left(A\left(\operatorname{Id}+E(s) A^{-1}\right)\right.}{d s}\right) \\
& =\operatorname{Tr}\left(A(\operatorname{Id}+E(s))^{-1} \frac{E(s)}{d s} A^{-1}\right)=\operatorname{Tr}\left((\operatorname{Id}+E(s))^{-1} \frac{E(s)}{d s}\right)
\end{aligned}
$$

This is also the derivative of the logarithm of $\operatorname{det}_{\mathrm{Fr}}(\operatorname{Id}+E(s))$ so the result follows.

### 3.16. Fredholm alternative

Since we have shown that $\operatorname{det}_{\mathrm{Fr}}: \operatorname{Id}+\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right) \longrightarrow \mathbb{C}$ is an entire function, we see that $G_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ is the complement of a (singular) holomorphic hypersurface, namely the surface $\left\{\operatorname{Id}+E ; \operatorname{det}_{\mathrm{Fr}}(\operatorname{Id}+E)=0\right\}$. This has the following consequence, which is sometimes call the 'Fredholm alternative' and also part of 'analytic Fredholm theory'.

## iml. 21

Lemma 3.12. If $\Omega \subset \mathbb{C}$ is an open, connected set and $A: \Omega \longrightarrow \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ is a holomorphic function then either $\operatorname{Id}+A(z)$ is invertible of all but a discrete subset of $\Omega$ and $(\operatorname{Id}+A(z))$ is meromorphic on $\Omega$ with all residues of finite rank, or else it is invertible at no point of $\Omega$.

Proof. Of course the point here is that $\operatorname{det}_{\mathrm{Fr}}(\operatorname{Id}+A(z))$ is a holomorphic function on $\Omega$. Thus, either $\operatorname{det}_{\mathrm{Fr}}(A(z))=0$ is a discrete set, $D \subset \Omega$ or else $\operatorname{det}_{\mathrm{Fr}}(\operatorname{Id}+A(z)) \equiv 0$ on $\Omega$; this uses the connectedness of $\Omega$. Since this corresponds exactly to the invertibility of $\operatorname{Id}+A(z)$ the main part of the lemma is proved. It remains only to show that, in the former case, $(\operatorname{Id}+A(z))^{-1}$ is meromorphic. Thus consider a point $p \in D$. Thus the claim is that near $p$

$$
\begin{equation*}
(\operatorname{Id}+A(z))^{-1}=\operatorname{Id}+E(z)+\sum_{j=1}^{N} z^{-j} E_{j}, E_{j} \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right) \text { of finite rank } \tag{3.101}
\end{equation*}
$$

and where $E(z)$ is locally holomorphic with values in $\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$.
If $N$ is sufficiently large and $\Pi_{N}$ is the projection onto the first $N$ eigenspaces of the harmonic oscillator then $B(z)=\operatorname{Id}+E(z)-\Pi_{N} E(z) \Pi_{N}$ is invertible near $p$ with the inverse being of the form $\operatorname{Id}+F(z)$ with $F(z)$ locally holomorphic. Now

$$
\begin{aligned}
(\operatorname{Id}+F(z))(\operatorname{Id}+E(z))=\operatorname{Id} & +(\operatorname{Id}+F(z)) \Pi_{N} E(z) \Pi_{N} \\
& =\left(\operatorname{Id}-\Pi_{N}\right)+\Pi_{N} M(z) \Pi_{N}+\left(\operatorname{Id}-\Pi_{N}\right) M^{\prime}(z) \Pi_{N}
\end{aligned}
$$

It follows that this is invertible if and only if $M(z)$ is invertible as a matrix on the range of $\Pi_{N}$. Since it must be invertible near, but not at, $p$, its inverse is a meromorphic matrix $K(z)$. It follows that the inverse of the product above can be written

$$
\begin{equation*}
\operatorname{Id}-\Pi_{N}+\Pi_{N} K(z) \Pi_{N}-\left(\operatorname{Id}-\Pi_{N}\right) M^{\prime}(z) \Pi_{N} K(z) \Pi_{N} \tag{3.102}
\end{equation*}
$$

This is meromorphic and has finite rank residues, so it follows that the same is true of $A(z)^{-1}$.

This result for the smoothing operators, which really follows from the corresponding result for matrices, gives a similar result for a holomorphic family of elliptic operators.

Proposition 3.16. Suppose $\Omega \subset \mathbb{C}$ is open and connected, $h(z)$ is holomorphic on $\Omega$ and $A(z) \in \Psi_{\text {iso }}^{h(z)}\left(\mathbb{R}^{n}\right)$ is a holomorphic and elliptic family. Then if $A(z)$ is invertible at one point it is invertible for all but a discrete set and its inverse is locally of the form $B(z)+H(z)$ where $B(z) \in \Psi_{\text {iso }}^{-h(z)}\left(\mathbb{R}^{n}\right)$ is a locally holomorphic family and $H(z)$ is meromorphic with values in $\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ and has all residues of finite rank.

Proof. By the results of Section $\frac{\text { S.Complex.order }}{3.13 \text { we can choose, at least locally near any }}$ point, a holomorphic parametrix $B(z)$ for the family in $\Psi_{\text {iso }}^{-h(z)}\left(\mathbb{R}^{n}\right)$. Suppose $p \in \Omega$ is a point at which $A(z)$ is invertible. Then the parametrix differs from the inverse by a smoothing operator, so modifying $B(z)$ by a constant smoothing operator if follows that $B(z) A(z)=\operatorname{Id}+E(z)$ is a holomorphic family in $\operatorname{Id}+\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ with $E(p)=0$. It follows from the Fredholm alternative that $\operatorname{Id}+E(z)$ is invertible with holomorphic inverse near $p$. Thus, $A(z)$ is invertible for $z$ in an open set around $p$; the set at which it is invertible is therefore open. Let $Z \subset \Omega$ be the closed set at which $A(z)$ is not invertible. If $p$ is a boundary point of $Z$ then it follows, using the notation above for a parametrix of $A(z)$ near $p$, that $\mathrm{Id}+E(z)$ is invertible near $p$ and hence $p$ is an isolated point of non-invertibility. Thus all boundary points of $Z$ are isolated. A closed subset of an open connected set in Euclidean space is either discrete or has a non-isolated boundary point, so this shows that the set of non-invertibility it is discrete. The local structure of the inverse follows from this discussion.

### 3.17. Resolvent and spectrum

One direct application of analytic Fredholm theory is to the resolvent of an elliptic operator of positive order. For simplicity we assume that $A \in \Psi_{\text {iso }}^{m}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$ with $m \in \mathbb{N}$, although the case of non-integral positive order is only slightly more complicated.

Proposition 3.17. If $A \in \Psi_{\text {iso }}^{m}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$, $m \in \mathbb{N}$, and there exists one point $\lambda^{\prime} \in \mathbb{C}$ such that $A-\lambda^{\prime}$ and $A^{*}-\overline{\lambda^{\prime}}$ both have trivial null space, then

$$
\begin{equation*}
(A-\lambda)^{-1} \in \Psi_{\text {iso }}^{-m}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right) \tag{3.103}
\end{equation*}
$$

is a meromorphic family with all residues finite rank smoothing operators; the span of the ranges of the residues at any $\tilde{\lambda}$ is the linear space of generalized eigenvalues, the solutions of

$$
\begin{equation*}
(A-\tilde{\lambda})^{p} u=0 \text { for some } p \in \mathbb{N} \tag{3.104}
\end{equation*}
$$

Proof. Since $A$ is elliptic and of positive integral order, $m, A-\lambda \in \Psi_{\text {iso }}^{m}\left(\mathbb{R}^{n}\right)$ is and entire ellipticcfamily. By hypothesis, its inverse exists for some $\lambda^{\prime} \in \mathbb{C}$. Thus, by Proposition $\frac{10}{3.16}(A-\lambda)^{-1} \in \Psi_{\text {iso }}^{-m}\left(\mathbb{R}^{n}\right)$ is a meromorphic family in the complex plane, with all residues finite rank smoothing operators.

Let $\tilde{\lambda}$ be a pole of $A-\lambda$. Since we can replace $A$ by $A-\tilde{\lambda}$ we may suppose without loss of generality that $\tilde{\lambda}=0$. Thus, for some $k$ the product $\lambda^{k}(A-\lambda)^{-1}$ is holomorphic near $\lambda=0$. Differentiating the identities

$$
(A-\lambda)\left[\lambda^{k}(A-\lambda)^{-1}\right]=\lambda^{k} \operatorname{Id}=\left[\lambda^{k}(A-\lambda)^{-1}\right](A-\lambda)
$$

up to $k$ times gives the relations

$$
\begin{align*}
& A \circ R_{k-j}=R_{k-j} \circ A=R_{k-j+1}, j=0, \cdots, k-1  \tag{3.105}\\
& \quad A \circ R_{0}=R_{0} \circ A=\mathrm{Id}+R_{1}, \text { where } \\
& (A-\lambda)^{-1}=R_{k} \lambda^{-k}+R_{k-1} \lambda^{-k+1}+\cdots+R_{0}+\cdots, R_{k+1}=0
\end{align*}
$$

Thus $A^{p} \circ R_{k-p+1}=0=R_{k-p+1} \circ A^{p}$ for $0<p \leq k$, which shows that all the residues, $R_{j}, 1 \leq j \leq k$, have ranges in the generalized eigenfunctions.

Notice also from ( $\left(\frac{\mathrm{im} 1 \mathrm{I} 128}{2}\right)$ that the range of $R_{k-j+1}$ is contained in the range of $R_{k-j}$ for each $j=0, \ldots, k-1$, and conversely for the null spaces

$$
\begin{aligned}
& \operatorname{Ran}\left(R_{k}\right) \subset \operatorname{Ran}\left(R_{k-1}\right) \subset \cdots \subset \operatorname{Ran}\left(R_{1}\right) \\
& \operatorname{Nul}\left(R_{k}\right) \supset \operatorname{Nul}\left(R_{k-1}\right) \supset \cdots \supset \operatorname{Nul}\left(R_{1}\right) .
\end{aligned}
$$

Thus,

> iml.29

$$
u \in \operatorname{Ran}\left(R_{p}\right), p \geq 1 \Longleftrightarrow \exists u_{1} \in \operatorname{Ran}\left(R_{1}\right) \text { s.t. } A^{p-1} u_{1}=u
$$

### 3.18. Residue trace

isotropic trace
We have shown, in Proposition 3.14 , the existence of a unique trace functional
 and developed by Guillemin $[5],[6]$ and Wodzicki $\frac{14],[13] \text { to investigate the traces }}{}$ on the full algebra $\Psi_{\text {iso }}^{\mathbb{Z}}\left(\mathbb{R}^{n}\right)$ of polyhomogeneous operators of integral order. We will prove the existence of a trace but defer until later the proof of its uniqueness.

Observe that for $A \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ the kernel can be written

$$
A(x, y)=(2 \pi)^{-n} \int e^{i(x-y) \xi} a_{L}(x, \xi) d \xi
$$

and hence the trace, from ( $3.8 \mathrm{eq}: 5$ ), becomes
just the integral of the left-reduced symbol. In fact this is true for any amplitude (of order $-\infty$ ) representing $A$ :
Feb.17.2000.eq:2

Feb.17.2000.eq:3

$$
\begin{equation*}
\operatorname{Tr}: \Psi_{\infty, \text { iso }}^{-2 n-\epsilon}\left(\mathbb{R}^{n}\right) \rightarrow \mathbb{C}, \text { for any } \epsilon>0 \tag{3.109}
\end{equation*}
$$

To extend it further we need somehow to regularize the resultant divergent integral in (3.107) (and to pay the price in terms of properties). ©ne elegant way to do this is to use a holomorphic family as discussed in Section 3.13. Notice that we are passing from the algebra-with-bounds in (3.109) to polyhomogeneous operators.

LEMMA 3.13. If $A(z) \in \Psi_{\text {iso }}^{z}\left(\mathbb{R}^{n}\right)$ is a holomorphic family then $f(z)=\operatorname{Tr}(A(z))$,
 at most simple poles on the divisor

$$
\{-2 n,-2 n+1, \ldots,-1,0,1, \ldots\} \subset \mathbb{C}
$$

Proof. We know that $A(z) \in \Psi_{\text {iso }}^{z}\left(\mathbb{R}^{n}\right)$ is a holomorphic family if and only if its left-reduced symbol is of the form

$$
\sigma_{L}(A(z))=\left(1+|x|^{2}+|z|^{2}\right)^{z / 2} a(z ; x, \xi)
$$

where $a(z ; x, y)$ is an entire function with values in $S_{\mathrm{phg}}^{0}\left(\mathbb{R}^{n}\right)$. For $\Re z<-2 n$ the trace of $A(z)$ is

$$
f(z)=(2 \pi)^{-n} \int_{\mathbb{R}^{2 n}}\left(1+|x|^{2}+|\xi|^{2}\right)^{z / 2} a(z ; x, \xi) d x d \xi
$$

Consider the part of this integral on the ball

$$
f_{1}(z)=(2 \pi)^{-n} \int_{|x|^{2}+|\xi|^{2} \leq 1}\left(1+|x|^{2}+|\xi|^{2}\right)^{z / 2} a(z, x, y) d x d \xi
$$

This is clearly an entire function of $z$, since the integrand is entire and the domain compact.

To analyze the remaining part $f_{2}(z)=f(z)-f_{1}(z)$ let us introduce polar coordinates

$$
r=\left(|x|^{2}+|\xi|^{2}\right)^{1 / 2}, \quad \theta=\frac{(x, \xi)}{r} \in \mathbb{S}^{2 n-1}
$$

The integral, convergent in $\Re z<-2 n$, becomes

$$
f_{2}(z)=(2 \pi)^{-n} \int_{1}^{\infty} \int_{\mathbb{S}^{2 n-1}}\left(1+r^{2}\right)^{z / 2} \tilde{a}(z ; r, \theta) d \theta r^{2 n-1} d r
$$

Let us now pass to the radical compactification of $\mathbb{R}^{2 n}$ or more prosaically, introduce $t=1 / r \in[0,1]$ as variable of integration, so

$$
f_{2}(z)=(2 \pi)^{-n} \int_{0}^{1} \int_{\mathbb{S}^{2 n-1}} t^{-z}\left(1+t^{2}\right)^{z / 2} \tilde{a}\left(z ; \frac{1}{t}, \theta\right) d \theta t^{-2 n} \frac{d t}{t} .
$$

Now the definition of $S_{\mathrm{phg}}^{0}\left(\mathbb{R}^{2 n}\right)$ reduces to the statement that

Feb.17.2000.eq:4

$$
\begin{equation*}
b(z ; t, \theta)=\left(1+t^{2}\right)^{z / 2} \tilde{a}\left(z ; \frac{1}{t}, \theta\right) \in \mathcal{C}^{\infty}\left(\mathbb{C} \times[0,1] \times \mathbb{S}^{2 n-1}\right) \tag{3.110}
\end{equation*}
$$

is holomorphic in $z$.
If we replace $b$ by its Taylor series at $t=0$ to high order,

Feb.17.2000.eq:5

$$
\begin{equation*}
b(z ; t, \theta)=\sum_{j=0}^{k} \frac{t^{j}}{j!} b_{j}(z ; \theta)+t^{k+1} b_{(k)}(z ; t, \theta), \tag{3.111}
\end{equation*}
$$

where $b_{(k)}(z ; t, \theta)$ has the same regularity ( $\begin{gathered}\text { Feb. 17.2000.eq: } 4 \\ 3.110), ~ t h e n ~ \\ f_{2}(z)\end{gathered}$ is decomposed as

Feb.17.2000.eq:6

$$
\begin{equation*}
f_{2}(z)=(2 \pi)^{-n} \sum_{j=0}^{k} \int_{0}^{1} \int_{\mathbb{S}^{2 n-1}} \frac{t^{-z+j}}{j!} b_{j}(z ; \theta) t^{-2 n} \frac{d t}{t}+f_{2}^{(k)}(z) \tag{3.112}
\end{equation*}
$$

The presence of this factor $t^{k}$ in the remainder in (Feb.17.2000.eq:5 ${ }^{\text {F.1IT) shows that } f_{2}^{(k)}(z) \text { is }}$ holmorphicin $\Re z_{0}<-2 n+k$. On the other hand the individual terms in the sum in ( 3.12 ) can be computed (for $\Re z<-2 n$ ) as

$$
\begin{aligned}
(2 \pi)^{-n}\left[\frac{t^{-z+j-2 n}}{(-z+j-2 n)}\right]_{0}^{1} \int_{\mathbb{S}^{2} n-1} b_{j}(z, \theta) & \frac{d \theta}{j!} \\
& =(2 \pi)^{-n} \frac{1}{(z-j+2 n)} \int_{\mathbb{S}^{2} n-1} b_{j}(z, \theta) \frac{d \theta}{j!}
\end{aligned}
$$

Each of these terms extends to be meromorphic in the entire complex plane, with a simple pole (at most) at $z=-2 n+j$. This shows that $f(z)$ has a meromorphic continuation as claimed.

By this argument we have actually computed the residues of the analytic continuation of $\operatorname{Tr}(A(z))$ as
$\qquad$

$$
\begin{equation*}
\lim _{z \rightarrow-2 n+j}(z-j+2 n) \operatorname{Tr}(A(z))=(2 \pi)^{-n} \int_{\mathbb{S}^{2 n-1}} a_{j}(\theta) d \theta \tag{3.113}
\end{equation*}
$$

when $a_{j}(\theta) \in \mathcal{C}^{\infty}\left(S^{2 n-1}\right)$ is the function occurring in the asymptotic expansion of the left symbol of $A(z)$ :

Feb.17.2000.eq:7

$$
\begin{align*}
\sigma_{L}(A(z)) & \sim \sum_{j=0}^{\infty}\left(|x|^{2}+|\xi|^{2}\right)^{z / 2-j} \tilde{a}_{j}(z, \theta)  \tag{3.114}\\
& |x|^{2}+|\xi|^{2} \rightarrow \infty, \theta=\frac{(x, \xi)}{\left(|x|^{2}+|\xi|^{2}\right)^{1 / 2}}, a_{j}(\theta)=\tilde{a}_{j}(-2 n+j, \theta)
\end{align*}
$$

More generally, if $m \in \mathbb{Z}$ and $A(z) \in \Psi_{\text {iso }}^{m+z}\left(\mathbb{R}^{n}\right)$ is a holomorphic family then

$$
\begin{aligned}
& \operatorname{Tr}(A(z)) \text { is meromorphic with at most } \\
& \quad \text { simple poles at }-2 n-m+\mathbb{N}_{0}
\end{aligned}
$$

Indeed this just follows by considering the family $A(z-m)$.
We are especially interested in the behavior at $z=0$. Since the residue there is an integral of the term of order $-2 n$, we know that

$$
\begin{align*}
A(z) & \in \Psi_{\text {iso }}^{m+z}\left(\mathbb{R}^{n}\right) \text { holomorphic with } A(0)=0  \tag{3.115}\\
& \Longrightarrow \operatorname{Tr}(A(z)) \text { is regular at } z=0
\end{align*}
$$

This allows us to make the following definition:

$$
\begin{gathered}
\operatorname{Tr}_{\mathrm{Res}}(A)=\lim _{z \rightarrow 0} z \operatorname{Tr}(A(z)) \text { if } \\
A(z) \in \Psi_{\text {iso }}^{m+z}\left(\mathbb{R}^{n}\right) \text { is holomorphic with } A(0)=A
\end{gathered}
$$

We know that such a holomorphic family exists, since we showed in Section ${ }^{\text {S.Complex.order }}$ existence of a holomorphic family $F(z) \in \Psi_{\text {iso }}^{z}\left(\mathbb{R}^{n}\right)$ with $F(0)=\mathrm{Id} ; A(z)=A F(z)$ is therefore an example. Similarly we know that $\operatorname{Tr}_{2}$. 2 . $A 00$ is independent of the choice of holomorphic family $A(z)$ because of (3.1i5) applied to the difference, which vanishes at zero.

Lemma 3.14. The residue functional $\operatorname{Tr}_{\operatorname{Res}}(A), A \in \Psi_{\mathrm{iso}}^{\mathbb{Z}}\left(\mathbb{R}^{n}\right)$, is a trace:

Feb.24.2000.eq:4

Feb.24.2000.eq:5

$$
\begin{equation*}
\operatorname{Tr}_{\mathrm{Res}}([A, B])=0 \forall A, B \in \Psi_{\mathrm{iso}}^{\mathbb{Z}}\left(\mathbb{R}^{n}\right) \tag{3.116}
\end{equation*}
$$

which vanishes on $\Psi_{\text {iso }}^{-2 n-1}\left(\mathbb{R}^{n}\right)$ and is given explicitly by

$$
\begin{equation*}
\operatorname{Tr}_{\text {Res }}(A)=(2 \pi)^{-n} \int_{\mathcal{S}^{2 n-1}} a_{-2 n}(\theta) d \theta \tag{3.117}
\end{equation*}
$$

where $a_{-2 n}(\theta)$ is the term of order $-2 n$ in the expansion of the left (or right) symbol of $a$.

Proof Wehave already shown that $\operatorname{Tr}_{\text {Res }}(A)$ is well-defined and (Feb.24.2000.eq:5 lows from (3.113) with $a_{-2 n}(\theta)$ the term of order $-2 n$ in the left-reduced symbol of $A=A(0)$. On the other hand, the same argument applies for the right-reduced symbol.

To see (Feb.24.2000.eq:4 (3.116) just note that if $A(z)$ and $B(z)$ are holomorphic families with $A(0)=A$, and $B(0)=B$ then $C(z)=[A(z), B(z)]$ is a holomorphic family with $C(0)=[A, B]$. On the other hand, $\operatorname{Tr}(C(z))=0$ when $\Re_{2000 \%}^{0} 0_{4}$ so the analytic continuation of $\operatorname{Tr}(C(z))$ vanishes identically and (3.116) follows.

As we shall see below, $\operatorname{Tr}_{\text {Res }}$ is the unique trace (up to a multiple of course) on $\Psi_{\text {iso }}^{\mathbb{Z}}\left(\mathbb{R}^{n}\right)$.

### 3.19. Exterior derivation

Let $A(z) \in \Psi_{\text {iso }}^{z}\left(\mathbb{R}^{n}\right)$ be a holomorphic family with $A(0)=\mathrm{Id}$. Then

$$
G(z)=A(z) \cdot A(-z) \in \Psi_{\text {iso }}^{0}\left(\mathbb{R}^{n}\right)
$$

is a holomorphic family of fixed order with $G(0)=$ Id. By analytic Fredholm theory

$$
\begin{equation*}
G^{-1}(z) \in \Psi_{\mathrm{iso}}^{0}\left(\mathbb{R}^{n}\right) \text { is a meromorphic family with finite rank poles. } \tag{3.118}
\end{equation*}
$$

It follows that $A^{-1}(z)=A(-z) G^{-1}(z)$ is a meromorphic family of order $-z$ with at most finite rank poles and regular near 0 . Set

$$
1.2 .2000 .286
$$

$$
\begin{equation*}
\Psi_{\mathrm{iso}}^{m}\left(\mathbb{R}^{n}\right) \ni B \mapsto A(z) B A^{-1}(z)=B(z) \tag{3.119}
\end{equation*}
$$

Thus $B(z)$ is a meromorphic family of order $m$ with $B(0)=B$. The derivative gives a linear map.

$$
\begin{equation*}
\Psi_{\mathrm{iso}}^{m}\left(\mathbb{R}^{n}\right) \ni B \mapsto D_{A} B=\left.\frac{d}{d z} A(z) B A^{-1}(z)\right|_{z=0} \in \Psi_{\mathrm{iso}}^{m}\left(\mathbb{R}^{n}\right) \tag{3.120}
\end{equation*}
$$

PBoposidion 3.18. For any holomornhic family of order $z$, with $A(0)=\mathrm{Id}$, the $\operatorname{map}\left(\frac{1}{3.2} 20\right)$, defined through $\left(\frac{1}{3.1 .2000 .286}\right.$ ) is a derivation and for two choices of $A(z)$ the derivations differ by an inner derivation.

Proof. Since

$$
A(z) B_{1} B_{2} A^{-1}(z)=A(z) B_{1} A^{-1}(z) A(z) B_{2} A^{-1}(z)
$$

it follows that

$$
\left.\frac{d}{d z} A(z) B_{1} B_{2} A^{-1}(z)\right|_{z=0}=\left(D_{A} B_{1}\right) \circ B_{2}+B_{1} \circ\left(D_{A} B_{2}\right)
$$

If $A_{1}(z)$ and $A_{2}(z)$ are two holomorphic families of order $z$ with $A_{1}(0)=A_{2}(0)=\operatorname{Id}$ then

$$
A_{2}(z)=A_{1}(z) G(z)
$$

when $G(z) \in \Psi_{\text {iso }}^{\infty}\left(\mathbb{R}^{n}\right)$ is a meromorphic family, with finite rank poles. Thus

$$
\begin{aligned}
A_{2}(z) B A_{2}^{-1}(z) & =A_{1}(z) G(z) B G^{-1}(z) A_{1}^{-1}(z) \\
& =A_{1}(z) B A^{-1}(z)+z A_{1}(z) H(z) A_{1}^{-1}(z)
\end{aligned}
$$

Here $H(z)=\left(G(z) B G^{-1}(z)-B\right) / z$ is a holomorphic family of degree $m$ with $H(0)=G^{\prime}(0) B-B G^{\prime}(0)$. Thus

$$
\left.\frac{d}{d z} A_{2}(z) B A_{2}^{-1}(z)\right|_{z=0}=\left.\frac{d}{d z} A_{1}(z) B A^{-1}(z)\right|_{z=0}+\left[G^{\prime}(0), B\right]
$$

which shows that the two derivations differ by an inner derivation, which is to say commutation with an element of $\Psi_{\text {iso }}^{0}\left(\mathbb{R}^{n}\right)$.

Note that in fact

$$
D_{A}: \Psi_{\text {iso }}^{m}\left(\mathbb{R}^{n}\right) \rightarrow \Psi_{\text {iso }}^{m-1}\left(\mathbb{R}^{n}\right) \forall m
$$

since the symbol of $A(z) B A^{-1}(z)$ is equal to the principal symbol of $B$ for all $z$.
For the specific choice of $A(z)=H(z)$ given by

$$
\sigma_{L}(H(z))=\left(1+|x|^{2}+|\xi|^{2}\right)^{z / 2}
$$

we shall set

$$
D_{A} B=D_{H} B .
$$

Observe that $\frac{1}{2} \log \left(1+|x|^{2}+|\xi|^{2}\right) \in S_{\infty}^{\epsilon}\left(\mathbb{R}^{2 n}\right) \forall \epsilon>0$. Thus $\log \left(1+|x|^{2}+|\xi|^{2}\right)$, defined by Weyl quantization, is an element of $\Psi_{\infty-\text { iso }}^{-\epsilon}\left(\mathbb{R}^{n}\right)$ for all $\epsilon>0$. By differentiation the symbols satisfy

$$
D_{H} B=\left[\frac{1}{2} \log \left(1+|x|^{2}+|D|^{2}\right), B\right]+[G, B]
$$

where $G \in \Psi_{\text {iso }}^{-1}\left(\mathbb{R}^{n}\right)$. Thus $D_{H}$ is not itself an interior derivation. It is therefore an exterior derivation.

### 3.20. Regularized trace

In Section $\frac{\text { S.Residue etrace }}{\mathbf{3} .18 \text { we defined }}$ the residue trace of $B$ as the residue at $z=0$ of the analytic continuation of $\operatorname{Tr}(B A(z))$, where $A(z)$ is a holomorphic family of order $z$ with $A(0)=\mathrm{Id}$. Next we consider the functional

Feb.24.2000.eq:E

$$
\begin{equation*}
\overline{\operatorname{Tr}}_{\mathrm{A}}(B)=\lim _{z=0}\left(\operatorname{Tr}(B A(z))-\frac{1}{z} \operatorname{Tr}_{\mathrm{Res}}(B)\right) \tag{3.121}
\end{equation*}
$$

In contrast to the residue trace, $\overline{\operatorname{Tr}}_{\mathrm{A}}(z)$ does depend on the choice of analytic family $A(z)$.

Lemma 3.15. If $A_{i}(z), i=1,2$, are two holomorphic families of order $z$ with $A_{i}(0)=\operatorname{Id}$ and $G^{\prime}(0)=\left.\frac{d}{d z} A_{2}(z) A_{1}^{-2}(z)\right|_{z=0}$ then

$$
\begin{equation*}
\overline{\operatorname{Tr}}_{\mathrm{A}_{2}}(B)-\overline{\operatorname{Tr}}_{\mathrm{A}_{1}}(B)=\operatorname{Tr}_{\mathrm{Res}}\left(B G^{\prime}(0)\right) \tag{3.122}
\end{equation*}
$$

Proof. Writing $G(z)=A_{2}(z) A_{1}^{-1}(z)$, which is a meromorphic family of order 0 with $G(0)=\mathrm{Id}$,

$$
\begin{aligned}
& \operatorname{Tr}\left(B A_{2}(z)\right)=\operatorname{Tr}\left(B G(z) A_{1}(z)\right) \\
& \quad=\operatorname{Tr}\left(B A_{1}(z)\right)+z \operatorname{Tr}\left(B G^{\prime}(0) A_{1}(z)\right)+z^{2} \operatorname{Tr}\left(H(z) A_{1}(z)\right)
\end{aligned}
$$

where $H(z)=\frac{B}{z^{2}}\left(G(z)-\operatorname{Id}-z G^{\prime}(0)\right)$ is then meromorphic with only finite rank poles and is regular near $z=0$. Thus the analytic continuation of $z^{2} \operatorname{Tr}(H(z) A(z))$ vanishes at zero from which ( $\frac{1 . \mathrm{eb} .24 .200 . \mathrm{eq}: \mathrm{F}}{3.122)}$ follows.

This regularized trace $\overline{\operatorname{Tr}}_{\mathrm{A}}(B)$ therefore only depends on the first order, in $z$, term in $A(z)$ at $z=0$. It is important to note that it is not itself a trace.

Lemma 3.16. If $B_{1}, B_{2} \in \Psi_{\text {iso }}^{\mathbb{Z}}\left(\mathbb{R}^{n}\right)$ then
Feb.24.2000.eq:G

$$
\begin{equation*}
\overline{\operatorname{Tr}}_{\mathrm{A}}\left(\left[B_{1}, B_{2}\right]\right)=\operatorname{Tr}_{\operatorname{Res}}\left(B_{2} D_{A} B_{1}\right) . \tag{3.123}
\end{equation*}
$$

$$
\begin{align*}
B_{1} B_{2} A(z)-B_{2} B_{1} A(z)= & B_{2}\left[A(z), B_{1}\right]+\left[B_{1}, B_{2} A(z)\right]  \tag{3.124}\\
& =B_{2}\left([A(z), B] A^{-1}(z)\right) A_{1}(z)+\left[B_{1} B_{2} A(z)\right]
\end{align*}
$$

$$
\frac{\mathrm{Feb} .24 .2000 . \mathrm{eq}: \mathrm{H}}{\mathrm{~F}}
$$

The second term on the right in $\left(\begin{array}{l}\text { Feb. } 24.2000 . e q: H\end{array}\right.$ has zero trace before analytic continuation. Thus $\overline{\operatorname{Tr}}_{A}\left(\left[B_{1}, B_{2}\right]\right)$ is the regularized value of the analytic continuation of the trace of $Q(z) A(z)$ where

$$
Q(z)=B_{2}\left[A(z), B_{1}\right] A^{-1}(z)=z D_{A} B_{1}+z^{2} L(z)
$$

with $L(z)$ meromorphic of fixed order and regular at $z=0$. Thus ( $\mathrm{Feb} .24 .2000 . \mathrm{eq}: \mathrm{G}$

Note that

$$
\begin{equation*}
\operatorname{Tr}_{\mathrm{Res}}\left(D_{A} B\right)=0 \forall B \in \Psi_{\mathrm{iso}}^{\mathbb{Z}}\left(\mathbb{R}^{n}\right) \tag{3.125}
\end{equation*}
$$

and any family $A$. Indeed the residue trace is the residue of $z=0$ of the analytic continuation of $\operatorname{Tr}(H(z) A(z))$ when $A(z)$ is any meromorphic family of fixed order with $H(0)=D_{A} B$. In particular we can take

$$
H(z)=\frac{1}{z}\left(A(z) B A^{-1}(z)-B\right) .
$$

Then $H(z) A(z)=\frac{1}{z}[A(z), B]$ so the trace vanishes before analytic continuation.

### 3.21. Projections

### 3.22. Complex powers

### 3.23. Index and invertibility

We have already seen that the elliptic elements

$$
\begin{equation*}
E_{\text {iso }}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right) \subset \Psi_{\text {iso }}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right) \hookrightarrow \mathcal{B}\left(L^{2}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)\right) \tag{3.126}
\end{equation*}
$$

define Fredholm operators. The index of such an operator

$$
1.2 .2000 .370
$$

$$
\begin{equation*}
\operatorname{Ind}(A)=\operatorname{dim} \operatorname{Nul}(A)-\operatorname{dim} \operatorname{Nul}\left(A^{*}\right) \tag{3.127}
\end{equation*}
$$

is a measure of its non-invertibility. Set

$$
\begin{equation*}
E_{\mathrm{iso}, k}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)=\left\{A \in E_{\mathrm{iso}}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right) ; \operatorname{Ind}(A)=k\right\}, k \in \mathbb{Z} \tag{3.128}
\end{equation*}
$$

Proposition 3.19. If $A \in E_{\text {iso }}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$ and $\operatorname{Ind}(A)=0$ then there exists $E \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$ such that $A+E$ is invertible in $\mathcal{B}\left(L^{2}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)\right)$ and the inverse then lies in $\Psi_{\mathrm{iso}}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$.

Proof. Let $B$ be the generalized inverse of $A$, assumed to be elliptic. The assumption that $\operatorname{Ind}(A)=0$ means that $\operatorname{Nul}(A)$ and $\operatorname{Nul}\left(A^{*}\right)$ have the same dimension. Let $e_{1}, \cdots, e_{p} \in \mathcal{S}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$ and $f_{1}, \cdots, f_{p} \in \mathcal{S}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$ be bases of $\operatorname{Nul}(A)$ and $\operatorname{Nul}\left(A^{*}\right)$. Then consider
1.2 .2000 .375

$$
\begin{equation*}
E=\sum_{j=1}^{p} f_{j}(x) \overline{e_{j}(y)} \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right) \tag{3.129}
\end{equation*}
$$

By construction $E$ is an isomorphism (in fact an arbitrary one) between $\operatorname{Nul}(A)$ and $\operatorname{Nul}\left(A^{*}\right)$. Thus $A+E$ is continuous, injective and surjective, hence has an inverse in $\mathcal{B}\left(L^{2}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)\right)$. Indeed this inverse is $B+E^{-1}$ where $E^{-1}$ is the inverse of $E$ as a
map from $\operatorname{Nul}(A)$ to $\operatorname{Nul}\left(A^{*}\right)$. This shows that $A$ can be perturbed by a smoothing operator to be invertible.

Let

$$
\begin{equation*}
G_{\mathrm{iso}}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right) \subset E_{\mathrm{iso}, 0}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right) \subset E_{\mathrm{iso}}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right) \subset \Psi 0_{\mathrm{iso}}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right) \tag{3.130}
\end{equation*}
$$

denote the group of the invertible elements (invertibility being either in $\mathcal{B}\left(L^{2}\left(\mathbb{R} ; \mathbb{C}^{N}\right)\right.$ or in $\left.\Psi_{\text {iso }}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)\right)$ in the ring of elliptic elements of index 0.
1.2 .2000 .373
1.2 .2000 .376

Corollary 3.7. The first inclusion in $\frac{(1.2 .2000 .374}{3.130) ~ i s ~ d e n s e ~ i n ~ t h e ~ t o p o l o g y ~ o f ~}$ $\Psi_{\text {iso }}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$.

Proof. This follows from the proof of Proposition $\frac{1.2 .2000 .372}{3.19}$, since $A+s E$ is invertible for all $s \neq 0$.

We next derive some simple formulæ for the index of an element of $E_{\text {iso }}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$. First observe that the trace of a finite dimensional projection is its rank, the dimension of its range. Thus
where the trace may be reinterpreted as the trace on smoothing operators. The identities, (3.34), satishied by the generalized inverse of $A$ shows that this can be rewritten

## 1.2 .2000 .377

$$
\begin{equation*}
\operatorname{Ind}(A)=-\operatorname{Tr}(B A-\mathrm{Id})+\operatorname{Tr}(A B-\mathrm{Id})=\operatorname{Tr}([A, B]) \tag{3.132}
\end{equation*}
$$

Here $[A, B]=\Pi_{\operatorname{Nul}(A)}-\Pi_{\operatorname{Nul}\left(A^{*}\right)}$ is a smoothing operator, even though both $A$ and $B$ are elliptic of order 0 .

LEMMA 3.17. If $A \in E_{\text {iso }}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$ the identity $\left(\frac{1.2 .2000 .377}{3.132), \text { which may be rewritten }}\right.$

$$
\begin{equation*}
\operatorname{Ind}(A)=\operatorname{Tr}([A, B]) \tag{3.133}
\end{equation*}
$$

holds for any parametrix $B$.
Proof. If $B^{\prime}$ is a parametrix and $B$ is the generalized inverse then $B^{\prime}-B=$ $E \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$. Thus

$$
\left[A, B^{\prime}\right]=[A, B]+[A, E]
$$

Since $\operatorname{Tr}\left([A, E]=0\right.$, one of the arguments being a smoothing operator, $\left(\frac{1.2 .2000 .379}{3.133)}\right.$ follows in general from the particular case ( $\overline{3} \cdot 132$ ).

Note that it follows from ( $\frac{1.2 .2000 .379}{3.133)} \operatorname{that} \operatorname{Ind}(A)=\operatorname{Ind}(A+E)$ if $E$ is smoothing. In fact the index is even more stable than this as we shall see, since it is locally constant on $E_{\text {iso }}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$. In any case this shows that
1.2 .2000 .381

$$
\begin{align*}
\text { Ind : } \mathcal{E}_{\text {iso }}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right) \longrightarrow & \mathbb{Z}, \operatorname{Ind}(a)=\operatorname{Ind}(A) \text { if } a=[A]  \tag{3.134}\\
\qquad \mathcal{E}_{\text {iso }}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)= & E_{\text {iso }}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right) / \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right) \\
& \subset \mathcal{A}_{\text {iso }}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)=\Psi_{\text {iso }}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right) / \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)
\end{align*}
$$

is well-defined.
The argument of the trace functional in $\frac{\left(\frac{1}{3} \cdot 2.2000 .379\right.}{3.133) \text { is a smoothing operator, but }}$ we may still rewrite the formula in terms of the regularized trace, with respect to the standard regularizer $H(z)$ with left symbol $\left(1+|x|^{2}+|\xi|^{2}\right)^{\frac{z}{2}}$. The advantage
of doing so is that we can then use the trace defect formula ( $\frac{\text { Feb. } 24.2000 . e \mathrm{eq}: \mathrm{G}}{(3.123) . \text { Thus for any }}$ elliptic isotropic operator of order 0

$$
\begin{equation*}
\operatorname{Ind}(A)=\operatorname{Tr}_{\operatorname{Res}}\left(B D_{H} A\right) \tag{3.135}
\end{equation*}
$$

Here $B$ is a parametrix for $A$. The residue trace is actually a functional

$$
\operatorname{Tr}_{\mathrm{Res}}: \mathcal{A}_{\mathrm{iso}}^{\mathbb{Z}}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right) \longrightarrow \mathbb{C}
$$

so if we write $a^{-1}$ for the inverse of $a$ in the $\operatorname{ring} \mathcal{E}_{\text {iso }}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$ then

$$
\begin{equation*}
\operatorname{Ind}(a)=\operatorname{Tr}_{\mathrm{Res}}\left(a^{-1} D_{H} a\right), D_{H}: \mathcal{A}_{\mathrm{iso}}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right) \longrightarrow \mathcal{A}_{\mathrm{iso}}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right) \tag{3.136}
\end{equation*}
$$

being the induced derivation (since $D_{H}$ clearly preserves the ideal $\Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$.
From this simple formula we can easily deduce two elementary properties of elliptic operators. These actually hold in general for Fredholm operators, although the proofs here are not valid in that generality. Namely
1.2 .2000 .383
1.2.2000.384

$$
\begin{equation*}
1.2 .2000 .385 \tag{3.139}
\end{equation*}
$$

is a multiplicative map which is the obstruction to perturbative invertibility. In the next two sections we will derive a closely related obstruction to the perturbative invertibility of a family of elliptic operators. Thus, suppose

$$
\begin{equation*}
Y \ni y \longmapsto A_{y} \in E_{\mathrm{iso}, 0}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right) \tag{3.141}
\end{equation*}
$$

is a family of elliptic operators depending smoothly on a parameter in the compact manifold $Y$. We are interested in the families perturbative invertibility question. That is, does there exist a smooth family
1.2 .2000 .388

The finuity of the 11.2 .2000 .382
The first of these follows the continuity of the formula ( $\mathbf{3 . 1 3 6}$ ) since under deformation of $a$ in $\mathcal{E}_{\text {iso }}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$ the inverse $a^{-1}$ varies continuously, so Ind is continuous and integer-valued, hence locally constant. Similarly the second, logarithmic additivity, property follows from the fact that $D_{H}$ is a derivation, so

$$
D_{H}\left(a_{1} a_{2}\right)=\left(D_{H} a_{1}\right) a_{2}+a_{1} D_{H} a_{2}
$$

and the the trace property of $\operatorname{Tr}_{\text {Res }}$ which shows that

$$
\begin{aligned}
& \operatorname{Ind}\left(a_{1} a_{2}\right)=\operatorname{Tr}_{\operatorname{Res}}\left(\left(a_{1} a_{2}\right)^{-1} D_{H}\left(a_{1} a_{2}\right)=\operatorname{Tr}\left(a_{2}^{-1} a_{1}^{-1}\left(\left(D_{H} a_{1}\right) a_{2}+a_{1} D_{H} a_{2}\right)\right.\right. \\
& =\operatorname{Tr}\left(a_{2}^{-1} a_{1}^{-1}\left(D_{H} a_{1}\right) a_{2}\right)+\operatorname{Tr}\left(a_{2}^{-1} D_{H} a_{2}\right)=\operatorname{Ind}\left(a_{1}\right)+\operatorname{Ind}\left(a_{2}\right)
\end{aligned}
$$

### 3.24. Variation 1-form

In the previous section we have seen that the index

$$
\begin{equation*}
\text { Ind : } E_{\mathrm{iso}}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right) \longrightarrow \mathbb{Z} \tag{3.140}
\end{equation*}
$$

$$
\begin{align*}
& \text { Ind }: \mathcal{E}_{\text {iso }}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right) \longrightarrow \mathbb{Z} \text { is locally constant and }  \tag{3.137}\\
& \operatorname{Ind}\left(a_{1} a_{2}\right)=\operatorname{Ind}\left(a_{1}\right)+\operatorname{Ind}\left(a_{2}\right) \forall a_{1}, a_{2} \in \mathcal{E}_{\text {iso }}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right) . \tag{3.138}
\end{align*}
$$

We have assumed that the operators have index zero since this is necessary (and sufficient) for $E_{y}$ to exist for any one $y \in Y$. Thus the issue is the smoothness (really just the continuity) of the perturbation $E_{y}$.

We shall start by essentially writing down such a putative obstruction directly and then subsequently we shall investigate its topological origins.
1.2 .2000 .390

Proposition 3.20. If a smooth family $\left(\frac{11}{3} .2 .2000 .387\right)$, parameterized by a compact manifold $Y_{0}$ is perturbatively invariant in the sense that there is a smooth family as in (3.142), then the closed 2 -form on $Y$

$$
\begin{align*}
\beta=\frac{1}{2} \operatorname{Tr}_{\mathrm{Res}}\left(a_{y}^{-1} d_{y} a_{y} \wedge a_{y}^{-1} d_{y} a_{y} \cdot a_{y}^{-1} D_{H} a_{y}\right) & \in \mathcal{C}^{\infty}\left(Y ; \Lambda^{2}\right)  \tag{3.143}\\
a_{y} & =\left[A_{y}\right] \in \mathcal{E}_{\mathrm{iso}, 0}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)
\end{align*}
$$

is exact.
Proof. Note first that $\beta$ is indeed a smooth form, since the full symbolic inverse depends smoothly on parameters. Next we show that $\beta$ is always closed. The 1-forms $a_{y}^{-1} d_{y} a_{y} a_{y}^{-1}$ and $d a_{y}$ are exact so differentiating directly gives

$$
\begin{align*}
d \beta= & \frac{1}{2} \operatorname{Tr}_{\operatorname{Res}}\left(a_{y}^{-1} d_{y} a_{y} \wedge a_{y}^{-1} d_{y} a_{y} \wedge d\left(a_{y}^{-1} D_{H} a_{y}\right)\right) \\
= & \left.-\frac{1}{2} \operatorname{Tr}_{\operatorname{Res}}\left(a_{y}^{-1} d_{y} a_{y} \wedge a_{y}^{-1} d_{y} a_{y} \wedge a_{y}^{-1} d_{y} a_{y}^{-1} D_{H} a_{y}\right)\right)  \tag{3.144}\\
& +\frac{1}{2} \operatorname{Tr}_{\operatorname{Res}}\left(a_{y}^{-1} d_{y} a_{y} \wedge a_{y}^{-1} d_{y} a_{y} \wedge a_{y}^{-1} D_{H}\left(d a_{y}\right)\right) \\
= & \frac{1}{2} \operatorname{Tr}_{\operatorname{Res}}\left(a_{y}^{-1} d_{y} a_{y} \wedge a_{y}^{-1} d_{y} a_{y} \wedge D_{H}\left(a_{y}^{-1} d a_{y}\right)\right)
\end{align*}
$$

Using the trace property and the commutativity of a 2 -form with other forms the last expression can be written

$$
\begin{equation*}
\frac{1}{6} \operatorname{Tr}_{\mathrm{Res}}\left(D_{H}\left(a_{y}^{-1} d_{y} a_{y} \wedge a_{y}^{-1} d_{y} a_{y} \wedge a_{y}^{-1} d a_{y}\right)\right)=0 \tag{3.145}
\end{equation*}
$$


Now, suppose that a smooth perturbation as in (1.2.2000.388 replace $A_{y}$ by $A_{y}+E_{y}$ without affecting $\beta$, since the residue trace vanishes on the ideal of smoothing operators. Thus we can assume that $A_{y}$ itself is invertible. Then consider the 1 -form defined using the regularized trace

$$
\begin{equation*}
\bar{\alpha}=\overline{\operatorname{Tr}}_{\mathrm{H}}\left(A_{y}^{-1} d_{y} A_{y}\right) \tag{3.146}
\end{equation*}
$$

This is an extension of the 1 -form $d \log \operatorname{det}_{F}$ on $G_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$. The extension is not in general closed, because the regularized trace does not satisfy the trace condition. Using the stanadard formula for the variation of the inverse, $d A_{y}^{-1}=-A_{y}^{-1} d A_{y} A_{y}^{-1}$, the exterior derivative is the 2 -form
1.2 .2000 .394

$$
\begin{equation*}
d \bar{\alpha}=-\overline{\operatorname{Tr}}_{\mathrm{H}}\left(A_{y}^{-1}\left(d_{y} A\right) A_{y}^{-1} d_{y} A_{y}\right) \tag{3.147}
\end{equation*}
$$

The 2-form argument is a commutator. Indeed, in terms of local coordinates we can write

$$
\begin{aligned}
& A_{y}^{-1}\left(d_{y} A\right) A_{y}^{-1} d_{y} A_{y}=\sum_{j, k=1}^{p} A_{y}^{-1}\left(\frac{\partial A}{\partial y_{j}}\right) A_{y}^{-1}\left(\frac{\partial A}{\partial y_{k}}\right) d y_{j} \wedge d y_{k} \\
&=\frac{1}{2} \sum_{j, k=1}^{p}\left(A_{y}^{-1}\left(\frac{\partial A}{\partial y_{j}}\right) A_{y}^{-1}\left(\frac{\partial A}{\partial y_{k}}\right)\right.\left.-A_{y}^{-1}\left(\frac{\partial A}{\partial y_{k}}\right) A_{y}^{-1}\left(\frac{\partial A}{\partial y_{j}}\right)\right) d y_{j} \wedge d y_{k} \\
&=\frac{1}{2} \sum_{j, k=1}^{p}\left[A_{y}^{-1}\left(\frac{\partial A}{\partial y_{j}}\right), A_{y}^{-1}\left(\frac{\partial A}{\partial y_{k}}\right)\right] d y_{j} \wedge d y_{k}
\end{aligned}
$$

Applying the trace defect formula ( $\left(\begin{array}{l}\text { Feb. 24.2000.eq:G } \\ 3.123) \\ \text { shows that }\end{array}\right.$
1.2 .2000 .395

$$
\begin{equation*}
d \bar{\alpha}=-\frac{1}{2} \operatorname{Tr}_{\operatorname{Res}}\left(A_{y}^{-1} d_{y} A_{y} \wedge D_{H}\left(A_{y}^{-1} d_{y} A_{y}\right)\right) \tag{3.148}
\end{equation*}
$$

locally and hence globally.
Expanding the action of the derivation $D_{H}$ gives

## 1.2 .2000 .396

$$
\begin{align*}
& d \bar{\alpha}=\beta-\frac{1}{2} \operatorname{Tr}_{\mathrm{Res}}\left(A_{y}^{-1} d_{y} A_{y} \wedge A_{y}^{-1} d_{y}\left(D_{H} A_{y}\right)\right)=\beta-d \gamma, \text { where }  \tag{3.149}\\
& \gamma=\frac{1}{2} \operatorname{Tr}_{\mathrm{Res}}\left(A_{y}^{-1} d_{y} A_{y} \wedge \cdot A_{y}^{-1} D_{H} A_{y}\right)
\end{align*}
$$

We conclude that if $A_{y}$ has an invertible lift then $\beta$ is exact.
Note that the form $\gamma$ in $\left(\frac{1.2 .2000 .396}{(3.149)}\right.$ is well-defined as a form on $\mathcal{E}_{\text {iso }, 0}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$, and is independent of the perturbation. Thus the cohomology class which we have constructed as the obstruction to perturbative invertibility can be written
1.2 .2000 .397

$$
\begin{equation*}
[\beta]=[\beta-d \gamma] \in H^{2}\left(\mathcal{E}_{\text {iso }, 0}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)\right) \tag{3.150}
\end{equation*}
$$

### 3.25. Determinant bundle

To better explain the topological origin of the cohomology class ( $\frac{1.2 .2000 .397}{3.150) \text { we con- }}$ struct the determinant bundle. This was originally introduced for families of Dirac operators by Quillen $\frac{10 . \text { Recall that the Fredholm determinant is a character }}{\text { IN }}$

$$
1.2 .2000 .398
$$

$$
\begin{align*}
& \operatorname{det}_{\mathrm{Fr}}: \operatorname{Id}+\Psi_{\mathrm{iso}}^{-2 n-1}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right) \longrightarrow \mathbb{C}  \tag{3.151}\\
& \quad \operatorname{det}_{\mathrm{Fr}}(A B)=\operatorname{det}_{\mathrm{Fr}}(A) \operatorname{det}_{\mathrm{Fr}}(B) \forall A, B \in \operatorname{Id}+\Psi_{\mathrm{iso}}^{-2 n-1}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)
\end{align*}
$$

As we shall see, it is not possible to extend the Fredholm determinant as a multiplicative function to $G_{\text {iso }}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$, essentially because of the non-extendibility of the trace.

However in trying to extend the determinant we can consider the possible values it would take on a point $A \in G_{\mathrm{iso}}^{0}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$ as the set of pairs $(A, z), z \in \mathbb{C}$. Thus we simple consider the product

$$
\begin{equation*}
D^{0}=G^{0} \times \mathbb{C} \tag{3.152}
\end{equation*}
$$

where from now on we simplify the notation and write $G^{0}=G_{\text {iso }}\left(\mathbb{R}^{n} ; \mathbb{C}^{N}\right)$ etc. Although it is not reasonable to expect full multiplicative of the determinant, it is more reasonable to expect the determinant of $A(\operatorname{Id}+B), B \in \Psi^{-2 n-1}$ to be related to the product of determinants. Thus it is natural to identify pairs in $D^{0}$,

$$
\begin{array}{|l|}
\hline 1.2 .2000 .400 \\
\hline
\end{array}
$$

$$
\begin{gather*}
(A, z) \sim_{p}\left(A^{\prime}, z^{\prime}\right) \text { if } \\
A, A^{\prime} \in G^{0}, A^{\prime}=A(\operatorname{Id}+B), z^{\prime}=\operatorname{det}_{\mathrm{Fr}}(\operatorname{Id}+B) z, B \in \Psi^{p}, p<-2 n \tag{3.153}
\end{gather*}
$$

The equivalence relations here are slightly different, depending on $p$. In all cases the action of the determinant is linear, so the quotient is a line bundle.

Lemma 3.18. For any integer $p<-2 n$, and also $p=-\infty$, the quotient

$$
\begin{equation*}
\mathcal{D}_{p}^{0}=D^{0} / \sim_{p} \tag{3.154}
\end{equation*}
$$

is a smooth line bundle over $\mathcal{G}_{p}^{0}=G^{0} / G^{p}$.

Proof. The projection is just the quotient in the first factor and this clearly defines a commutative square
21.2 .1998 .123
21.2 .1998 .124
21.2 .1998 .126
21.2.1998.125
21.2 .1998 .127
21.2 .1998 .128


### 3.26. Index bundle

### 3.27. Index formulæ

susceptible

### 3.28. Isotropic essential support

### 3.29. Isotropic wavefront set

### 3.30. Isotropic FBI transform

### 3.31. Problems

Problem 3.1. Define the isotropic Sobolev spaces of integral order by (3.156)
$H_{\text {iso }}^{k}\left(\mathbb{R}^{n}\right)= \begin{cases}\left\{u \in L^{2}\left(\mathbb{R}^{n}\right) ; x^{\alpha} D_{x}^{\beta} u \in L^{2}\left(\mathbb{R}^{n}\right) \forall|\alpha|+|\beta| \leq k\right\} & k \in \mathbb{N} \\ \left\{u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) ; u=\sum_{|\alpha|+|\beta| \leq-k} x^{\alpha} D_{x}^{\beta} u_{\alpha, \beta}, u_{\alpha, \beta} \in L^{2}\left(\mathbb{R}^{n}\right)\right\} & k \in-\mathbb{N} .\end{cases}$
Show that if $A \in \Psi_{\text {iso }}^{p}\left(\mathbb{R}^{n}\right)$ with $p$ an integer, then $A: H_{\text {iso }}^{k}\left(\mathbb{R}^{n}\right) \longrightarrow H_{\text {iso }}^{k-p}\left(\mathbb{R}^{n}\right)$ for any integral $k$. Deduce (using the properties of elliptic isotropic operators) that the general definition

$$
\begin{equation*}
H_{\text {iso }}^{m}\left(\mathbb{R}^{n}\right)=\left\{u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) ; A u \in L^{2}\left(\mathbb{R}^{n}\right), \forall A \in \Psi_{\text {iso }}^{-m}\left(\mathbb{R}^{n}\right)\right\}, m \in \mathbb{R} \tag{3.157}
\end{equation*}
$$ is consistent with ( ${ }^{21.2 .156)}$ ) and has the properties

$$
\begin{gather*}
A \in \Psi_{\mathrm{iso}}^{M}\left(\mathbb{R}^{n}\right) \Longrightarrow A: H_{\mathrm{iso}}^{m}\left(\mathbb{R}^{n}\right) \longrightarrow H_{\mathrm{iso}}^{m-M}\left(\mathbb{R}^{n}\right)  \tag{3.158}\\
\bigcap_{m} H_{\mathrm{iso}}^{m}\left(\mathbb{R}^{n}\right)=\mathcal{S}\left(\mathbb{R}^{n}\right), \bigcup_{m} H_{\mathrm{iso}}^{m}\left(\mathbb{R}^{n}\right)=\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)  \tag{3.159}\\
A \in \Psi_{\mathrm{iso}}^{m}\left(\mathbb{R}^{n}\right), u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right), A u \in H^{m^{\prime}}\left(\mathbb{R}^{n}\right) \Longrightarrow u \in H^{m^{\prime}-m}\left(\mathbb{R}^{n}\right), \tag{3.160}
\end{gather*}
$$

Problem 3.2. Show that if $\epsilon>0$ then

$$
H_{\text {iso }}^{\epsilon}\left(\mathbb{R}^{n}\right) \subsetneq(1+|x|)^{-\epsilon} L^{2}\left(\mathbb{R}^{n}\right) \cap H^{\epsilon}\left(\mathbb{R}^{n}\right)
$$

Deduce that $H_{\mathrm{iso}}^{\epsilon}\left(\mathbb{R}^{n}\right) \hookrightarrow L^{2}\left(\mathbb{R}^{n}\right)$ is a compact inclusion (i.e. the image of a bounded set is precompact).

Problem 3.3. Using Problem $\frac{21.2 .1998 .127}{3.2, \text { or other wise, show that each element of }}$ $\Psi_{\text {iso }}^{-\epsilon}\left(\mathbb{R}^{n}\right), \epsilon>0$, defines a compact operator on $L^{2}\left(\mathbb{R}^{n}\right)$.

$$
\begin{equation*}
C_{j}=D_{j}+i x_{j}, A_{j}=D_{j}-i x_{j}=C_{j}^{*}, \tag{3.161}
\end{equation*}
$$

and show that
21.2 .1998 .132

$$
\begin{align*}
& \text { 162) } H=\sum_{j=1}^{n} C_{j} A_{j}+n=\sum_{j=1}^{n} A_{j} C_{j}-n,  \tag{3.162}\\
& {\left[A_{j}, H\right]=2 A_{j},\left[C_{j}, H\right]=-2 C_{j},\left[C_{l}, C_{j}\right]=0,\left[A_{l}, A_{j}\right]=0,\left[A_{l}, C_{j}\right]=2 \delta_{l k} \mathrm{Id},}
\end{align*}
$$

where $[A, B]=A \circ B-B \circ A$ is the commutator bracket and $\delta_{l k}$ is the Kronecker symbol. Knowing that $(H-\lambda) u=0$, for $\lambda \in \mathbb{C}$ and $u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ implies $u \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ (why?) show that
21.2 .1998 .133
21.2 .1998 .134

$$
\begin{align*}
& E_{\lambda}=\left\{u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) ;(H-\lambda) u=0\right\} \neq\{0\} \Longleftrightarrow \lambda \in n+2 \mathbb{N}_{0} \\
& \text { and } E_{-n+2 k}=\left\{\sum_{|\alpha|=k} c_{\alpha} C^{\alpha} \exp \left(-|x|^{2} / 2\right), c_{\alpha} \in \mathbb{C}\right\}, k \in \mathbb{N}_{0} . \tag{3.164}
\end{align*}
$$

Problem 3.7. [Definition of determinant of matrices.]
Problem 3.8. [Proof that $d \alpha=0$ in ( (3ml $^{3.93)}$.] . To prove that the 1 -form is closed it suffices to show that it is closed when restricted to any 2-dimensional submanifold. Thus we may suppose that $A=A(s, t)$ depends on 2 parameters. In terms of these parameters
iml. 7

$$
\begin{equation*}
\alpha=\operatorname{Tr}\left(A(s, t)^{-1} \frac{d A(s, t)}{d s}\right) d s+\operatorname{Tr}\left(A(s, t)^{-1} \frac{d A(s, t)}{d t}\right) d t \tag{3.165}
\end{equation*}
$$

Show that the exterior derivative can be written
iml. 8

$$
\begin{equation*}
d \alpha=\operatorname{Tr}\left(\left[A(s, t)^{-1} \frac{d A(s, t)}{d t}, A(s, t)^{-1} \frac{d A(s, t)}{d s}\right]\right) d s \wedge d t \tag{3.166}
\end{equation*}
$$

and hence that it vanishes.

Problem 3.9. If $E$ and $F$ are vector spaces, show that the space of operators $\Psi_{\text {iso }}^{m}\left(\mathbb{R}^{n} ; E, F\right)$ from $\mathcal{S}^{\prime}\left(\mathbb{R}^{n} ; E\right)$ to $\mathcal{S}^{\prime}\left(\mathbb{R}^{n} ; F\right)$ is well-defined as the matrices with entries in $\Psi_{\text {iso }}^{m}\left(\mathbb{R}^{n}\right)$ for any choice of bases of $E$ and $F$.

Problem 3.10. Necessity of ellipticity for a psuedodifferential operator to be Fredholm on the isotropic Sobolev spaces.
(1) Reduce to the case of operators of order 0.
(2) Construct a sequence in $L^{2}$ such that $\left\|u_{n}\right\|=1, u_{n} \rightarrow 0$ weakly and $A u_{n} \rightarrow 0$ strongly in $L^{2}$.

Problem 3.11. [Koszul complex] Consider the form bundles over $\mathbb{R}^{n}$. That is $\Lambda^{k} \mathbb{R}^{n}$ is the vector space of dimension $\binom{n}{k}$ consisting of the totally antisymmetric $k$-linear forms on $\mathbb{R}^{n}$. If $e_{1}, e_{2}, \ldots, e_{n}$ is the standard basis for $\mathbb{R}^{n}$ then for a $k$-tuple $\alpha e^{\alpha}$ defined on basis elements by

$$
e^{\alpha}\left(e_{i_{1}}, \ldots, e_{i_{k}}\right)=\prod_{j=1}^{k} \delta_{1_{j} \alpha_{j}}
$$

extends uniquely to a $k$-linear map. Elements $d x^{\alpha} \in \Lambda^{k} \mathbb{R}^{n}$ are defined by the total antisymmetrization of the $e^{\alpha}$. Explicitly,

$$
d x^{\alpha}\left(v_{1}, \ldots, v_{k}\right)=\sum_{\pi} \operatorname{sgn} \pi e^{\alpha}\left(v_{\pi_{1}}, \ldots, v_{\pi_{n}}\right)
$$

where the sum is over permutations $\pi$ of $\{1, \ldots, n\}$ and $\operatorname{sgn} \pi$ is the parity of $\pi$. The $d x^{\alpha}$ for strictly increasing $k$-tuples $\alpha$ of elements of $\{1, \ldots, n\}$ give a basis for $\Lambda^{k} \mathbb{R}^{n}$. The wedge product is defined by $d x^{\alpha} \wedge d x^{\beta}=d x^{\alpha, \beta}$.

Now let $\mathcal{S}^{\prime}\left(\mathbb{R}^{n} ; \Lambda^{k}\right)$ be the tensor product, that is $u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n} ; \Lambda^{k}\right)$ is a finite sum

$$
\begin{equation*}
u=\sum_{\alpha} u_{\alpha} d x^{\alpha} \tag{3.167}
\end{equation*}
$$

The annihilation operators in $\left(\frac{21.2 .1998 .131}{3.161)}\right.$ define an operator, for each $k$,

$$
D: \mathcal{S}^{\prime}\left(\mathbb{R}^{n} ; \Lambda^{k}\right) \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n} ; \Lambda^{k+1}\right), \quad D u=\sum_{j=1}^{n} A_{j} u_{\alpha} d x^{j} \wedge d x^{\alpha}
$$

Show that $D^{2}=0$. Define inner products on the $\Lambda^{k} \mathbb{R}^{n}$ by declaring the basis introduced above to be orthonormal. Show that the adjoint of $D$, defined with respect to these inner products and the $L^{2}$ pairing is

$$
D^{*}: \mathcal{S}^{\prime}\left(\mathbb{R}^{n} ; \Lambda^{k}\right) \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n} ; \Lambda^{k-1}\right), D^{*} u=\sum_{j=1}^{n} C_{j} u_{\alpha} \iota_{j} d x^{\alpha}
$$

Here, $\iota_{j}$ is 'contraction with $e_{j}$ ', it is the adjoint of $d x^{j}$. Shory that $D+D^{*}$ is an elliptic element of $\Psi_{\text {iso }}^{1}\left(\mathbb{R}^{n} ; \Lambda^{*}\right)$. Maybe using Problem 3.6 show that the null space of $D+D^{*}$ on $\mathcal{S}^{\prime}\left(\mathbb{R}^{n} ; \Lambda^{*} \mathbb{R}^{n}\right)$ is 1-dimensional. Deduce that
21.2 .1998 .138

$$
\begin{align*}
\left\{u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) ; D u=0\right\} & =\mathbb{C} \exp \left(-|x|^{2} / 2\right)  \tag{3.168}\\
& \left\{u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n} ; \Lambda^{k}\right) ; D u=0\right\}=\left(\mathcal{S}^{\prime}\left(\mathbb{R}^{n} ; \Lambda^{k-1}\right), k \geq 1\right.
\end{align*}
$$

Observe that, as an operator from $\mathcal{S}^{\prime}\left(\mathbb{R}^{n} ; \Lambda^{\text {odd }}\right)$ to $\mathcal{S}^{\prime}\left(\mathbb{R}^{n} ; \Lambda^{\text {even }}\right), D+D^{*}$ is an elliptic element of $\Psi_{\text {iso }}^{1}\left(\mathbb{R}^{n} ; \Lambda^{\text {odd }}, \Lambda^{\text {even }}\right)$ and has index 1 .

Problem 3.12. [Isotropic essential support] For an element of $S^{m}\left(\mathbb{R}^{n}\right)$ define (isotropic) essential support, or operator wavefront set, of $A \in \Psi_{\text {iso }}^{m}\left(\mathbb{R}^{n}\right)$ by

$$
\begin{equation*}
\mathrm{WF}_{\text {iso }}(A)=\operatorname{cone} \operatorname{supp}\left(\sigma_{L}(A)\right) \subset \mathbb{R}^{2 n} \backslash\{0\} . \tag{3.169}
\end{equation*}
$$

Show that $\mathrm{WF}_{\text {iso }}(A)=\operatorname{cone} \operatorname{supp}\left(\sigma_{L}(A)\right)$ and check the following

$$
\begin{gather*}
\mathrm{WF}_{\text {iso }}^{\prime}(A+B) \cup \mathrm{WF}_{\text {iso }}^{\prime}(A \circ B) \subset \mathrm{WF}_{\text {iso }}^{\prime}(A) \cap \mathrm{WF}_{\text {iso }}^{\prime}(B),  \tag{3.170}\\
\mathrm{WF}_{\text {iso }}^{\prime}(A)=\emptyset \Longleftrightarrow A \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right) . \tag{3.171}
\end{gather*}
$$

Problem 3.13. [Isotropic partition of unity] Show that if $U_{i} \subset \mathbb{S}^{n-1}$ is an open cover of the unit sphere and $\tilde{U}_{i}=\left\{Z \in \mathbb{R}^{2 n} \backslash\{0\} ; \frac{Z}{|Z|} \in U_{i}\right\}$ is the corresponding conic open cover of $\mathbb{R}^{2 n} \backslash\{0\}$ then there exist (finitely many) operators $A_{i} \in$ $\Psi_{\text {iso }}^{0}\left(\mathbb{R}^{n}\right)$ with $\mathrm{WF}_{\text {iso }}^{\prime}\left(A_{i}\right) \subset \tilde{U}_{i}$, such that

$$
\begin{equation*}
\mathrm{Id}-\sum_{i} A_{i} \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right) \tag{3.172}
\end{equation*}
$$

Problem 3.14. Suppose $A \in \Psi_{\text {iso }}^{m}\left(\mathbb{R}^{n}\right)$, is elliptic and has index zero as an operator on $\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$. Show that there exists $E \in \Psi_{\text {iso }}^{-\infty}\left(\mathbb{R}^{n}\right)$ such that $A+E$ is an isomorphism of $\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$.

Problem 3.15. [Isotropic wave front set] For $u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ define

$$
\begin{equation*}
\mathrm{WF}_{\text {iso }}(u)=\bigcap\left\{\mathrm{WF}_{\text {iso }}^{\prime}(A) ; A \in \Psi_{\text {iso }}^{0}\left(\mathbb{R}^{n}\right), A u \in \mathcal{S}\left(\mathbb{R}^{n}\right)\right\} . \tag{3.173}
\end{equation*}
$$

## CHAPTER 4

## Microlocalization

### 4.1. Calculus of supports

Recall that we have already defined the support of a tempered distribution in the slightly round-about way:
(4.1) if $u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right), \operatorname{supp}(u)=\left\{x \in \mathbb{R}^{n} ; \exists \phi \in \mathcal{S}\left(\mathbb{R}^{n}\right), \phi(x) \neq 0, \phi u=0\right\}^{\text {C }}$.

Now if $A: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ is any continuous linear operator we can consider the support of the kernel:

$$
\begin{equation*}
\operatorname{supp}(A)=\operatorname{supp}\left(K_{A}\right) \subset \mathbb{R}^{n} \times \mathbb{R}^{n}=\mathbb{R}^{2 n} \tag{tabular}
\end{equation*}
$$

We write out the space as a product here to point to the fact that any subset of the product defines (is) a relation i.e. a map on subsets:

$$
\begin{equation*}
G \subset \mathbb{R}^{n} \times \mathbb{R}^{n}, \quad S \subset \mathbb{R}^{n} \Longrightarrow \tag{4.3}
\end{equation*}
$$

One can write this much more geometrically in terms of the two projection maps

$$
\begin{equation*}
7.4 \tag{4.4}
\end{equation*}
$$



Thus $\pi_{R}(x, y)=y, \pi_{L}(x, y)=x$. Then $\left(\frac{7.3}{4.3}\right)$ can be written in terms of the action of maps on sets as

$$
\begin{equation*}
G \circ S=\pi_{L}\left(\pi_{R}^{-1}(S) \cap G\right) \tag{4.5}
\end{equation*}
$$

From this it follows that if $S$ is compact and $G$ is closed, then $G \circ S$ is closed, since its intersection with any compact set is the image of a compact set under a continuous map, hence compact. Now, by the calculus of supports we mean the 'trivial' result.

Proposition 4.1. If $A: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ is a continuous linear map then

$$
\begin{equation*}
\operatorname{supp}(A \phi) \subset \operatorname{supp}(A) \circ \operatorname{supp}(\phi) \forall \phi \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right) \tag{4.6}
\end{equation*}
$$

Proof. Since we want to bound $\operatorname{supp}(A \phi)$ we can use $\left(\frac{7.1}{4.1}\right)$ directly, i.e. show that

$$
\begin{equation*}
x \notin \operatorname{supp}(A) \circ \operatorname{supp}(\phi) \Longrightarrow x \notin \operatorname{supp}(A \phi) \tag{4.7}
\end{equation*}
$$

Since we know $\operatorname{supp}(A) \circ \operatorname{supp}(\phi)$ to be closed, the assumption that $x$ is outside this set means that there exists $\psi \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)$ with

$$
\psi(x) \neq 0 \text { and } \operatorname{supp}(\psi) \cap \operatorname{supp}(A) \circ \operatorname{supp}(\phi)=\emptyset
$$

From $\left(\frac{7.3}{4.3}\right)$ or $\left(\frac{7.5}{4.5}\right)$ this means

$$
\begin{equation*}
\operatorname{supp}(A) \cap(\operatorname{supp}(\psi) \times \operatorname{supp}(\phi))=\emptyset \text { in } \mathbb{R}^{2 n} \tag{4.8}
\end{equation*}
$$

But this certainly implies that

> | 7.10 |
| :--- |

$$
\begin{gathered}
K_{A}(x, y) \psi(x) \phi(y)=0 \\
\Longrightarrow \psi A(\phi)=\int K_{A}(x, y) \psi(x) \phi(y) d y=0
\end{gathered}
$$

Thus we have proved $\left(\frac{7.7}{4.6}\right)$ and the lemma.
Diff ops.

### 4.2. Singular supports

As well as the support of a tempered distribution we can consider the singular support:

$$
\begin{equation*}
\operatorname{sing} \operatorname{supp}(u)=\left\{x \in \mathbb{R}^{n} ; \exists \phi \in \mathcal{S}\left(\mathbb{R}^{n}\right), \phi(x) \neq 0, \phi u \in \mathcal{S}\left(\mathbb{R}^{n}\right)\right\}^{\complement} \tag{4.10}
\end{equation*}
$$

Again this is a closed set since $x \notin \operatorname{sing} \operatorname{supp}(u) \Longrightarrow \exists \phi \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ with $\phi u \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ and $\phi(x) \neq 0$ so $\phi\left(x^{\prime}\right) \neq 0$ for $\left|x-x^{\prime}\right|<\epsilon$, some $\epsilon>0$ and hence $x^{\prime} \notin \operatorname{sing} \operatorname{supp}(u)$ i.e. the complement of $\operatorname{sing} \operatorname{supp}(u)$ is open.

Directly from the definition we have

$$
\begin{align*}
& \operatorname{sing} \operatorname{supp}(u) \subset \operatorname{supp}(u) \forall u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \text { and }  \tag{4.11}\\
& \quad \operatorname{sing} \operatorname{supp}(u)=\emptyset \Longleftrightarrow u \in \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right) \tag{4.12}
\end{align*}
$$

Examples

### 4.3. Pseudolocality

We would like to have a result like $\left(\frac{7.7}{4.6}\right)$ for singular support, and indeed we can get one for pseudodifferential operators. First let us work out the singular support of the kernels of pseudodifferential operators.

Proposition 4.2. If $A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ then

$$
\begin{equation*}
\operatorname{sing} \operatorname{supp}(A)=\operatorname{sing} \operatorname{supp}\left(K_{A}\right) \subset\left\{(x, y) \in \mathbb{R}^{2 n} ; x=y\right\} \tag{4.13}
\end{equation*}
$$

Proof. The kernel is defined by an oscillatory integral

$$
\begin{equation*}
I(a)=(2 \pi)^{-n} \int e^{i(x-y) \cdot \xi} a(x, y, \xi) d \xi \tag{4.14}
\end{equation*}
$$

If the order $m$ is $<-n$ we can show by integration by parts that

$$
\begin{equation*}
(x-y)^{\alpha} I(a)=I\left(\left(-D_{\xi}\right)^{\alpha} a\right) \tag{4.15}
\end{equation*}
$$

and then this must hold by continuity for all orders. If $a$ is of order $m$ and $|\alpha|>$ $m+n$ then $\left(-D_{\xi}\right)^{\alpha} a$ is of order less than $-n$, so

$$
\begin{equation*}
(x-y)^{\alpha} I(a) \in \mathcal{C}_{\infty}^{0}\left(\mathbb{R}^{n}\right),|\alpha|>m+n \tag{4.16}
\end{equation*}
$$

In fact we can also differentiate under the integral sign:
7.18

$$
\begin{equation*}
D_{x}^{\beta} D_{y}^{\gamma}(x-y)^{\alpha} I(a)=I\left(D_{x}^{\beta} D_{y}^{\gamma}\left(-D_{\xi}\right)^{\alpha} a\right) \tag{4.17}
\end{equation*}
$$

so generalizing $\left(\frac{7.17}{4.16}\right)$ to

$$
\begin{equation*}
(x-y)^{\alpha} I(a) \in \mathcal{C}_{\infty}^{k}\left(\mathbb{R}^{n}\right) \text { if }|\alpha|>m+n+k \tag{4.18}
\end{equation*}
$$

This implies that $I(A)$ is $\mathcal{C}^{\infty}$ on the complement of the diagonal, $\{x=y\}$. This proves (4.13).

An operator is said to be pseudolocal if it satisfies the condition

Proposition 4.3. Pseudodifferential operators are pseudolocal.
Proof. Suppose $u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ has compact support and $\bar{x} \notin \operatorname{sing} \operatorname{supp}(u)$. Then we can choose $\phi \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ with $\phi \equiv 1$ near $\bar{x}$ and $\phi u \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ (by definition). Thus

$$
\begin{equation*}
u=u_{1}+u_{2}, u_{1}=(1-\phi) u, \quad u_{2} \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{4.20}
\end{equation*}
$$

Since $A: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}\left(\mathbb{R}^{n}\right), A u_{2} \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ so

$$
\begin{equation*}
\operatorname{sing} \operatorname{supp}(A u)=\operatorname{sing} \operatorname{supp}\left(A u_{1}\right) \text { and } \bar{x} \notin \operatorname{supp}\left(u_{1}\right) \tag{4.21}
\end{equation*}
$$

Choose $\psi \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ with compact support, $\psi(\bar{x})=1$ and

$$
\begin{equation*}
\operatorname{supp}(\psi) \cap \operatorname{supp}(1-\phi)=\emptyset \tag{4.22}
\end{equation*}
$$

Thus

### 7.25

$$
\begin{equation*}
\psi A u_{1}=\psi A(1-\phi) u=\tilde{A} u \tag{4.23}
\end{equation*}
$$

where

$$
\begin{equation*}
K_{\tilde{A}}(x, y)=\psi(x) K_{A}(x, y)(1-\phi(y)) \tag{4.24}
\end{equation*}
$$

Combining $\left(\frac{7.24}{4.22}\right)$ and $\left(\frac{7.15}{4.13}\right)$ shows that $K_{\tilde{A}} \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$ so, by Lemma $\frac{5.42}{2.8}, \tilde{A} u \in$ $\mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right)$ and $\bar{x} \notin \operatorname{sing} \operatorname{supp}(A u)$ by $(4.13)(?)$. This proves the proposition.

### 4.4. Coordinate invariance

If $\Omega \subset \mathbb{R}^{n}$ is an open set, put

$$
\begin{align*}
\mathcal{C}_{c}^{\infty}(\Omega) & =\left\{u \in \mathcal{S}\left(\mathbb{R}^{n}\right) ; \operatorname{supp}(u) \Subset \Omega\right\} \\
\mathcal{C}_{c}^{-\infty}(\Omega) & =\left\{u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) ; \operatorname{supp}(u) \Subset \Omega\right\} \tag{4.25}
\end{align*}
$$

respectively the space of $\mathcal{C}^{\infty}$ functions of compact support in $\Omega$ and of distributions of compact support in $\Omega$. Here $K \Subset \Omega$ indicates that $K$ is a compact subset of $\Omega$. Notice that if $u \in \mathcal{C}_{c}^{-\infty}(\Omega)$ then $u$ defines a continuous linear functional

$$
\begin{equation*}
\mathcal{C}^{\infty}(\Omega) \ni \phi \longmapsto u(\phi)=u(\psi \phi) \in \mathbb{C} \tag{4.26}
\end{equation*}
$$

where if $\psi \in \mathcal{C}_{c}^{\infty}(\Omega)$ is chosen to be identically one near $\operatorname{supp}(u)$ then $\left(\frac{\sqrt{4} \cdot 27}{4.26}\right)$ is independent of $\psi$. [Think about what continuity means here!]

Now suppose

$$
\begin{equation*}
F: \Omega \longrightarrow \Omega^{\prime} \tag{4.27}
\end{equation*}
$$

is a diffeomorphism between open sets of $\mathbb{R}^{n}$. The pull-back operation is

$$
\begin{equation*}
F^{*}: \mathcal{C}_{c}^{\infty}\left(\Omega^{\prime}\right) \longleftrightarrow \mathcal{C}_{c}^{\infty}(\Omega), F^{*} \phi=\phi \circ F \tag{4.28}
\end{equation*}
$$

Lemma 4.1. If $F$ is a diffeomorphismb $\left(\frac{7.28}{4.27}\right)$, between open sets of $\mathbb{R}^{n}$ then there is an extension by continuity of $(\overline{4.28})$ to

$$
\begin{equation*}
F^{*}: \mathcal{C}_{c}^{-\infty}\left(\Omega^{\prime}\right) \longleftrightarrow \mathcal{C}_{c}^{-\infty}(\Omega) \tag{4.29}
\end{equation*}
$$

Proof. The density of $\mathcal{C}_{c}^{\infty}(\Omega)$ in $\mathcal{C}_{c}^{-\infty}(\Omega)$, in the weak topology given by the seminorms from ( 4.26$)_{p, 7}$ can be proved in the same way as the density of $\mathcal{S}\left(\mathbb{R}^{n}\right)$ in $\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ (see Problem 4.5 ). Thus, we only need to show continuity of ( 4.29 ) in this sense. Suppose $u \in \mathcal{C}_{c}^{\infty}(\Omega)$ and $\phi \in \mathcal{C}_{c}^{\infty}\left(\Omega^{\prime}\right)$ then

$$
\begin{align*}
\left(F^{*} u\right)(\phi) & =\int u(F(x)) \phi(x) d x \\
& =\int u(y) \phi(G(y))\left|J_{G}(y)\right| d y \tag{4.30}
\end{align*}
$$

where $J_{G}(y)=\left(\frac{\partial G(y)}{\partial y}\right)$ is the Jacobian of $G$, the inverse of $F$. Thus ( $\left(\frac{7.29}{4.28}\right)$ can be written

$$
\begin{equation*}
F^{*} u(\phi)=\left(\left|J_{G}\right| u\right)\left(G^{*} \phi\right) \tag{4.31}
\end{equation*}
$$

and since $G^{*}: \mathcal{C}^{\infty}(\Omega) \longrightarrow \mathcal{C}^{\infty}\left(\Omega^{\prime}\right)$ is continuous (!) we conclude that $F^{*}$ is continuous as desired.

Now suppose that

$$
A: \mathcal{S}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)
$$

has

$$
\begin{equation*}
\operatorname{supp}(A) \Subset \Omega \times \Omega \subset \mathbb{R}^{2 n} \tag{4.32}
\end{equation*}
$$

Then

$$
\begin{equation*}
A: \mathcal{C}_{c}^{\infty}(\Omega) \longrightarrow \mathcal{C}_{c}^{-\infty}(\Omega) \tag{4.33}
\end{equation*}
$$

by Proposition $\frac{17.6}{4.1}$. Applying a diffeomorphism, $F$, as in $\left(\frac{\sqrt{4} .28}{4.27}\right)$ set

$$
\begin{equation*}
A_{F}: \mathcal{C}_{c}^{\infty}\left(\Omega^{\prime}\right) \longrightarrow \mathcal{C}_{c}^{-\infty}\left(\Omega^{\prime}\right), A_{F}=G^{*} \circ A \circ F^{*} \tag{4.34}
\end{equation*}
$$

Lemma 4.2. If $A$ satisfies $\left(\frac{7.34}{4.32}\right)$ and $F$ is a diffeomorphism $\left(\frac{7.28}{4.27}\right)$ then

$$
\begin{equation*}
K_{A_{F}}(x, y)=(G \times G)^{*} K \cdot\left|J_{G}(y)\right| \text { on } \Omega^{\prime} \times \Omega^{\prime} \tag{4.35}
\end{equation*}
$$

has compact support in $\Omega^{\prime} \times \Omega^{\prime}$.
Proof. Essentially the same as that of ( $\frac{7.32}{4.30}$ ).
Proposition 4.4. Suppose $A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ has kernel satisfying $\left(\frac{7.34}{4.32}\right)$ and $F$ is a diffeomorphism as in $\left(\frac{1.28}{4.27}\right)$ then $A_{F}$, defined by $\left(\frac{1.36}{4.34}\right)$, is an element of $\Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$.

Proof. Since $A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$,

$$
\begin{equation*}
K_{A}(x, y)=(2 \pi)^{-n} \int e^{i(x-y) \cdot \xi} a(x, \xi) d \xi \tag{4.36}
\end{equation*}
$$

for some $a \in S_{\infty}^{m}\left(\mathbb{R}^{n} ; \mathbb{R}_{7}^{n}\right)_{34}$ Now choose $\psi \in \mathcal{C}_{c}^{\infty}(\Omega)$ such that $\psi(x) \psi(y)=1$ on $\operatorname{supp}\left(K_{A}\right)$, possible by ( 4.32 ). Then

$$
\begin{equation*}
K_{A}=I(\psi(x) \psi(y) a(x, \xi)) \tag{4.37}
\end{equation*}
$$

In fact suppose $\mu_{\epsilon}(x, y) \in \mathcal{C}^{\infty}\left(\mathbb{R}^{2 n}\right)$ and $\mu \equiv 1$ in $|x-y|<\epsilon$ for $\epsilon>0, \mu(x, y)=0$ in $|x-y|>2 \epsilon$. Then if

$$
\begin{equation*}
K_{A_{\epsilon}}=I\left(\mu_{\epsilon}(x, y) \psi(x) \psi(y) a(x, \xi)\right) \tag{4.38}
\end{equation*}
$$

we know that if
7.43

$$
\begin{equation*}
A_{\epsilon}^{\prime}=A-A_{\epsilon} \text { then } K_{A_{\epsilon}^{\prime}}=\left(1-\mu_{\epsilon}(x, y)\right) K_{A} \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right) \tag{4.39}
\end{equation*}
$$

by $\left(\frac{7.15}{4.13}\right)$. Certainly $A_{\epsilon}^{\prime}$ also satisfies $\left(\frac{7.34}{4.32}\right)$ and from $\left(\frac{7.38}{4.35}\right)$
$\left(A_{\epsilon}^{\prime}\right)_{F} \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$.
So we only need to consider $A_{\epsilon}$ defined by ( $\frac{7.42}{4.38}$ ). Again using ( 14.38 ) (assuming $m<-n$ )

$$
\begin{equation*}
K_{\left(A_{\epsilon}\right)_{F}}(x, y)=(2 \pi)^{-n} \int e^{i(G(x)-G(y)) \cdot \xi} b(G(x), G(y), \xi)\left|\frac{\partial G}{\partial y}\right| d \xi \tag{4.41}
\end{equation*}
$$

where $b(x, y, \xi)=\mu_{\epsilon}(x-y) \psi(x) \psi(y) a(x, \xi)$. Applying Taylor's formula,

$$
\begin{equation*}
G(x)-G(y)=(x-y) \cdot T(x, y) \tag{4.42}
\end{equation*}
$$

where $T(x, y)$ is an invertible $\mathcal{C}^{\infty}$ matrix on $K \times K \cap\{|x-y|<\epsilon\}$ for $\epsilon<\epsilon(K)$, where $\epsilon(K)>0$ depends on the compact set $K \Subset \Omega^{\prime}$. Thus we can set(??)

$$
\begin{equation*}
\eta=T^{t}(x, y) \cdot \xi \tag{4.43}
\end{equation*}
$$

and rewrite $\left(\frac{7.45}{4.4 \mathrm{I}}\right)$ as
7.47

$$
\begin{gather*}
K_{\left(A_{\epsilon}\right)_{F}}(x, y)=(2 \pi)^{-n} \int e^{i(x-y) \cdot \eta} c(x, y, \eta) d \eta \\
c(x, y, \eta)=b\left(G(x), G(y),\left(T^{t}\right)^{-1}(x, y) \eta\right)\left|\frac{\partial G}{\partial y}\right| \cdot|\operatorname{det} T(x, y)|^{-1} \tag{4.44}
\end{gather*}
$$

So it only remains to show that $c \in S_{\infty}^{m}\left(\mathbb{R}^{2 n} ; \mathbb{R}^{n}\right)$ and the proof is complete. We can drop all the $\mathcal{C}^{\infty}$ factors, given by $|\partial G / \partial y|$ etc. and proceed to show that

$$
\begin{equation*}
\left|D_{x}^{\alpha} D_{y}^{\beta} D_{\xi}^{\gamma} a(G(x), G(y), S(x, y) \xi)\right| \leq C(1+|\xi|)^{m-|\gamma|} \quad \text { on } K \times K \times \mathbb{R}^{n} \tag{4.45}
\end{equation*}
$$

where $K \subset \subset \Omega^{\prime}$ and $S$ is $\mathcal{C}^{\infty}$ with $|\operatorname{det} S| \geq \epsilon$. The estimates with $\alpha=\beta=0$ follow easily and the general case by induction:

$$
\begin{aligned}
& D_{x}^{\alpha} D_{y}^{\beta} D_{\xi}^{\gamma} a(G(x), G(y), S(x, y) \xi) \\
&=\sum_{\substack{|\mu| \leq|\alpha|+|\beta|+|\gamma| \\
\left|\alpha^{\prime}\right| \leq|\alpha|,|\beta| \leq|\beta| \\
|\nu|+|\gamma| \leq|\mu|}} M_{\alpha, \beta, \gamma, \nu}^{\alpha^{\prime}, \rho^{\prime}, \mu^{\prime}}(x, y) \xi^{\nu}\left(D^{\alpha^{\prime}} D^{\beta^{\prime}} D^{\mu} a\right)(G(x), G(y), S \xi)
\end{aligned}
$$

where the coefficients are $\mathcal{C}^{\infty}$ and the main point is that $|\nu| \leq|\mu|$.

## P.7.1

$$
\begin{equation*}
S_{\infty}^{\infty}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right) \ni a \longmapsto q_{W}(a)=(2 \pi)^{-n} \int e^{i(x-y) \cdot \xi} a\left(\frac{x+y}{2}, \xi\right) d \xi \tag{4.46}
\end{equation*}
$$

is well-defined by continuity from $S_{\infty}^{-\infty}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right)$ and induces an isomorphism

$$
\begin{equation*}
S_{\infty}^{m}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right) \underset{q_{W}}{\stackrel{\sigma_{W}}{\rightleftarrows}} \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right) \forall m \in \mathbb{R} \tag{4.47}
\end{equation*}
$$

Find an asymptotic formula relating $q_{W}(A)$ to $q_{L}(A)$ for any $A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$.

### 4.5. Problems

Problem 4.1. Show that Weyl quantization

Problem 4.2. Show that if $A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ then $A^{*}=A$ if and only if $\sigma_{W}(A)$ is real-valued.

Problem 4.3. Is it true that every $E \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$ defines a map from $\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ to $\mathcal{S}\left(\mathbb{R}^{n}\right)$ ?

Problem 4.4. Show that $\mathcal{S}\left(\mathbb{R}^{n}\right)$ is dense in $L^{2}\left(\mathbb{R}^{n}\right)$ by proving that if $\phi \in$ $\mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right)$ has compact support and is identically equal to 1 near the origin then

$$
\begin{equation*}
u_{n}(x)=(2 \pi)^{-n} \phi\left(\frac{x}{n}\right) \int e^{i x \cdot \xi} \phi(\xi / n) \hat{u}(\xi) d \xi \in \mathcal{S}\left(\mathbb{R}^{n}\right) \text { if } u \in L^{2}\left(\mathbb{R}^{n}\right) \tag{4.48}
\end{equation*}
$$

and $u_{n} \rightarrow u$ in $L^{2}\left(\mathbb{R}^{n}\right)$. Can you see any relation to pseudodifferential operators here?

$$
\begin{equation*}
H^{k}\left(\mathbb{R}^{n}\right)=\left\{u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) ; u=\sum_{|\alpha| \leq-k} D^{\alpha} u_{\alpha}, u_{\alpha} \in L^{2}\left(\mathbb{R}^{n}\right)\right\} \tag{4.49}
\end{equation*}
$$

for $-k \in \mathbb{N}$ one does have

$$
\begin{equation*}
u \in H^{k}\left(\mathbb{R}^{n}\right) \Longleftrightarrow\langle D\rangle^{k} u \in L^{2}\left(\mathbb{R}^{n}\right) \tag{4.50}
\end{equation*}
$$

as claimed in the text.
Problem 4.6. Suppose that $a(x) \in \mathcal{C}_{\infty}^{\infty}\left(\mathbb{R}^{n}\right)$ and that $a(x) \geq 0$. Show that the operator

$$
\begin{equation*}
A=\sum_{j=1}^{n} D_{x_{j}}^{2}+a(x) \tag{4.51}
\end{equation*}
$$

can have no solution which is in $L^{2}\left(\mathbb{R}^{n}\right)$.
5.49 in the weak topology.

Problem 4.8. Use formula $\left(\frac{7.47}{4.44)}\right.$ to find the principal symbol of $A_{F}$; more precisely show that if $F^{*}: T^{*} \Omega^{\prime} \longrightarrow T^{*} \omega$ is the (co)-differential of $F$ then

$$
\sigma_{m}\left(A_{F}\right)=\sigma_{m}(A) \circ F^{*}
$$

We have now studied special distributions, the Schwartz kernels of pseudodifferential operators. We shall now apply this knowledge to the study of general distributions. In particular we shall examine the wavefront set, a refinement of singular support, of general distributions. This notion is fundamental to the general idea of 'microlocalization.'

### 4.6. Characteristic variety

If $A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$, the left-reduced symbol is elliptic at $(\bar{x}, \bar{\xi}) \in \mathbb{R}^{n} \times\left(\mathbb{R}^{n} \backslash\{0\}\right)$ if there exists $\epsilon>0$ such that
8.1

$$
\begin{gather*}
\left|\sigma_{L}(A)(x, \xi)\right| \geq \epsilon|\xi|^{m} \quad \text { in } \\
\left\{(x, \xi) \in \mathbb{R}^{n} \times\left(\mathbb{R}^{n} \backslash\{0\}\right) ;|x-\bar{x}| \leq \epsilon,\left|\frac{\xi}{|\xi|}-\frac{\bar{\xi}}{|\xi|}\right| \leq \epsilon,|\xi| \geq \frac{1}{\epsilon}\right\} \tag{4.52}
\end{gather*}
$$

Directly from the definition, ellipticity at $(\bar{x}, \bar{\xi})$ is actually a property of the principal symbol, $\sigma_{m}(A)$ and if $A$ is elliptic at $(\bar{x}, \bar{\xi})$ then it is elliptic at $(\bar{x}, t \bar{\xi})$ for any $t>0$. Clearly

$$
\left\{(\bar{x}, \bar{\xi}) \in \mathbb{R}^{n} \times\left(\mathbb{R}^{n} \backslash\{0\}\right) ; A \text { is elliptic (of order } m \text { ) at }(\bar{x}, \bar{\xi})\right\}
$$

is an open cone in $\mathbb{R}^{n} \times\left(\mathbb{R}^{n} \backslash\{0\}\right)$. The complement
(4.53) $\Sigma_{m}(A)=\left\{(\bar{x}, \bar{\xi}) \in \mathbb{R}^{n} \times\left(\mathbb{R}^{n} \backslash\{0\}\right) ; A\right.$ is not elliptic of order $m$ at $\left.(\bar{x}, \bar{\xi})\right\}$
is therefore a closed conic subset of $\mathbb{R}^{n} \times\left(\mathbb{R}^{n} \backslash\{0\}\right)$; it is the characteristic set (or variety) of $A$. Since the product of two symbols is only elliptic at $(\bar{x}, \bar{\xi})$ if they are both elliptic there, if follows from the composition properties of pseudodifferential operators that

$$
\Sigma_{m+m^{\prime}}(A \circ B)=\Sigma_{m}(A) \cup \Sigma_{m^{\prime}}(B)
$$

### 4.7. Wavefront set

We adopt the following bald definition:

$$
\text { If } u \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right)=\left\{u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) ; \operatorname{supp}(u) \Subset \mathbb{R}^{n}\right\} \text { then }
$$

$$
\begin{equation*}
\mathrm{WF}(u)=\bigcap\left\{\Sigma_{0}(A) ; A \in \Psi_{\infty}^{0}\left(\mathbb{R}^{n}\right) \text { and } A u \in \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right)\right\} \tag{4.55}
\end{equation*}
$$

Thus $\mathrm{WF}(u) \subset \mathbb{R}^{n} \times\left(\mathbb{R}^{n} \backslash\{0\}\right)$ is always a closed conic set, being the intersection of such sets. The first thing we wish to show is that $\mathrm{WF}(u)$ is a refinement of sing $\operatorname{supp}(u)$. Let

$$
\begin{array}{|l|}
\hline 8.5 \\
\hline
\end{array}
$$

$$
\begin{equation*}
\pi: \mathbb{R}^{n} \times\left(\mathbb{R}^{n} \backslash\{0\}\right) \ni(x, \xi) \longmapsto x \in \mathbb{R}^{n} \tag{4.56}
\end{equation*}
$$

be projection onto the first factor.

$$
\begin{equation*}
\pi(\mathrm{WF}(u))=\operatorname{sing} \operatorname{supp}(u) \tag{4.57}
\end{equation*}
$$

Proof. The inclusion $\pi(\operatorname{WF}(u)) \subset \operatorname{sing} \operatorname{supp}(w)$ is straightforward. Indeed, if $\bar{x} \notin \operatorname{sing} \operatorname{supp}(u)$ then there exists $\phi \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)$ with $\phi(\bar{x}) \neq 0$ such that $\phi u \in$ $\mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right)$. Of course as a multiplication operator, $\phi \in \Psi_{\infty}^{0}\left(\mathbb{R}^{n}\right)$ and $\Sigma_{0}(\phi) \not \supset(\bar{x}, \bar{\xi})$ for any $\bar{\xi} \neq 0$. Thus the definition ( $\overline{4.55) \text { ) shows that }(\bar{x}, \bar{\xi}) \notin \mathrm{WF}(u) \text { for all } \bar{\xi} \in \mathbb{R}^{n} \backslash 0}$ proving the inclusion.

Using the calculus of pseudodifferential operators, the opposite inclusion,

$$
\begin{equation*}
\pi(\mathrm{WF}(u)) \supset \operatorname{sing} \operatorname{supp}(u) \tag{4.58}
\end{equation*}
$$

is only a little more complicated. Thus we have to show that if $(\bar{x}, \bar{\xi}) \notin \mathrm{WF}(u)$ for all $\bar{\xi} \in \mathbb{R}^{n} \backslash 0$ then $\bar{x} \notin \operatorname{sing} \operatorname{supp}(u)$. The hypothesis is that for each $(\bar{x}, \bar{\xi}), \bar{\xi} \in \mathbb{R}^{n} \backslash 0$, there exists $A \in \Psi_{\infty}^{0}\left(\mathbb{R}^{n}\right)$ such that $A$ is elliptic at $(\bar{x}, \bar{\xi})$ and $A u \in \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right)$. The set of elliptic points is open so there exists $\epsilon=\epsilon(\bar{\xi})>0$ such that $A$ is elliptic on

$$
\begin{equation*}
\left\{(x, \xi) \in \mathbb{R}^{n} \times\left(\mathbb{R}^{n} \backslash 0\right) ;|x-\bar{x}|<\epsilon,\left|\frac{\xi}{|\xi|}-\frac{\bar{\xi}}{|\bar{\xi}|}\right|<\epsilon\right\} . \tag{4.59}
\end{equation*}
$$

Let $B_{j}, j=1, \ldots, N$ be a finite set of such operators associated to $\bar{\xi}_{j}$ and such that the corresponding sets in ( $\overline{4.59)}$ cover $\{\bar{x}\} \times\left(\mathbb{R}^{n} \backslash 0\right)$; the finiteness follows from the compactness of the sphere. Then consider

$$
B=\sum_{j=1}^{N} B_{j}^{*} B_{j} \Longrightarrow B u \in \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right)
$$

This operator $B$ is elliptic at $(\bar{x}, \xi)$, for all $\xi \neq 0$. Thus if $\phi \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right), 0 \leq \phi(x) \leq 1$, has support sufficiently close to $\bar{x}, \phi(x)=1$ in $|x-\bar{x}|<\epsilon / 2$ then, since $B$ has nonnegative principal symbol

$$
\begin{equation*}
B+(1-\phi) \in \Psi_{\infty}^{0}\left(\mathbb{R}^{n}\right) \tag{4.60}
\end{equation*}
$$

is globally elliptic. Thus, by Lemma $\frac{15.31}{2.7}$ there exists $G \in \Psi_{\infty}^{0}\left(\mathbb{R}^{n}\right)$ which is a parametrix for $B+(1-\phi)$ :

$$
\begin{equation*}
\mathrm{Id} \equiv G \circ B+G(1-\phi) \quad \bmod \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right) \tag{4.61}
\end{equation*}
$$

Let $\psi \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)$ be such that $\operatorname{supp}(\psi) \subset\{\phi=1\}$ and $\psi(\bar{x}) \neq 0$. Then, from the reduction formula

$$
\psi \circ G \circ(1-\phi) \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)
$$

Thus from (18.9 1.61 ) we find

$$
\psi u=\psi G \circ B u+\psi G(1-\phi) u \in \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right)
$$

Thus $\bar{x} \notin \operatorname{sing} \operatorname{supp}(u)$ and the proposition is proved.
We extend the definition to general tempered distributions by setting

$$
\begin{equation*}
\mathrm{WF}(u)=\bigcup_{\phi \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)} \mathrm{WF}(\phi u), u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \tag{4.62}
\end{equation*}
$$

Then $\left(\frac{8.7}{4.57}\right)$ holds for every $u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$.

### 4.8. Essential support

Next we shall consider the notion of the essential support of a pseudodifferential operator. If $a \in S_{\infty}^{m}\left(\mathbb{R}^{N} ; \mathbb{R}^{n}\right)$ we define the cone support of $a$ by

$$
\begin{align*}
\operatorname{cone} \operatorname{supp}(a)= & \left\{(\bar{x}, \bar{\xi}) \in \mathbb{R}^{N} \times\left(\mathbb{R}^{n} \backslash 0\right) ; \exists \epsilon>0 \text { and } \forall M \in \mathbb{R}, \exists C_{M}\right. \text { s.t. } \\
& \left.|a(x, \xi)| \leq C_{M}\langle\xi\rangle^{-M} \text { if }|x-\bar{x}| \leq \epsilon,\left|\frac{\xi}{|\xi|}-\frac{\bar{\xi}}{|\bar{\xi}|}\right| \leq \epsilon\right\}^{\complement} \tag{4.63}
\end{align*}
$$

8.11

This is clearly a closed conic set in $\mathbb{R}^{N} \times\left(\mathbb{R}^{n} \backslash 0\right)$. By definition the symbol decays rapidly outside this cone, in fact even more is true.

Lemma 4.3. If $a \in S_{\infty}^{\infty}\left(\mathbb{R}^{N} ; \mathbb{R}^{n}\right)$ then

$$
(\bar{x}, \bar{\eta}) \notin \text { cone } \operatorname{supp}(a) \Longrightarrow
$$

$$
\begin{equation*}
\exists \epsilon>0 \text { s.t. } \forall M, \alpha, \beta \exists C_{M} \text { with } \tag{4.64}
\end{equation*}
$$

$$
\left|D_{x}^{\alpha} D_{\xi}^{\beta} a(x, \eta)\right| \leq C_{M}\langle\eta\rangle^{-M} \text { if }|x-\bar{x}|<\epsilon,\left|\frac{\eta}{|\eta|}-\frac{\bar{\eta}}{|\bar{\eta}|}\right|<\epsilon
$$

Proof. To prove ( $\left(\frac{8.15}{4.64}\right)$ it suffices to show it to be valid for $D_{x_{j}} a, D_{\xi_{k}} a$ and then use an inductive argument, i.e. to show that

$$
\begin{equation*}
\operatorname{cone} \operatorname{supp}\left(D_{x_{j}} a\right), \text { conesupp }\left(D_{\xi_{k}} a\right) \subset \operatorname{conesupp}(a) \tag{4.65}
\end{equation*}
$$

Arguing by contradiction suppose that $D_{x_{\ell}} a$ does not decay to order $M$ in any cone around $(\bar{x}, \bar{\xi}) \notin$ conesupp. Then there exists a sequence $\left(x_{j}, \xi_{j}\right)$ with

$$
\left\{\begin{array}{l}
x_{j} \longrightarrow \bar{x},\left|\frac{\xi_{j}}{\left|\xi_{j}\right|}-\frac{\bar{\xi}}{|\bar{\xi}|}\right| \longrightarrow 0,\left|\xi_{j}\right| \longrightarrow \infty  \tag{4.66}\\
\text { and }\left|D_{x_{\ell}} a\left(x_{j}, \xi_{j}\right)\right|>j\left\langle\xi_{j}\right\rangle^{M}
\end{array}\right.
$$

We can assume that $M<m$, since $a \in S_{\infty}^{m}\left(\mathbb{R}^{n} ; \mathbb{R}^{N}\right)$. Applying Taylor's formula with remainder, and using the symbol bounds on $D_{x_{j}}^{2} a$, gives
(4.67) $a\left(x_{j}+t e_{\ell}, \xi_{j}\right)=a\left(x_{j}, \xi_{j}\right)+i t\left(D_{x_{j}} a\right)\left(x_{j}, \xi_{j}\right)+O\left(t^{2}\left\langle\xi_{j}\right\rangle^{m}\right),\left(e_{\ell}\right)_{j}=\delta_{\ell j}$
providing $|t|<1$ Taking $t=\left\langle\xi_{j}\right\rangle^{M-m} \longrightarrow 0$ as $j \longrightarrow \infty$, the first and third terms on the right in (4.67) are small compared to the second, so

$$
\begin{equation*}
\left|a\left(x_{j}+\left\langle\xi_{j}\right\rangle^{\frac{M-m}{2}}, \xi_{j}\right)\right|>\left\langle\xi_{j}\right\rangle^{2 M-m} \tag{4.68}
\end{equation*}
$$

contradicting the assumption that $(\bar{x}, \bar{\xi}) \notin$ cone $\operatorname{supp}(a)$. A similar argument applies to $D_{\xi_{\ell}} a$ so ( $\left.\overline{4.64}\right)$, and hence the lemma, is proved.

For a pseudodifferential operator we define the essential support by
8.12

$$
\begin{equation*}
\mathrm{WF}^{\prime}(A)=\operatorname{conesupp}\left(\sigma_{L}(A)\right) \subset \mathbb{R}^{n} \times\left(\mathbb{R}^{n} \backslash 0\right) \tag{4.69}
\end{equation*}
$$

Lemma 4.4. For every $A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$

$$
\begin{equation*}
\mathrm{WF}^{\prime}(A)=\operatorname{cone} \operatorname{supp}\left(\sigma_{R}(A)\right) \tag{4.70}
\end{equation*}
$$

Proof. Using ( $\frac{8.15}{1.64}$ ) and the formula relating $\sigma_{R}(A)$ to $\sigma_{L}(A)$ we conclude that

$$
\begin{equation*}
\operatorname{cone} \operatorname{supp}\left(\sigma_{L}(A)\right)=\operatorname{cone} \operatorname{supp}\left(\sigma_{R}(A)\right) \tag{4.71}
\end{equation*}
$$

from which $\left(\frac{8.14}{4.70}\right)$ follows.
A similar argument shows that

$$
\begin{equation*}
\mathrm{WF}^{\prime}(A \circ B) \subset \mathrm{WF}^{\prime}(A) \cap \mathrm{WF}^{\prime}(B) \tag{4.72}
\end{equation*}
$$

Indeed the asymptotic formula for $\sigma_{L}(A \circ B)$ in terms of $\sigma_{L}(A)$ and $\sigma_{L}(B)$ shows that

$$
\begin{equation*}
\operatorname{cone} \operatorname{supp}\left(\sigma_{L}(A \circ B)\right) \subset \operatorname{cone} \operatorname{supp}\left(\sigma_{L}(A)\right) \cap \operatorname{cone\operatorname {supp}}\left(\sigma_{L}(B)\right) \tag{4.73}
\end{equation*}
$$

which is the same thing.

### 4.9. Microlocal parametrices

The concept of essential support allows us to refine the notion of a parametrix for an elliptic operator to that of a microlocal parametrix.
9.1 Lemma 4.5. If $A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ and $z \notin \Sigma_{m}(A)$ then there exists a microlocal parametrix at $z, B \in \Psi_{\infty}^{-m}\left(\mathbb{R}^{n}\right)$ such that

$$
\begin{equation*}
z \notin \mathrm{WF}^{\prime}(\operatorname{Id}-A B) \text { and } z \notin \mathrm{WF}^{\prime}(\operatorname{Id}-B A) . \tag{4.74}
\end{equation*}
$$

Proof. If $z=(\bar{x}, \bar{\xi}), \bar{\xi} \neq 0$, consider the symbol

$$
\begin{equation*}
\gamma_{\epsilon}(x, \xi)=\phi\left(\frac{x-\bar{x}}{\epsilon}\right)(1-\phi)(\epsilon \xi) \phi\left(\left(\frac{\xi}{|\xi|}-\frac{\bar{\xi}}{|\bar{\xi}|}\right) / \epsilon\right) \tag{4.75}
\end{equation*}
$$

where as usual $\phi \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right), \phi(\zeta)=1$ in $|\zeta| \leq \frac{1}{2}, \phi(\zeta)=0$ in $|\zeta| \geq 1$. Thus $\gamma_{\epsilon} \in S_{\infty}^{0}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right)$ has support in

$$
\begin{equation*}
|x-\bar{x}| \leq \epsilon,|\xi| \geq \frac{1}{2 \epsilon},\left|\frac{\xi}{|\xi|}-\frac{\bar{\xi}}{|\bar{\xi}|}\right| \leq \epsilon \tag{4.76}
\end{equation*}
$$

and is identically equal to one, and hence elliptic, on a similar smaller set

$$
\begin{equation*}
|x-\bar{x}|<\frac{\epsilon}{2},|\xi| \geq \frac{1}{\epsilon},\left|\frac{\xi}{|\xi|}-\frac{\bar{\xi}}{|\bar{\xi}|}\right| \leq \frac{\epsilon}{2} . \tag{4.77}
\end{equation*}
$$

Define $L_{\epsilon} \in \Psi_{\infty}^{0}\left(\mathbb{R}^{n}\right)$ by $\sigma_{L}\left(L_{\epsilon}\right)=\gamma_{\epsilon}$. Thus, for any $\epsilon>0$,
(4.78) $\quad z \notin \mathrm{WF}^{\prime}\left(\operatorname{Id}-L_{\epsilon}\right), \mathrm{WF}^{\prime}\left(L_{\epsilon}\right) \subset\left\{(x, \xi) ;|x-\bar{x}| \leq \epsilon\right.$ and $\left.\left|\frac{\xi}{|\xi|}-\frac{\bar{\xi}}{|\bar{\xi}|}\right| \leq \epsilon\right\}$.

Let $G_{2 m} \in \Psi_{\infty}^{2 m}\left(\mathbb{R}^{n}\right)$ be a globally elliptic operator with positive principal symbol. For example take $\sigma_{L}\left(G_{2 m}\right)=\left(1+|\xi|^{2}\right)^{m}$, so $G_{s} \circ G_{t}=G_{s+t}$ for any $s$, $t \in \mathbb{R}$. Now consider the operator

$$
\begin{equation*}
J=\left(\operatorname{Id}-L_{\epsilon}\right) \circ G_{2 m}+A^{*} A \in \Psi_{\infty}^{2 m}\left(\mathbb{R}^{n}\right) \tag{4.79}
\end{equation*}
$$

The principal symbol of $J$ is $\left(1-\gamma_{\epsilon}\right)\left(1+|\xi|^{2}\right)^{m}+\left|\sigma_{m}(A)\right|^{2}$ whigh is globally elliptic if $\epsilon>0$ is small enough (so that $\sigma_{m}(A)$ is elliptic on the set (4.76)). According to Lemma $\frac{5.1}{2.75}, J$ has a global parametrix $H \in \Psi_{\infty}^{-2 m}\left(\mathbb{R}^{n}\right)$. Then

$$
\begin{equation*}
B=H \circ A^{*} \in \Psi_{\infty}^{-m}\left(\mathbb{R}^{n}\right) \tag{4.80}
\end{equation*}
$$

is a microlocal right parametrix for $A$ in the sense that $B \circ A-\mathrm{Id}=R_{R}$ with $z \notin \mathrm{WF}^{\prime}\left(R_{R}\right)$ since

$$
\begin{align*}
R_{R}=B \circ A-\mathrm{Id}=H \circ A^{*} \circ A-\mathrm{Id} &  \tag{4.81}\\
& =(H \circ J-\mathrm{Id})+H \circ\left(\mathrm{Id}-L_{\epsilon}\right) G_{2 m} \circ A
\end{align*}
$$

and the first term on the right is in $\Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$ whilst $z$ is not in the operator wavefront set of $\left(\operatorname{Id}-L_{\epsilon}\right)$ and hence not in the operator wavefront set of the second term.

By a completely analogous construction we can find a left microlocal parametrix. Namely $\left(\operatorname{Id}-L_{\epsilon}\right) \circ G_{2 m}+A \circ A^{*}$ is also globally elliptic with parametrix $H^{\prime}$ and then $B^{\prime}=A^{*} \circ H^{\prime}$ satisfies

$$
\begin{equation*}
B^{\prime} \circ A-\mathrm{Id}=R_{L}, z \notin \mathrm{WF}^{\prime}\left(R_{L}\right) \tag{4.82}
\end{equation*}
$$

Then, as usual,

$$
\begin{equation*}
B=\left(B^{\prime} \circ A-R_{L}\right) B=B^{\prime}(A \circ B)-R_{L} B=B^{\prime}+B^{\prime} R_{R}-R_{L} B \tag{4.83}
\end{equation*}
$$

so $z \notin \mathrm{WF}^{\prime}\left(B-B^{\prime}\right)$, which implies that $B$ is both a left and right microlocal parametrix.

In fact this argument shows that such a left parametrix is essentially unique. See Problem $\frac{.9 .2 .}{4 .}$

### 4.10. Microlocality

Now we can consider the relationship between these two notions of wavefront set.

$$
\begin{equation*}
\mathrm{WF}(A u) \subset \mathrm{WF}^{\prime}(A) \cap \mathrm{WF}(u) \quad \forall A \in \Psi_{\infty}^{\infty}\left(\mathbb{R}^{n}\right), u \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right) \tag{4.84}
\end{equation*}
$$

Proof. We need to show that

$$
\begin{equation*}
\mathrm{WF}(A u) \subset \mathrm{WF}^{\prime}(A) \text { and } \mathrm{WF}(A u) \subset \mathrm{WF}(u) \tag{4.85}
\end{equation*}
$$

the second being the usual definition of microlocality. The first inclusion is easy. Suppose $(\bar{x}, \bar{\xi}) \notin$ cone $\operatorname{supp} \sigma_{L}(A)$. If we choose $B \in \Psi_{\infty}^{0}\left(\mathbb{R}^{n}\right)$ with $\sigma_{L}(B)$ supported in a small cone around $(\bar{x}, \bar{\xi})$ then we can arrange

$$
\begin{equation*}
(\bar{x}, \bar{\xi}) \notin \Sigma_{0}(B), \mathrm{WF}^{\prime}(B) \cap \mathrm{WF}^{\prime}(A)=\emptyset \tag{4.86}
\end{equation*}
$$

Then from ( $\frac{8.20}{4.72}$ ), $\mathrm{WF}^{\prime}(B A)=\emptyset$ so $B A \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$ and $B A u \in \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right)$. Thus $(\bar{x}, \bar{\xi}) \notin \mathrm{WF}(A u)$.

Similarly suppose $(\bar{x}, \bar{\xi}) \notin \mathrm{WF}(u)$. Then there exists $G \in \Psi_{\infty}^{0}\left(\mathbb{R}^{n}\right)$ which is elliptic at $(\bar{x}, \bar{\xi})$ with $G u \in \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right)$. Let $B$ be a microlocal parametrix for $G$ at $(\bar{x}, \bar{\xi})$ as in Lemma $\frac{9.1}{4.5}$. Thus

$$
\begin{equation*}
u=B G u+S u, \quad(\bar{x}, \bar{\xi}) \notin \mathrm{WF}^{\prime}(S) \tag{4.87}
\end{equation*}
$$

Now apply $A$ to this identity. Since, by assumption, $G u \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)$ the first term on the right in

$$
\begin{equation*}
A u=A B G u+A S u \tag{4.88}
\end{equation*}
$$

is smooth. Since, by $\left(\frac{8.20}{4.72}\right),(\bar{x}, \bar{\xi}) \notin \mathrm{WF}^{\prime}(A S)$ it follows from the first part of the argument above that $(\bar{x}, \bar{\xi}) \notin \mathrm{WF}(A S u)$ and hence $(\bar{x}, \bar{\xi}) \notin \mathrm{WF}(A u)$.

We can deduce from the existence of microlocal parametrices at elliptic points a partial converse of (8.24).

Proposition 4.7. For any $u \in \mathcal{C}^{-\infty}\left(\mathbb{R}^{n}\right)$ and any $A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$

$$
\begin{equation*}
\mathrm{WF}(u) \subset \mathrm{WF}(A u) \cup \Sigma_{m}(A) \tag{4.89}
\end{equation*}
$$

Prgo. . If $(\bar{x}, \bar{\xi}) \notin \Sigma_{m}(A)$ then, by definition, $A$ is elliptic at $(\bar{x}, \bar{\xi})$. Thus, by Lemma $\frac{9.1}{4.5}, A$ has a microlocal parametrix $B$, so

$$
\begin{equation*}
u=B A u+S u,(\bar{x}, \bar{\xi}) \notin \mathrm{WF}^{\prime}(S) \tag{4.90}
\end{equation*}
$$

It follows that $(\bar{x}, \bar{\xi}) \notin \mathrm{WF}(A u)$ implies that $(\bar{x}, \bar{\xi}) \notin \mathrm{WF}(u)$ proving the Proposition.

### 4.11. Explicit formulations

From this discussion of $\mathrm{WF}^{\prime}(A)$ we can easily find a 'local coordinate' formulations of $\mathrm{WF}(u)$ in general.
8.25 Lemma 4.6. If $(\bar{x}, \bar{\xi}) \in \mathbb{R}^{n} \times\left(\mathbb{R}^{n} \backslash 0\right)$ and $u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ then $(\bar{x}, \bar{\xi}) \notin \mathrm{WF}(u)$ if and only if there exists $\phi \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)$ with $\phi(\bar{x}) \neq 0$ such that for some $\epsilon>0$, and for all $M$ there exists $C_{M}$ with
8.26

$$
\begin{equation*}
|\widehat{\phi u}(\xi)| \leq C_{M}\langle\xi\rangle^{M} \text { in }\left|\frac{\xi}{|\xi|}-\frac{\bar{\xi}}{|\xi|}\right|<\epsilon . \tag{4.91}
\end{equation*}
$$

Proof. If $\zeta \in \mathcal{C}^{\infty}(\mathbb{R}), \zeta(\xi) \equiv 1$ in $|\xi|<\frac{\epsilon}{2}$ and $\operatorname{supp}(\zeta) \subset\left[\frac{-3 \epsilon}{4}, \frac{3 \epsilon}{4}\right]$ then

$$
\begin{equation*}
\gamma(\xi)=(1-\zeta)(\xi) \cdot \zeta\left(\frac{\xi}{|\xi|}-\frac{\bar{x}}{|\bar{x}|}\right) \in S_{\infty}^{0}\left(\mathbb{R}^{n}\right) \tag{4.92}
\end{equation*}
$$

is elliptic at $\bar{\xi}$ and from $\left(\frac{8.26}{1.91}\right)$

$$
\begin{equation*}
\gamma(\xi) \cdot \widehat{\phi u}(\xi) \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{4.93}
\end{equation*}
$$

Thus if $\sigma_{R}(A)=\phi_{1}(x) \gamma(\xi)$ then $A\left(\phi_{2} u\right) \in \mathcal{C}^{\infty}$ where $\phi_{1} \phi_{2}=\phi, \phi_{1}(\bar{x}), \phi_{2}(\bar{x}) \neq 0$, $\phi_{1}, \phi_{2} \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)$. Thus $(\bar{x}, \bar{\xi}) \notin \mathrm{WF}(u)$. Conversely. if $(\bar{x}, \bar{\xi}) \notin \mathrm{WF}(u)$ and $A$ is chosen as above then $A\left(\phi_{1} u\right) \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ and Lemma 4.6 holds.

### 4.12. Wavefront set of $K_{A}$

At this stage, a natural thing to look at is the wavefront set of the kernel of a pseudodifferential operator, since these kernels are certainly an interesting class of distributions.

Proposition 4.8. If $A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ then

$$
\begin{align*}
\mathrm{WF}\left(K_{A}\right)= & \left\{(x, y, \xi, \eta) \in \mathbb{R}^{2 n} \times\left(\mathbb{R}^{2 n} \backslash 0\right)\right. \\
& \left.x=y, \xi+\eta=0 \text { and }(x, \xi) \in \mathrm{WF}^{\prime}(A)\right\} . \tag{4.94}
\end{align*}
$$

In particular this shows that $\mathrm{WF}^{\prime}(A)$ determines $\mathrm{WF}\left(K_{A}\right)$ and conversely.
Proof. Using Proposition $\frac{8.6}{4.5}$ we know that $\pi\left(\operatorname{WF}\left(K_{A}\right)\right) \subset\{(x, x)\}$ so

$$
\mathrm{WF}\left(K_{A}\right) \subset\{(x, x ; \xi, \eta)\}
$$

To find the wave front set more precisely consider the kernel

$$
K_{A}(x, y)=(2 \pi)^{-n} \int e^{i(x-y) \cdot \xi} b(x, \xi) d \xi
$$

where we can assume $|x-y|<1$ on $\operatorname{supp}\left(K_{A}\right)$. Thus is $\phi \in \mathcal{C}_{c}^{\infty}(X)$ then

$$
g(x, y)=K_{A}(x, y) \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right)
$$

and

$$
\begin{aligned}
\hat{g}(\zeta, \eta) & =(2 \pi)^{-n} \int e^{-i \zeta x-i \eta y} e^{i(x-y) \cdot \zeta}(\phi b)(x, \xi) d \zeta d x d y \\
& =\int e^{-i(\zeta+\eta) \cdot x}(\phi b)(x,-\eta) d x \\
& =\widehat{\phi b}(\zeta+\eta,-\eta)
\end{aligned}
$$

The fact that $\phi b$ is a symbol of compact support in $x$ means that for every $M$

$$
|\widehat{\phi b}(\zeta+\eta,-\eta)| \leq C_{M}(\langle\zeta+\eta\rangle)^{-M}\langle\eta\rangle^{m}
$$

This is rapidly decreasing if $\zeta \neq-\eta$, so

$$
\mathrm{WF}\left(K_{A}\right) \subset\{(x, x, \eta,-\eta)\} \text { as claimed. }
$$

Moreover if $(\bar{x}, \bar{\eta}) \notin \mathrm{WF}^{\prime}(A)$ then choosing $\phi$ to have small support near $\bar{x}$ makes $\widehat{\phi b}$ rapidly decreasing near $-\bar{\eta}$ for all $\zeta$. This proves Proposition $\frac{8.27}{4.8}$

### 4.13. Elementary calculus of wavefront sets

We want to achieve a reasonable understanding, in terms of wavefront sets, of three fundamental operations. These are

> Pull-back: $F^{*} u$
> Push-forward: $F_{*} u$ and
> Multiplication: $u_{1} \cdot u_{2}$.

In order to begin to analyze these three operations we shall first introduce and discuss some other more "elementary" operations:
proj
rest

## expr

inv

$$
\begin{equation*}
\text { Pairing: }(u, v) \longrightarrow\langle u, v\rangle=\int u(x) \overline{v(x)} d x \tag{4.98}
\end{equation*}
$$

$$
\begin{equation*}
\text { Projection: } u(x, y) \longmapsto \int u(x, y) d y \tag{4.99}
\end{equation*}
$$

$$
\begin{equation*}
\text { Restriction: } u(x, y) \longmapsto u(x, 0) \tag{4.100}
\end{equation*}
$$

$$
\begin{equation*}
\text { Exterior product: }(u, v) \longmapsto(u \boxtimes v)(x, y)=u(x) v(y) \tag{4.101}
\end{equation*}
$$

 combination of ( $\frac{\text { mut }}{4.97}$ ) and ( $\frac{1.95}{4.95) \text {. Conversely the three fundamental operations can }}$ be expressed in terms of these elementary ones. We can give direct definitions of the latter which we then use to analyze the former. We shall start with the pairing in (4.98).

### 4.14. Pairing

We know how to 'pair' a distribution and a $\mathcal{C}^{\infty}$ function. If both are $\mathcal{C}^{\infty}$ and have compact supports then

$$
\begin{equation*}
\left\langle u_{1}, u_{2}\right\rangle=\int u_{1}(x) \overline{u_{2}(x)} d x \tag{4.103}
\end{equation*}
$$

and in general this pairing extends by continuity to either $\mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right) \times \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right)$ or $\mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right) \times \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right)$ Suppose both $u_{1}$ and $u_{2}$ are distributions, when can we pair them?
then if $A \in \Psi_{\infty}^{0}\left(\mathbb{R}^{n}\right)$ has

$$
\begin{equation*}
\mathrm{WF}\left(u_{1}\right) \cap \mathrm{WF}^{\prime}(A)=\emptyset, \quad \mathrm{WF}\left(u_{2}\right) \cap \mathrm{WF}^{\prime}\left(\operatorname{Id}-A^{*}\right)=\emptyset \tag{4.105}
\end{equation*}
$$

the bilinear form
9.11

$$
\begin{equation*}
\left\langle u_{1}, u_{2}\right\rangle=\left\langle A u_{1}, u_{2}\right\rangle+\left\langle u_{1},\left(\operatorname{Id}-A^{*}\right) u_{2}\right\rangle \tag{4.106}
\end{equation*}
$$

is independent of the choice of $A$.
Notice that $A$ satisfying $\left(\frac{9.10}{4.105}\right)$ does indeed exist, just choose $a \in S_{\infty}^{0}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right)$ to be identically 1 on $\operatorname{WF}\left(u_{2}\right)$, but to have conesupp $(a) \cap \mathrm{WF}\left(u_{1}\right)=\emptyset$, possible because of (4.104), and set $A=q_{L}(a)$.

Proof. Of course $\left(\frac{9.11}{4.106)}\right.$ makes sense because $A u_{1},\left(\operatorname{Id}-A^{*}\right) u_{2} \in \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right)$. To prove that this definition is independent of the choice of $A$, suppose $A^{\prime}$ also satisfies (4.105). Set

$$
\begin{equation*}
\left\langle u_{1}, u_{2}\right\rangle^{\prime}=\left\langle A^{\prime} u_{1}, u_{2}\right\rangle+\left\langle u_{1},\left(\operatorname{Id}-A^{\prime}\right)^{*} u_{2}\right\rangle \tag{4.107}
\end{equation*}
$$

Then

$$
\begin{equation*}
\mathrm{WF}^{\prime}\left(A-A^{\prime}\right) \cap \mathrm{WF}\left(u_{1}\right)=\mathrm{WF}^{\prime}\left(\left(A-A^{\prime}\right)^{*}\right) \cap \mathrm{WF}\left(u_{2}\right)=\emptyset \tag{4.108}
\end{equation*}
$$

The difference can be written
9.12

$$
\begin{equation*}
\left\langle u_{1}, u_{2}\right\rangle^{\prime}-\left\langle u_{1}, u_{2}\right\rangle=\left\langle\left(A-A^{\prime}\right) u_{1}, u_{2}\right\rangle-\left\langle u_{1},\left(A-A^{\prime}\right)^{*} u_{2}\right\rangle \tag{4.109}
\end{equation*}
$$

Naturally we expect this to be zero, but this is not quite obvious since $u_{1}$ and $u_{2}$ are both distributions. We need an approximation argument to finish the proof.

Choose $B \in \Psi_{\infty}^{0}\left(\mathbb{R}^{n}\right)$ with

$$
\begin{gather*}
\mathrm{WF}^{\prime}(B) \cap \mathrm{WF}\left(u_{1}\right)=\mathrm{WF}^{\prime}(B) \cap \mathrm{WF}\left(u_{2}\right)=\emptyset \\
\mathrm{WF}^{\prime}(\mathrm{Id}-B) \cap \mathrm{WF}\left(A-A^{\prime}\right)=\emptyset \tag{4.110}
\end{gather*}
$$

If $v_{n} \longrightarrow u_{2}$, in $\mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right), v_{n} \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)$ then

$$
\begin{equation*}
w_{n}=\phi\left[(\operatorname{Id}-B) v_{n}+B u_{2}\right] \longrightarrow u_{2} \tag{4.111}
\end{equation*}
$$

if $\phi \equiv 1$ in a neighbourhood of $\operatorname{supp}\left(u_{2}\right), \phi \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)$. Here $B u_{2} \in \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right)$, so (4.112)
$\left(A-A^{\prime}\right) w_{n}=\left(A-A^{\prime}\right) \phi(\operatorname{Id}-B) \cdot v_{n}+\left(A-A^{\prime}\right) \phi B u_{2} \longrightarrow\left(A-A^{\prime}\right) u_{2}$ in $\mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right)$,
since $\left(A-A^{\prime}\right) \phi(\operatorname{Id}-B) \in \Psi_{\infty}^{-\infty}\left(\mathbb{R}^{n}\right)$. Thus

$$
\begin{aligned}
& \left\langle\left(A-A^{\prime}\right) u_{1}, u_{2}\right\rangle \longrightarrow\left\langle\left(A-A^{\prime}\right) u_{1}, u_{2}\right\rangle \\
& \left\langle u_{1},\left(A-A^{\prime}\right)^{*} w_{n}\right\rangle \longrightarrow\left\langle u_{1},\left(A-A^{\prime}\right)^{*} u_{2}\right\rangle
\end{aligned}
$$

since $w_{n} \longrightarrow u_{2}$ in $\mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right)$ and $\left(A-A^{\prime}\right)^{*} w_{n} \longrightarrow\left(A-A^{\prime}\right)^{*} u_{2}$ in $\mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right)$. Thus

$$
\begin{equation*}
\left\langle u_{1}, u_{2}\right\rangle^{\prime}-\left\langle u_{1}, u_{2}\right\rangle=\lim _{n \rightarrow \infty}\left[\left\langle\left(A-A^{\prime}\right) u_{1}, w_{n}\right\rangle-\left\langle u_{1},\left(A-A^{\prime}\right)^{*} w_{n}\right]=0\right. \tag{4.113}
\end{equation*}
$$

Here we are using the complex pairing. If we define the real pairing by

$$
\begin{equation*}
\left(u_{1}, u_{2}\right)=\left\langle u_{1}, \bar{u}_{2}\right\rangle \tag{tabular}
\end{equation*}
$$

then we find

$$
\begin{equation*}
(x, \xi) \in \mathrm{WF}\left(u_{1}\right) \Longrightarrow(x,-\xi) \notin \mathrm{WF}\left(u_{2}\right) \tag{4.115}
\end{equation*}
$$

then the real pairing, defined by

$$
\begin{equation*}
\left(u_{1}, u_{2}\right)=\left(A u_{1}, u_{2}\right)+\left(u_{1}, A^{t} u_{2}\right) \tag{4.116}
\end{equation*}
$$

where $A$ satisfies $\left(\frac{9.10}{4.105}\right)$, is independent of $A$.
Proof. Notice that

$$
\begin{equation*}
\mathrm{WF}(\bar{u})=\left\{(x,-\xi) \in \mathbb{R}^{n} \times\left(\mathbb{R}^{n} \backslash 0\right) ;(x, \xi) \in \mathrm{WF}(u)\right\} \tag{4.117}
\end{equation*}
$$

We can write ( $\frac{9.17}{4.175)}$, using ( $\frac{9.15}{4.11} 4$ ), as

$$
\begin{equation*}
\left(u_{1}, u_{2}\right)=\left\langle A u_{1}, \bar{u}_{2}\right\rangle+\left\langle u_{1}, \overline{A^{t} u_{2}}\right\rangle \tag{4.118}
\end{equation*}
$$

Since, by definition, $\overline{A^{t} u_{2}}=A^{*} \bar{u}_{2}$,

$$
\begin{equation*}
\left(u_{1}, u_{2}\right)=\left\langle A u_{1}, \bar{u}_{2}\right\rangle+\left\langle u_{1}, A^{*} \bar{u}_{2}\right\rangle=\left\langle u_{1}, \bar{u}_{2}\right\rangle \tag{4.119}
\end{equation*}
$$

is defined by (9.11 $\left(\frac{9}{4.106)}\right.$ ) since $\left(\frac{9.17}{4.15} 5\right)$ translates to (9.9 14.104$)$.

### 4.15. Multiplication of distributions

The pairing result ( $\left(\frac{9.18}{4.116}\right)$ can be used to define the product of two distributions under the same hypotheses, (4.115).

Proposition 4.11. If $u_{1}, u_{2} \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right)$ satisfy

$$
\begin{equation*}
(x, \xi) \in \mathrm{WF}\left(u_{1}\right) \Longrightarrow(x,-\xi) \notin \mathrm{WF}\left(u_{2}\right) \tag{4.120}
\end{equation*}
$$

then the product of $u_{1}$ and $u_{2} \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right)$ is well-defined by

$$
\begin{align*}
& \text { (4.121) } u_{1} u_{2}(\phi)=\left(u_{1}, \phi u_{2}\right)=\left(\phi u_{1}, u_{2}\right) \forall \phi \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)  \tag{4.121}\\
& \text { using (19.18} 4.116) .
\end{align*}
$$

Proof. We only need to observe that if $u \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right)$ and $A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ has $\mathrm{WF}^{\prime}(A) \cap \mathrm{WF}(u)=\emptyset$ then for any fixed $\psi \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)$

$$
\begin{equation*}
\|\psi A \phi u\|_{C^{k}} \leq C\|\phi\|_{C^{p}} \quad p=k+N \tag{4.122}
\end{equation*}
$$

for s.me $N$, depending on $m$. This implies the continuity of $\phi \longmapsto u_{1} u_{2}(\phi)$ defined by (4.121).

### 4.16. Projection

Here we write $\mathbb{R}_{z}^{n}=\mathbb{R}_{x}^{p} \times \mathbb{R}_{y}^{k}$ and define a continuous linear map, which we write rather formally as an integral
10.1

$$
\begin{equation*}
\mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right) \ni u \longmapsto \int u(x, y) d y \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{p}\right) \tag{4.123}
\end{equation*}
$$

by pairing. If $\phi \in \mathcal{C}^{\infty}\left(\mathbb{R}^{p}\right)$ then
10.2

$$
\begin{equation*}
\pi_{1}^{*} \phi \in \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right), \quad \pi_{1}: \mathbb{R}^{n} \ni(x, y) \longmapsto x \in \mathbb{R}^{p} \tag{4.124}
\end{equation*}
$$

and for $u \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right)$ we define the formal 'integral' in (10.1 $\left.\frac{11}{4.12} 3\right)$ by
10.3

$$
\begin{equation*}
\left(\int u(x, y) d y, \phi\right)=\left(\left(\pi_{1}\right)_{*} u, \phi\right):=u\left(\pi_{1}^{*} \phi\right) . \tag{4.125}
\end{equation*}
$$

In this sense we see that the projection is dual to pull-back (on functions) under $\pi_{1}$, so is "push-forward under $\pi_{1}$," a special case of (4.96). The support of the projection satisfies

$$
\begin{equation*}
\operatorname{supp}\left(\left(\pi_{1}\right)_{*} u\right) \subset \pi_{1}(\operatorname{supp}(u)) \forall u \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right) \tag{4.126}
\end{equation*}
$$

as follows by duality from
10.5 (4.127)

$$
\operatorname{supp}\left(\pi_{1}^{*} \phi\right) \subset \pi_{1}^{-1}(\operatorname{supp} \phi)
$$

Proposition 4.12. Let $\pi_{1}: \mathbb{R}^{p+k} \longrightarrow \mathbb{R}^{p}$ be projection, then for every $u \in$ $\mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{p+k}\right)$
10.7

$$
\begin{align*}
& \mathrm{WF}\left(\left(\pi_{1}\right)_{*} u\right) \subset\left\{(x, \xi) \in \mathbb{R}^{p} \times\left(\mathbb{R}^{p} \backslash 0\right)\right. \\
& \left.\quad \exists y \in \mathbb{R}^{k} \text { with }(x, y, \xi, 0) \in \mathrm{WF}(u)\right\} \tag{4.128}
\end{align*}
$$

Proof. First notice that

$$
\begin{equation*}
\left(\pi_{1}\right)_{*}: \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{p}\right) \tag{4.129}
\end{equation*}
$$

Combining this with $\left(\frac{10.4}{4.12} 6\right)$ we see that

$$
\begin{equation*}
\operatorname{sing} \operatorname{supp}\left(\left(\pi_{1}\right)_{*} u\right) \subset \pi_{1}(\operatorname{sing} \operatorname{supp} u) \tag{tabular}
\end{equation*}
$$

which is at least consistent with Proposition ${ }^{\frac{10.6}{10 .} \text {. }}$. To prove the proposition in full let me restate the local characterization of the wavefront set, in terms of the Fourier transform:
10.10
10.11

$$
\begin{gather*}
u \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right), \mathrm{WF}(u) \cap(K \times \Gamma)=\emptyset, A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right), \mathrm{WF}^{\prime}(A) \subset K \times \Gamma \\
\Longrightarrow A u \in \mathcal{S}\left(\mathbb{R}^{n}\right) \tag{4.131}
\end{gather*}
$$

In particular

$$
\begin{equation*}
u \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right), \operatorname{WF}(u) \cap(K \times \Gamma)=\emptyset, \phi \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right), \operatorname{supp}(\phi) \subset K \tag{4.132}
\end{equation*}
$$

$\Longrightarrow \widehat{\phi u}(\xi)$ is rapidly decreasing in $\Gamma$.
Conversely suppose $\Gamma \subset \mathbb{R}^{n} \backslash 0$ is a closed cone and $u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ is such that for some $\phi \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)$
10.13

$$
\begin{equation*}
\widehat{\phi u}(\xi) \text { is rapidly decreasing in } \Gamma \tag{4.133}
\end{equation*}
$$

then

$$
\begin{equation*}
\mathrm{WF}(u) \cap\left\{x \in \mathbb{R}^{n} ; \phi(x) \neq 0\right\} \times \operatorname{int}(\Gamma)=\emptyset \tag{tabular}
\end{equation*}
$$

With these local tools at our disposal, let us attack ( 10.7128 ). We need to show that

$$
\begin{align*}
(\bar{x}, \bar{\xi}) \in \mathbb{R}^{p} \times & \left(\mathbb{R}^{p} \backslash 0\right) \text { s.t. }(\bar{x}, y, \bar{\xi}, 0) \notin \mathrm{WF}(u) \forall y \in \mathbb{R}^{n} \\
& \Longrightarrow(\bar{x}, \bar{\xi}) \notin \mathrm{WF}\left(\left(\pi_{1}\right)_{*} u\right) \tag{4.135}
\end{align*}
$$

Notice that, $\mathrm{WF}(u)$ being conic and $\pi(\mathrm{WF}(u))$ being compact, $\mathrm{WF}(u) \cap\left(\mathbb{R}^{n} \times S^{n-1}\right)$ is compact. The hypothesis ( $\overline{4.135)}$ is the statement that

$$
\begin{equation*}
\{\bar{x}\} \times \mathbb{R}^{k} \times \mathbb{S}^{n-1} \times\{0\} \cap \mathrm{WF}(u)=\emptyset \tag{4.136}
\end{equation*}
$$

Thus $\bar{x}$ has an open neighbourhood, $W$, in $\mathbb{R}^{p}$, and $(\bar{\xi}, 0)$ a conic neighbourhood $\gamma_{1}$ in $\left(\mathbb{R}^{n} \backslash 0\right)$ such that

$$
\begin{equation*}
\left(W \times \mathbb{R}^{k} \times \gamma_{1}\right) \cap \mathrm{WF}(u)=\emptyset \tag{4.137}
\end{equation*}
$$

Now if $\phi \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{p}\right)$ is chosen to have support in $W$

$$
\begin{equation*}
\widehat{\left(\pi_{1}^{*} \phi\right)} u(\xi, \eta) \text { is rapidly decreasing in } \gamma_{1} \tag{4.138}
\end{equation*}
$$

Set $v=\phi\left(\pi_{1}\right)_{*} u$. From the definition of projection and the identity

$$
\begin{equation*}
v=\phi\left(\pi_{1}\right)_{*} u=\left(\pi_{1}\right)_{*}\left[\left(\pi_{1}^{*} \phi\right) u\right] \tag{4.139}
\end{equation*}
$$

we have
10.17

$$
\begin{equation*}
\widehat{v}(\xi)=v\left(e^{-i x \cdot \xi}\right)=\left(\widehat{\left(\pi_{1}^{*} \phi\right) u}\right)(\xi, 0) \tag{4.140}
\end{equation*}
$$

Now ( $\left(\frac{10.16}{4.138}\right)$ shows that $\widehat{v}(\xi)$ is rapidly decreasing in $\gamma_{1} \cap\left(\mathbb{R}^{p} \times\{0\}\right)$, which is a cone around $\bar{\xi}$ in $\mathbb{R}^{p}$. Since $v=\phi\left(\pi_{1}\right)_{*} u$ this shows that $(\bar{x}, \bar{\xi}) \notin \mathrm{WF}\left(\left(\pi_{1}\right)_{*} u\right)$, as claimed.

Before going on to talk about the other operations, let me note a corollary of this which is useful and, even more, helps to explain what is going on:
10.18 Corollary 4.1. If $u \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right)$ and

$$
\begin{equation*}
\mathrm{WF}(u) \cap\left\{(x, y, \xi, 0) ; x \in \mathbb{R}^{p}, y \in \mathbb{R}^{k}, \xi \in \mathbb{R}^{p} \backslash 0\right\}=\emptyset \tag{4.141}
\end{equation*}
$$

then $\left(\pi_{1}\right)_{*}(u) \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)$.

Proof. Indeed, $\left(\frac{10 \cdot 7}{4.12} 8\right)$ says WF $\left(\left(\pi_{1}\right)_{*} u\right)=\emptyset$.
Here, the vectors $(x, y, \xi, 0)$ are the ones "normal" (as we shall see really conormal) to the surfaces over which we are integrating. Thus Lemma 4.7 and Corollary 4.1 both state that the only singularities that survive integration are the ones which are conormal to the surface along which we integrating; the ones even partially in the direction of integration are wiped out. This in particular fits with the fact that if we integrate in all variables then there are no singularities left.

### 4.17. Restriction

Next we wish to consider the restriction of a distribution to a subspace

$$
\begin{equation*}
\mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right) \ni u \longmapsto u \upharpoonright\{y=0\} \underset{c}{\in \mathcal{C}_{c}^{-\infty}}\left(\mathbb{R}^{p}\right) \tag{4.142}
\end{equation*}
$$

This is not always defined, i.e. no reasonable map ( $\left(\frac{10}{4.142}\right)$ exists for all distributions. However under an appropriate condition on the wavefront set we can interpret ( ${ }^{10.142)}$ in terms of pairing, using our definition of products. Thus let

$$
\begin{equation*}
\iota: \mathbb{R}^{p} \ni x \longmapsto(x, 0) \in \mathbb{R}^{n} \tag{4.143}
\end{equation*}
$$

be the inclusion map. We want to think of $u \upharpoonright\{y=0\}$ as $\iota^{*} u$. If $u \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)$ then for any $\phi^{\prime} \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)$ the identity

$$
\begin{equation*}
\iota^{*} u\left(\iota^{*} \phi^{\prime}\right)=u\left(\phi^{\prime} \delta(y)\right) \tag{4.144}
\end{equation*}
$$

holds.
The restriction map $\iota^{*}: \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right) \longrightarrow \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{p}\right)$ is surjective. If $u \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right)$ satisfies the condition

$$
\begin{equation*}
\mathrm{WF}(u) \cap\left\{(x, 0,0, \eta) ; x \in \mathbb{R}^{p}, \eta \in \mathbb{R}^{n-p}\right\}=\emptyset \tag{4.145}
\end{equation*}
$$

then we can interpret the pairing

$$
\begin{array}{r}
\iota^{*} u(\phi)=u\left(\phi^{\prime} \delta(y)\right) \forall \phi \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{p}\right) \\
\text { where } \phi^{\prime} \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right) \text { and } \iota^{*} \phi^{\prime}=\phi \tag{4.146}
\end{array}
$$

to define $\iota^{*} u$. Indeed, the right side makes sense by Proposition $\frac{9.20}{4.11}$.
Thus we have directly proved the first part of
Proposition 4.13. Set $\mathcal{R}=\left\{u \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right) ;\left(\frac{10.21}{4.145}\right)\right.$ holds $\}$ then $\left(\frac{10.20}{4.146)}\right.$ defines a linear restriction map $\iota^{*}: \mathcal{R} \longrightarrow \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{p}\right)$ and
(4.147) $\mathrm{WF}\left(\iota^{*} u\right) \subset\left\{(x, \xi) \in \mathbb{R}^{p} \times\left(\mathbb{R}^{p} \backslash 0\right) ; \exists \eta \in \mathbb{R}^{n}\right.$ with $\left.(x, 0, \xi, \eta) \in \mathrm{WF}(u)\right\}$.

Proof. First note that $\left(\frac{10.21}{4.145}\right)$ means precisely that
(4.148) $\quad \hat{u}(\xi, \eta)$ is rapidly decreasing in a cone around $\{0\} \times \mathbb{R}^{k} \backslash 0$.

When $u \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)$ taking Fourier transforms in $\left(\frac{10.35}{4.144)}\right.$ gives
10.25

$$
\begin{equation*}
\widehat{\iota^{*} u}(\xi)=\frac{1}{(2 \pi)^{k}} \int \hat{u}(\xi, \eta) d \eta \tag{4.149}
\end{equation*}
$$

In general ( $\left(\frac{10}{4.24} .24\right)$ ensures that the integral in $\left(\frac{10}{4} .249\right)$ converges, it will then hold by continuity.

We actually apply $\left(\frac{10.25}{4.149}\right)$ to a localized version of $u$; if $\psi \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{p}\right)$ then

$$
\begin{equation*}
\widehat{\psi \iota^{*}(u)}(\xi)=(2 \pi)^{-k} \int \widehat{\psi}(\xi) \widehat{u}(\xi, \eta) d \eta \tag{4.150}
\end{equation*}
$$

Thus suppose $(\bar{x}, \bar{\xi}) \in \mathbb{R}^{p} \times\left(\mathbb{R}^{p} \backslash 0\right)$ is such that $(\bar{x}, 0, \bar{\xi}, \eta) \notin \mathrm{WF}(u)$ for any $\eta$. If $\psi$ has support close to $\bar{x}$ and $\zeta \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n-p}\right)$ has support close to 0 this means
(4.151) $\widehat{\psi \zeta u}(\xi, \eta)$ is rapidly decreasing in a cone around each $(\bar{\xi}, \eta)$.

We also have rapid decrease around $(0, \eta)$ from $\left(\frac{10.24}{4.148}\right)$ (make sure you understand this point) as

$$
\begin{equation*}
\widehat{\psi \zeta u}(\xi, \eta) \text { is rapidly decreasing in } \gamma \times \mathbb{R}^{p} \tag{4.152}
\end{equation*}
$$

for a cone, $\gamma$, around $\bar{\xi}$. From ( $\left(\frac{10.25}{14.149)}\right.$

$$
\begin{equation*}
\widehat{\psi \iota^{*}(\zeta u)}(\xi) \text { is rapidly decreasing in } \gamma \tag{4.153}
\end{equation*}
$$

Thus $(\bar{x}, \bar{\xi}) \notin \mathrm{WF}\left(\iota^{*}(\zeta u)\right)$. Of course if we choose $\zeta(y)=1$ near $0, \iota^{*}(\zeta u)=\iota^{*}(4)$ so $(\bar{x}, \bar{\xi}) \notin \mathrm{WF}(u)$, provided $(\bar{x}, 0, \bar{\xi}, \eta) \notin \mathrm{WF}(u)$, for all $\eta$. This is what ( $\overline{4.147}$ ) says.

Try to picture what is going on here. We can restate the main conclusion of Proposition $\frac{10}{4.13}$ as follows.

Take $\mathrm{WF}(u) \cap\left\{(x, 0, \xi, \eta) \in \mathbb{R}^{p} \times\{0\} \times\left(\mathbb{R}^{n} \backslash 0\right)\right\}$ and let $Z$ denote projection off the $\eta$ variable:

$$
\begin{equation*}
\mathbb{R}^{p} \times\{0\} \times \mathbb{R}^{p} \times \mathbb{R}^{k} \xrightarrow{Z} \mathbb{R}^{p} \times \mathbb{R}^{p} \tag{4.154}
\end{equation*}
$$

then

$$
\begin{equation*}
\mathrm{WF}\left(\iota^{*} u\right) \subset Z(\mathrm{WF}(u) \cap\{y=0\}) \tag{4.155}
\end{equation*}
$$

We will want to think more about these operations later.

### 4.18. Exterior product

This is maybe the easiest of the elementary operators. It is always defined

$$
\begin{equation*}
\left(u_{1} \boxtimes u_{2}\right)(\phi)=u_{1}\left(u_{2}(\phi(x, \cdot))=u_{2}\left(u_{1}(\phi(\cdot, y))\right.\right. \tag{4.156}
\end{equation*}
$$

Moreover we can easily compute the Fourier transform:

$$
\begin{equation*}
{\widehat{u_{1} \boxtimes u_{2}}}_{2}(\xi, \eta)=\hat{u}_{1}(\xi) \hat{u}_{2}(\eta) \tag{4.157}
\end{equation*}
$$

Proposition 4.14. The (exterior) product

$$
\begin{equation*}
\mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{p}\right) \times \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{k}\right) \longleftarrow \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{p+k}\right) \tag{4.158}
\end{equation*}
$$

is a bilinear map such that

$$
\begin{align*}
\mathrm{WF}\left(u_{1} \boxtimes u_{2}\right) \subset[ & \left.\left(\operatorname{supp}\left(u_{1}\right) \times\{0\}\right) \times \mathrm{WF}\left(u_{2}\right)\right]  \tag{4.159}\\
& \cup\left[\mathrm{WF}\left(u_{1}\right) \times\left(\operatorname{supp}\left(u_{2}\right) \times\{0\}\right)\right] \cup\left[\mathrm{WF}\left(u_{1}\right) \times \operatorname{WF}\left(u_{2}\right)\right] .
\end{align*}
$$

Proof. We can localize near any point $(\bar{x}, \bar{y})$ with $\phi_{1}(x) \phi_{2}(y)$, where $\phi_{1}$ is supported near $\bar{x}$ and $\phi_{2}$ is supported near $\bar{y}$. Thus we only need examine the decay of

$$
\phi_{1} \widehat{u_{1} \boxtimes \phi_{2} u_{2}}=\widehat{\phi_{1} u_{1}}(\xi) \cdot \widehat{\phi_{2} u_{2}}(\eta)
$$

Notice that if $\widehat{\phi_{1} u_{1}}(\xi)$ is rapidly decreasing around $\bar{\xi} \neq 0$ then the product is rapidly decreasing around any $(\bar{\xi}, \eta)$. This gives (14.33 14.159 .

### 4.19. Diffeomorphisms

We next turn to the question of the extension of $F^{*}$, where $F: \Omega_{1} \longrightarrow \Omega_{2}$ is a $\mathcal{C}^{\infty}$ map, from $\mathcal{C}^{\infty}\left(\Omega_{2}\right)$ to some elements of $\mathcal{C}^{-\infty}\left(\Omega_{2}\right)$. The simplest example of pull-back is that of transformation by a diffeomorphism.

We have already noted how pseudodifferential operators behave under a diffeomorphism: $F: \Omega_{1} \longrightarrow \Omega_{2}$ between open sets of $\mathbb{R}^{n}$. Suppose $A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ has Schwartz kernel of compact support in $\Omega_{1} \times \Omega_{1}$ then we define

$$
\begin{equation*}
A_{F}: \mathcal{C}_{c}^{\infty}\left(\Omega_{2}\right) \longrightarrow \mathcal{C}_{c}^{\infty}\left(\Omega_{2}\right) \tag{4.161}
\end{equation*}
$$

by $A_{F}=G^{*} \cdot A \cdot F^{*}, G=F^{-1}$. In $\S \frac{\text { Sect. CooInv }}{4.4 \text { we showed that } A_{F} \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right) \text {. In fact we }}$ showed much more, namely we computed a (very complicated) formula for the full symbols. Recall the definition of the cotangent bundle of $\mathbb{R}^{n}$

$$
\begin{equation*}
T^{*} \mathbb{R}^{n} \simeq \mathbb{R}^{n} \times \mathbb{R}^{n} \tag{4.162}
\end{equation*}
$$

identified as pairs of points $(\bar{x}, \bar{\xi})$, where $\bar{x} \in \mathbb{R}^{n}$ and

$$
\begin{equation*}
\bar{\xi}=d f(\bar{x}) \text { for some } f \in \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right) \tag{4.163}
\end{equation*}
$$

The differential $d f(\bar{x})$ of $f$ at $\bar{x} \in \mathbb{R}^{n}$ is just the equivalence class of $f(x)-f(\bar{x}) \in \mathcal{I}_{\bar{x}}$ modulo $\mathcal{I}_{\bar{x}}^{2}$. Here

$$
\left\{\begin{array}{l}
\mathcal{I}_{\bar{x}}=\left\{g \in \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right) ; g(\bar{x})=0\right\}  \tag{4.164}\\
\mathcal{I}_{\bar{x}}^{2}=\left\{\sum_{\text {finite }} g_{i} h_{i}, g_{i}, h_{i} \in \mathcal{I}_{\bar{x}}\right\} .
\end{array}\right.
$$

The identification of $\bar{\xi}$, given by ( $\left(\frac{11.2}{1.162)}\right.$ ) and $\left(\frac{11.3}{4.16} 3\right)$, with a point in $\mathbb{R}^{n}$ is obtained using Taylor's formula. Thus if $f \in \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right)$

$$
\begin{equation*}
f(x)=f(\bar{x})+\sum_{i=1}^{n} \frac{\partial f}{\partial x_{j}}(\bar{x})(x-\bar{x})_{j}+\sum_{i, j=1} g_{i j}(x) x_{i} x_{j} \tag{4.165}
\end{equation*}
$$

The double sum here is in $\mathcal{I}_{\bar{x}}^{2}$, so the residue class of $f(x)-f(\bar{x})$ in $\mathcal{I}_{\bar{x}} / \mathcal{I}_{\bar{x}}^{2}$ is the same as that of

$$
\begin{equation*}
\sum_{i=1}^{n} \frac{\partial f}{\partial x_{j}}(\bar{x})(x-\bar{x})_{j} \tag{4.166}
\end{equation*}
$$

That is, $d(x-\bar{x})_{j}=d x_{j}, j=1, \ldots, n$ form a basis for $T_{\bar{x}}^{*} \mathbb{R}^{n}$ and in terms of this basis

$$
\begin{equation*}
d f(\bar{x})=\sum_{i=1}^{n} \frac{\partial f}{\partial x_{j}}(\bar{x}) d x_{j} \tag{4.167}
\end{equation*}
$$

Thus the entries of $\bar{\xi}$ are just $\left(\frac{\partial f}{\partial x_{1}}, \ldots \frac{\partial f}{\partial x_{n}}\right)$ for some $f$. Another way of saying this is that the linear functions $\xi \cdot x=\xi_{1} x_{1}+\xi_{2} x_{2} \cdots \xi_{n} x_{n}$ have differentials spanning $T_{x}^{*} \mathbb{R}^{n}$.

So suppose $F: \Omega_{1} \longrightarrow \Omega_{2}$ is a $\mathcal{C}^{\infty}$ map. Then

$$
\begin{equation*}
F^{*}: T_{\bar{y}}^{*} \Omega_{2} \longrightarrow T_{\bar{x}}^{*} \Omega_{1}, \bar{y}=F(\bar{x}) \tag{4.168}
\end{equation*}
$$

is defined by $F^{*} d f(\bar{y})=d\left(F^{*} f\right)(\bar{x})$ since $F^{*}: \mathcal{I}_{\bar{y}} \longrightarrow \mathcal{I}_{\bar{x}}, F^{*}: \mathcal{I}_{\bar{y}}^{2} \longrightarrow \mathcal{I}_{\bar{x}}^{2}$. In coordinates $F(x)=y \Longrightarrow$

$$
\begin{equation*}
\frac{\partial}{\partial x_{j}}\left(F^{*} f(x)\right)=\frac{\partial}{\partial y} f(F(x))=\sum_{k=1}^{n} \frac{\partial f}{\partial x_{k}}(y) \frac{\partial F_{k}}{\partial x_{j}} \tag{4.169}
\end{equation*}
$$

i.e. $F^{*}(\eta \cdot d y)=\xi \cdot d x$ if

$$
\begin{equation*}
\xi_{j}=\sum_{k=1}^{n} \frac{\partial F_{k}}{\partial x_{j}}(x) \cdot \eta_{k} \tag{4.170}
\end{equation*}
$$

Qf course if $F$ is a diffeomorphism then the Jacobian matrix $\frac{\partial F}{\partial x}$ is invertible and $\left(\frac{1.170}{4.170)}\right.$ is a linear isomorphism. In this case

$$
\begin{gather*}
F^{*}: T_{\Omega_{2}}^{*} \mathbb{R}^{n} \longleftrightarrow T_{\Omega_{1}}^{*} \mathbb{R}^{n}  \tag{4.171}\\
(x, \xi) \longleftrightarrow(F(x), \eta)
\end{gather*}
$$

with $\xi$ and $\eta$ connected by $\left(\frac{11.6}{4.17} 0\right)$. Thus $\left(F^{*}\right)^{*}: \mathcal{C}^{\infty}\left(T^{*} \Omega_{1}\right) \longrightarrow \mathcal{C}^{\infty}\left(T^{*} \Omega_{2}\right)$.
Proposition 4.15. If $F: \Omega_{1} \longrightarrow \Omega_{2}$ is a diffeomorphism of open sets of $\mathbb{R}^{n}$ and $A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ has Schwartz kernel with compact support in $\Omega_{1} \times \Omega_{2}$ then

$$
\begin{equation*}
\sigma_{m}\left(A_{F}\right)=\left(F^{*}\right)^{*} \sigma_{m}(A) \tag{4.172}
\end{equation*}
$$

and
11.10

$$
\begin{equation*}
F^{*}\left(\mathrm{WF}^{\prime}\left(A_{F}\right)\right)=\mathrm{WF}^{\prime}(A) \tag{4.173}
\end{equation*}
$$

It follows that symbol $\sigma_{m}(A)$ of $A$ is well-defined as an element of $S_{\infty}^{m-[1]}\left(T^{*} \mathbb{R}^{n}\right)$ independent of coordinates and $\mathrm{WF}^{\prime}(A) \subset T^{*} \mathbb{R}^{n} \backslash 0$ is a well-defined closed conic set, independent of coordinates. The elliptic set and the characteristic set $\Sigma_{m}$ are therefore also well-defined complementary conic subsets of $T^{*} \Omega \backslash 0$.

Proof. Look at the formulae.
The main use we make of this invariance result is the freedom it gives us to choose local coordinates adapted to a particular problem. It also suggests that there should be neater ways to write various formulae, e.g. the wavefront sets of push-forward and pull-backs.

Proposition 4.16. If $u \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right)$ has $\operatorname{supp}(u) \subset \Omega_{2}$ and $F: \Omega_{1} \longrightarrow \Omega_{2}$ is a diffeomorphism then
11.13 (4.174)

$$
\mathrm{WF}\left(F^{*} u\right) \subset\left\{(x, \xi) \in \mathbb{R}^{n} \times\left(\mathbb{R}^{n} \backslash 0\right) ;(F(x), \eta) \in \mathrm{WF}(u), \eta_{j}=\sum_{i} \frac{\partial F_{i}}{\partial x_{j}}(x) \xi_{i}\right\}
$$

Proof. Just use the standard definition

$$
\begin{equation*}
\mathrm{WF}\left(F^{*} u\right)=\bigcap\left\{\Sigma(A) ; A\left(F^{*} u\right) \in \mathcal{C}^{\infty}\right\} \tag{4.175}
\end{equation*}
$$

To test the wavefront set of $F^{*} u$ it suffices to consider $A$ 's with kernels supported in $\Omega_{1} \times \Omega_{1}$ since $\operatorname{supp}\left(F^{*} u \Subset \Omega_{1}\right.$ and for a general pseudodifferential operator $A^{\prime}$
there exists $A$ with kernel supported in $\Omega_{1}$ such that $A^{\prime} u-A u \in \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right)$. Then $A F^{*} u \in \mathcal{C}_{c}^{\infty}\left(\Omega_{1}\right) \Longleftrightarrow A_{F} u \in \mathcal{C}_{c}^{\infty}\left(\Omega_{2}\right)$. Thus

$$
\begin{align*}
\mathrm{WF}\left(F^{*} u\right) & =\bigcap\left\{\Sigma(A) ; A_{F} u \in \mathcal{C}^{\infty}\right\}  \tag{4.176}\\
& =\bigcap\left\{F^{*}\left(\Sigma\left(A_{F}\right)\right) ; A_{F} u \in \mathcal{C}^{\infty}\right\}  \tag{4.177}\\
& =F^{*} \mathrm{WF}(u) \tag{4.178}
\end{align*}
$$

since, for $u$, it is enough to consider operators with kernels supported in $\Omega_{2} \times \Omega_{2}$.

### 4.20. Products

Although we have discussed the definition of the product of two distributions we have not yet analyzed the wavefront set of the result.

Proposition 4.17. If $u_{1}, u_{2} \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right)$ are such that

$$
\begin{equation*}
(x, \xi) \in \mathrm{WF}\left(u_{1}\right) \Longrightarrow(x,-\xi) \notin \mathrm{WF}\left(u_{2}\right) \tag{4.179}
\end{equation*}
$$

then the product $u_{1} u_{2} \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right)$, defined by Proposition $\frac{9.20}{4.11}$ satisfies

$$
\begin{align*}
\mathrm{WF}\left(u_{1} u_{2}\right) \subset & \left\{(x, \xi) ; x \in \operatorname{supp}\left(u_{1}\right) \text { and }(x, \xi) \in \mathrm{WF}\left(u_{2}\right)\right\} \\
& \cup\left\{(x, \xi) ; x \in \operatorname{supp}\left(u_{2}\right) \text { and }(x, \xi) \in \mathrm{WF}\left(u_{1}\right)\right\}  \tag{4.180}\\
& \cup\left\{(x, \xi) ; \xi=\eta_{1}+\eta_{2},\left(x, \eta_{i}\right) \in \mathrm{WF}\left(u_{i}\right), i=1,2\right\} .
\end{align*}
$$

Proof. We can represent the product in terms of three 'elementary' operations.

$$
\begin{equation*}
u_{1} u_{2}(x)=\iota^{*}\left[F^{*}\left(u_{1} \boxtimes u_{2}\right)\right] \tag{4.181}
\end{equation*}
$$

where $F: \mathbb{R}^{2 n} \longrightarrow \mathbb{R}^{2 n}$ is the linear transformation

$$
\begin{equation*}
F(x, y)=(x+y, x-y) \tag{4.182}
\end{equation*}
$$

and $\iota: \mathbb{R}^{n} \hookrightarrow \mathbb{R}^{n} \times\{0\} \subset \mathbb{R}^{2 n}$ is inclusion as the first factor. Thus (11.151) corresponds to the 'informal' notation

$$
\begin{equation*}
u_{1} u_{2}(x)=u_{1}(x+y) u_{2}(x-y) \upharpoonright\{y=0\} \tag{4.183}
\end{equation*}
$$

and will follow by continuity once. we analyse the wavefront set properties.
We know from Proposition 4.14 that

$$
\begin{gather*}
\mathrm{WF}\left(u_{1} \boxtimes u_{2}\right) \subset\left\{(X, Y, \Xi, H) ; X \in \operatorname{supp}\left(u_{1}\right), \Xi=0,(Y, H) \in \mathrm{WF}\left(u_{2}\right)\right\} \\
\cup\left\{(X, Y, \Xi, H) ;(X, \Xi) \in \mathrm{WF}\left(u_{1}\right), Y \in \operatorname{supp}\left(u_{2}\right), H=0\right\}  \tag{4.184}\\
\cup\left\{(X, Y, \Xi, H) ;(X, \Xi) \in \operatorname{WF}\left(u_{1}\right),(Y, H) \in \mathrm{WF}\left(u_{2}\right)\right\} .
\end{gather*}
$$

Since $F$ is a diffeomorphism, by Proposition $\frac{11.12}{1.16,}$

$$
\begin{aligned}
& \mathrm{WF}\left(F^{*}\left(u_{1} \boxtimes u_{2}\right)\right)=\left\{(x, y, \xi, \eta) ;\left(F^{t}(x, y), \Xi, H\right) \in \mathrm{WF}\left(u_{1} \boxtimes u_{2}\right),\right. \\
&\left.(\xi, \eta)=A^{t}(\Xi, H)\right\} .
\end{aligned}
$$

where $F^{t}$ is the transpose of $F$ as a linear map. In fact $F^{t}=F$, so

$$
\begin{aligned}
& \mathrm{WF}\left(F^{*}\left(u_{1} \boxtimes u_{2}\right)\right) \subset \\
& \qquad\left\{(x, y, \xi, \eta) ; x+y \in \operatorname{supp}\left(u_{1}\right), \xi+\eta=0,\left(x-y, \frac{1}{2}(\xi-\eta)\right) \in \mathrm{WF}\left(u_{2}\right)\right\} \\
& \quad \cup\left\{(x, y, \xi, \eta) ;\left(x+y, \frac{1}{2}(\xi+\eta)\right) \in \mathrm{WF}\left(u_{1}\right),\left(x-y, \frac{1}{2}(\xi-\eta)\right) \in \mathrm{WF}\left(u_{2}\right)\right\}
\end{aligned}
$$

and so using Proposition $\frac{10.22}{4.13}$

$$
\begin{aligned}
\operatorname{WF}\left(F^{*}\left(u_{1} \boxtimes u_{2}\right)\right) & \upharpoonright\{y=0\} \\
& \subset\left\{(x, 0, \xi,-\xi) ; x \in \operatorname{supp}\left(u_{1}\right),(x, \xi) \in \mathrm{WF}\left(u_{2}\right)\right\} \\
& \cup\left\{(x, 0, \xi, \eta) ;\left(x \in \operatorname{supp}\left(u_{2}\right),(x, \xi) \in \mathrm{WF}\left(u_{2}\right)\right\}\right. \\
\cup & \left\{(x, 0, \xi, \eta) ;\left(x, \frac{1}{2}(\xi+\eta)\right) \in \mathrm{WF}\left(u_{2}\right),\left(x, \frac{1}{2}(\xi-\eta)\right) \in \mathrm{WF}\left(u_{1}\right)\right\}
\end{aligned}
$$

Notice that

$$
\begin{equation*}
(x, 0,0, \eta) \in \mathrm{WF}\left(F^{*}\left(u_{1} \boxtimes u_{2}\right)\right) \Longrightarrow\left(x, \frac{1}{2} \eta\right) \in \mathrm{WF}\left(u_{1}\right) \text { and }\left(x, \frac{1}{2} \eta\right) \mathrm{WF}\left(u_{2}\right) \tag{4.185}
\end{equation*}
$$

which introduces the assumption under which $u_{1} u_{2}$ is defined. Finally then we see that
11.17 (4.186)

$$
\begin{aligned}
\mathrm{WF}\left(u_{1} u_{2}\right) \subset\{ & \left\{(x, \xi) ; x \in \operatorname{supp}\left(u_{1}\right),(x, \xi) \in \mathrm{WF}\left(u_{2}\right)\right\} \\
& \cup\left\{(x, \xi) ; x \in \operatorname{supp}\left(u_{2}\right),(x, \xi) \in \mathrm{WF}\left(u_{1}\right)\right\} \\
& \cup\left\{(x, \xi) ;\left(x, \eta_{1}\right) \in \mathrm{WF}\left(u_{1}\right),\left(x, \eta_{2}\right) \in \mathrm{WF}\left(u_{2}\right) \text { and } \xi=\eta_{1}+\eta_{2}\right\} .
\end{aligned}
$$

which is another way of writing the conclusion of Proposition $\frac{11.14}{1.17 .}$

### 4.21. Pull-back

Now let us consider a general $\mathcal{C}^{\infty}$ map

$$
F: \Omega_{1} \longrightarrow \Omega_{2}, \Omega_{1} \subset \mathbb{R}^{n}, \Omega_{2} \subset \mathbb{R}^{m}
$$

Thus even the dimension of domain and range spaces can be different. When can we define $F^{*} u$, for $u \in \mathcal{C}_{c}^{-\infty}\left(\Omega_{2}\right)$ and what can we say about $\mathrm{WF}\left(F^{*} u\right)$ ? For a general map $F$ it is not possible to give a sensible, i.e. consistent, definition of $F^{*} u$ for all distributions $u \in \mathcal{C}^{-\infty}\left(\Omega_{2}\right)$.

For smooth functions we have defined

$$
\begin{equation*}
F^{*}: \mathcal{C}_{c}^{\infty}\left(\Omega_{2}\right) \longrightarrow \mathcal{C}^{\infty}\left(\Omega_{1}\right) \tag{4.188}
\end{equation*}
$$

but in general $F^{*} \phi$ does not have compact support, even if $\phi$ does. We therefore impose the condition that $F$ be proper

$$
\begin{equation*}
F^{-1}(K) \Subset \Omega_{2} \forall K \Subset \Omega_{2} \tag{4.189}
\end{equation*}
$$

(mostly just for convenience). In fact if we want to understand $F^{*} u$ near $\bar{x}_{1} \in \Omega_{1}$ we only need to consider $u$ near $F\left(\bar{x}_{1}\right)_{18} \in \Omega_{2}$.

The problem is that the map ( 4.187 ) may be rather complicated. However any smooth map can be decomposed into a product of simpler maps, which we can analyze locally. Set

$$
\begin{equation*}
\operatorname{graph}(F)=\left\{(x, y) \in \Omega_{1} \times \Omega_{2} ; y=F(x)\right\} \xrightarrow{\iota_{F}} \Omega_{1} \times \Omega_{2} \tag{4.190}
\end{equation*}
$$

This is always an embedded submanifold of $\Omega_{1} \times \Omega_{2}$ the functions $y_{i}-F_{i}(x)$, $i=1, \ldots, N$ are independent defining functions for $\operatorname{graph}(F)$ and $x_{1}, \ldots, x_{n}$ are coordinates on it. Now we can write

### 11.20

$$
\begin{equation*}
F=\pi_{2} \circ \iota_{F} \circ g \tag{4.191}
\end{equation*}
$$

where $g: \Omega_{1} \longleftrightarrow \operatorname{graph}(F)$ is the diffeomorphism onto its range $x \longmapsto(x, F(x))$. This decomposes $F$ as a projection, an inclusion and a diffeomorphism. Now consider

$$
\begin{equation*}
F^{*} \phi=g^{*} \cdot \iota_{F}^{*} \cdot \pi_{2}^{*} \phi \tag{4.192}
\end{equation*}
$$

i.e. $F^{*} \phi$ is obtained by pulling $\phi$ back from $\Omega_{2}$ to $\Omega_{1} \times \Omega_{2}$, restricting to $\operatorname{graph}(F)$ and then introducing the $x_{i}$ as coordinates. We have directly discussed ( $\pi_{2}^{*} \phi$ ) but we can actually write it as

$$
\begin{equation*}
\pi_{2}^{*} \phi=1 \boxtimes \phi(y) \tag{4.193}
\end{equation*}
$$

so the result we have proved can be applied to it. So let us see what writing ( $\frac{11.21}{4.192}$ ) as

$$
\begin{equation*}
F^{*} \phi=g^{*} \circ \iota_{F}^{*}(1 \boxtimes \phi) \tag{4.194}
\end{equation*}
$$

tells us. If $u \in \mathcal{C}_{c}^{-\infty}\left(\Omega_{2}\right)$ then
11.24 (4.195)

$$
\underset{21}{\mathrm{WF}}(1 \boxtimes u) \subset\{(x, y, 0, \eta) ;(y, \eta) \in \mathrm{WF}(u)\}
$$

by Proposition $\frac{10.31}{4.14}$. So we have to discuss $\iota_{F}^{*}(1 \boxtimes u)$, i.e. restriction to $y=F(x)$. We can do this by making a diffeomorphism:

$$
\begin{equation*}
T_{F}(x, y)=(x, y+F(x)) \tag{4.196}
\end{equation*}
$$

so that $T_{F}^{-1}(\operatorname{graph}(F))=\{(x, 0)\}$. Notice that $g \circ T_{F}=\pi_{1}$, so
11.26

$$
\begin{equation*}
F^{*} \phi=\iota_{\{y=0\}}^{*}\left(T_{F}^{*}(1 \boxtimes u)\right) \tag{4.197}
\end{equation*}
$$

Now from Proposition $\frac{11.12}{4.16}$ we know that
$\qquad$

$$
\begin{align*}
& \mathrm{WF}\left(T_{F}^{*}(1 \boxtimes u)\right)=T_{F}^{*}(\mathrm{WF}(1 \boxtimes u))  \tag{4.198}\\
&=\{(X, Y, \Xi, H) ;(X, Y+F(X), \xi, \eta) \in \mathrm{WF}(1 \boxtimes u) \\
& \eta\left.=H, \xi_{i}=\Xi_{i}+\Sigma \frac{\partial F_{j}}{\partial x_{i}} H_{j}\right\}
\end{align*}
$$

i.e.
11.28
(4.199) $\mathrm{WF}\left(T_{F}^{*}(1 \boxtimes u)\right)=\left\{(x, y, \xi, \eta) ; \xi_{i}=\sum_{j} \frac{\partial F_{j}}{\partial x_{j}}(x) \eta_{j},(F(x), \eta) \in \mathrm{WF}(u)\right\}$.

So consider our existence condition for restriction to $y=0$, that $\xi \neq 0$ on $\mathrm{WF}\left(T_{F}^{*}(1 \boxtimes\right.$ u)) i.e.

$$
11.29
$$

(4.200)

$$
(F(x), \eta) \in \mathrm{WF}(u) \Longrightarrow \sum_{j} \frac{\partial F_{j}}{\partial x_{i}}(x) \eta_{j} \neq 0
$$

If $\left(\frac{11.29}{4.200}\right)$ holds then, from $\left(\frac{11.27}{4.198}\right)$ and Proposition $\frac{10.22}{1.13}$
11.30 (4.201)

$$
\mathrm{WF}\left(F^{*} u\right) \subset\left\{(x, \xi) ; \exists(F(x), \eta) \in \mathrm{WF}(u) \text { and } \xi_{j}=\sum_{j} \frac{\partial F_{j}}{\partial x_{i}}(x) \eta_{j}\right\}
$$

We can reinterpret $\left(\frac{11.29}{4.200}\right)$ and $\left(\frac{11.30}{4.201}\right)$ more geometrically. The differential of $F$ gives a map

$$
\begin{gather*}
F^{*}: T_{F(x)}^{*} \Omega_{2} \longrightarrow T_{x}^{*} \Omega_{1} \forall x \in \Omega_{1} \\
(F(x), \eta) \longmapsto(x, \xi) \text { where } \xi_{i}=\Sigma \frac{\partial F_{j}}{\partial x_{i}} \eta_{j} \tag{4.202}
\end{gather*}
$$

Thus $\left(\frac{11.29}{4.200}\right)$ can be restated as:

$$
\begin{equation*}
\forall x \in \Omega_{1}, \text { the null space of } F_{x}^{*}: T_{F(x)}^{*} \Omega_{2} \longrightarrow T_{x}^{*} \Omega_{1} \tag{4.203}
\end{equation*}
$$

and then $\left(\frac{11.30}{4.201}\right)$ becomes
11.32

$$
\begin{equation*}
\mathrm{WF}\left(F^{*} u\right) \subset \bigcup_{x \in \Omega_{1}} F_{x}^{*}\left[\mathrm{WF}(u) \cap T_{F(x)}^{*} \Omega_{2}\right]=F^{*}(\mathrm{WF}(u)) \tag{4.204}
\end{equation*}
$$

(proved we are a little careful in that $F^{*}$ is not a map; it is a relation between $T^{*} \Omega_{2}$ and $\left.T^{*} \Omega_{1}\right)$ and in this sense ( $(4.203)$ holds. Notice that ( 14.201 ) is a sensible "consequence" of ( 4.203 ), since otherwise $\mathrm{WF}\left(F^{*} u\right)$ would contain some zero directions.
11.33 Proposition 4.18. If $F: \Omega_{1} \longrightarrow \Omega_{2}$ is a proper $\mathcal{C}^{\infty}$ map then $F^{*}$ extends (by continuity) from $\mathcal{C}_{c}^{\infty}\left(\Omega_{2}\right)$ to
11.34 (4.205) $\left\{u \in \mathcal{C}_{c}^{-\infty}\left(\Omega_{2}\right) ; F^{*}(\mathrm{WF}(u)) \cap\left(\Omega_{1} \times 0\right)=\emptyset\right.$ in $\left.T^{*} \Omega_{1}\right\}$
and $\left(\frac{11.32}{4.204}\right)$ holds.

### 4.22. The operation $F_{*}$

Next we will look at the dual operation, that of push-forward. Notice the basic properties of pull-back:

| 12.1 |
| :--- |
| 12.2 |

Maps $\mathcal{C}_{c}^{\infty}$ to $\mathcal{C}_{c}^{\infty}$ (if $F$ is proper)
Not always defined on distributions.
Dually we find
Proposition 4.19. If $F: \Omega_{1} \longrightarrow \Omega_{2}$ is a $\mathcal{C}^{\infty}$ map of an open subset of $\mathbb{R}^{n}$ into an open subset of $\mathbb{R}^{n}$ then for any $u \in \mathcal{C}_{c}^{-\infty}\left(\Omega_{1}\right)$

$$
\begin{equation*}
F_{*}(u)(\phi)=u\left(F^{*} \phi\right) \tag{4.208}
\end{equation*}
$$

is a distribution of compact support and

$$
\begin{equation*}
F_{*}: \mathcal{C}_{c}^{-\infty}\left(\Omega_{1}\right) \longrightarrow \mathcal{C}_{c}^{-\infty}\left(\Omega_{2}\right) \tag{4.209}
\end{equation*}
$$

has the property:

## 12.6

$$
\mathrm{WF}\left(F_{*} u\right) \subset\left\{(y, \eta) ; y \in F(\operatorname{supp}(u)), y=F(x), F_{x}^{*} \eta=0\right\} \cup
$$

$$
\begin{equation*}
\left\{(y, \eta) ; y=F(x),\left(x, F_{x}^{*} \eta\right) \in \mathrm{WF}(u)\right\} \tag{4.210}
\end{equation*}
$$

Proof. Notice that the 'opposite' of $\left(\frac{12.1}{4.206}\right)$ and $\left(\frac{12.2}{4.207}\right)$ hold, i.e. $F_{*}$ is always defined but even if $u \in \mathcal{C}_{c}^{\infty}\left(\Omega_{1}\right)$ in general $F_{*} u \notin \mathcal{C}_{c}^{\infty}\left(\Omega_{2}\right)$. All we really have to prove is ( $\sqrt{4.210} 0)$. As usual we look for a formula in terms of elementary operations. So suppose $u \in \mathcal{C}_{c}^{\infty}\left(\Omega_{1}\right)$

$$
\begin{align*}
F_{*} u(\phi) & =u\left(F^{*} \phi\right) \quad \phi \in \mathcal{C}_{c}^{\infty}\left(\Omega_{2}\right) \\
& =\int u(x) \phi(F(x)) d x  \tag{4.211}\\
& =\int u(x) \delta(y-F(x)) \phi(y) d y d x
\end{align*}
$$

Thus, we see that

$$
\begin{equation*}
F_{*} u=\pi_{*} H^{*}(u \boxtimes \delta) \tag{4.212}
\end{equation*}
$$

where $\delta=\delta(y) \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{m}\right), H$ is the diffeomorphism

$$
\begin{equation*}
H(x, y)=(x, y-F(x)) \tag{4.213}
\end{equation*}
$$

and $\pi: \mathbb{R}^{n+m} \longrightarrow \mathbb{R}^{m}$ is projection off the first factor.
Thus ( $\frac{12.8}{4.212)}$ ) is the desired decomposition into elementary operations, since $\left.u \boxtimes \delta \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n+m}\right), \pi_{*} H_{12 .{ }^{*}}{ }^{*} u \boxtimes \delta\right)$ is always defined and indeed the map is continuous, which actually proves ( 4.212 ).

So all we need ${ }^{\text {to }} 31$ do is estimate the wavefront set using our earlier results. From Proposition $\sqrt{4.14}$ it follows that

$$
\begin{align*}
& \mathrm{WF}(u \boxtimes \delta) \subset\subset(x, 0, \xi, \eta) ; x \in \operatorname{supp}(u), \xi=0\} \cup\{(x, 0, \xi, 0) ;(x, \xi) \in \mathrm{WF}(u)\}  \tag{4.214}\\
& \cup\{(x, 0, \xi, \eta) ;(x, \xi) \in \mathrm{WF}(u)\} \\
&=\{(x, 0, \xi, \eta) ; x \in \operatorname{supp}(u), \xi=0\} \cup\{(x, 0, \xi, \eta) ;(x, \xi) \in \mathrm{WF}(u)\} .
\end{align*}
$$

Then consider what happens under $H^{*}$. This is a diffeomorphism so the wavefront set transforms under the pull-back:

$$
\begin{align*}
& \operatorname{WF}\left(H^{*}(u \boxtimes \delta)\right)=\mathrm{WF}(u(x) \delta(y-F(x)) \\
& =\left\{(x, F(x), \Xi, \eta) ; \Xi_{i}=\xi_{i}-\sum_{j} \frac{\partial F_{j}}{\partial x_{i}}(x) \eta_{j},(x, 0, \xi, \eta) \in \mathrm{WF}(u \boxtimes \delta)\right\} \\
& \left.=\left\{(x, F(x), \Xi, \eta) ; x \in \operatorname{supp}(u), \Xi_{i}=-\sum_{j} \frac{\partial F_{j}}{\partial x_{i}}(x) \eta_{j}\right)\right\}  \tag{4.215}\\
& \\
& \left.\quad \cup\left\{(x, F(x), \Xi, \eta) ; \eta \in \mathbb{R}^{m},(x, \xi) \in \mathrm{WF}(u)\right), \Xi_{i}=\xi_{i}-\sum_{j} \frac{\partial F_{i}}{\partial x_{j}} \eta_{j}\right\} .
\end{align*}
$$

Finally recall the behaviour of wavefront sets under projection, to see that

$$
\begin{aligned}
& \quad \mathrm{WF}\left(F_{*} u\right) \subset\left\{(y, \eta) ; \exists(x, y, 0, \eta) \in \mathrm{WF}\left(H^{*}(u \boxtimes \delta)\right)\right\} \\
& =\{(y, \eta) ; y=F(x) \text { for some } x \in \operatorname{supp}(u) \text { and } \\
& \left.\qquad \sum_{j} \frac{\partial F_{j}}{\partial x_{i}} \eta_{j}=0, i=1, \ldots, n\right\} \\
& \cup\{(y, \eta) ; y=F(x) \text { for some }(x, \xi) \in \mathrm{WF}(u) \text { and } \\
& \left.\left.\qquad \xi_{i}=\sum_{j} \frac{\partial F_{i}}{\partial x_{i}} \eta_{j}, i=1, \ldots, n\right)\right\} .
\end{aligned}
$$

This says

$$
\begin{align*}
\mathrm{WF}\left(F_{*} u\right) \subset\{ & \left.(y, \eta) ; y \in F(\operatorname{supp}(u)) \text { and } F_{x}^{*}(\eta)=0\right\}  \tag{4.216}\\
& \cup\left\{(y, \eta) ; y=F(x) \text { with }\left(x, F_{x}^{*} \eta\right) \in \mathrm{WF}(u)\right\} \tag{4.217}
\end{align*}
$$

which is just ( $\left.\frac{12.6}{4.210} 0\right)$.
As usual one should note that the two terms here are "really the same".
Now let us look at $F_{*}$ as a linear map,

$$
\begin{equation*}
F_{*}: \mathcal{C}_{c}^{\infty}\left(\Omega_{1}\right) \longrightarrow \mathcal{C}_{c}^{-\infty}\left(\Omega_{2}\right) \tag{4.218}
\end{equation*}
$$

As such it has a Schwartz kernel, indeed ( $\left(\frac{12.8}{4.212}\right)$ is just the usual formula for an operator in terms of its kernel:
12.10

$$
\begin{equation*}
F_{*} u(y)=\int K(y, x) u(x) d x, K(y, x)=\delta(y-F(x)) \tag{4.219}
\end{equation*}
$$

So consider the wavefront set of the kernel:
(4.220) $\quad \mathrm{WF}\left(\delta(y-F(x))=\mathrm{WF}\left(H^{*} \delta(y)\right)=\left\{(y, x ; \eta, \xi) ; y=F(x), \xi=F_{x}^{*} \eta\right\}\right.$.

Now changing the order of the factors we can regard this as a subset

> | 12.12 |
| :--- |

(4.221) $\mathrm{WF}^{\prime}(K)=\left\{((y, \eta),(x, \xi)) ; y=F(x), \xi=F^{*} \eta\right\} \subset\left(\Omega_{2} \times \mathbb{R}^{m}\right) \times\left(\Omega_{1} \times \mathbb{R}^{n}\right)$.

As a subset of the product we can regard $\mathrm{WF}^{\prime}(K)$ as a relation: if $\Gamma \subset \Omega_{2} \times$ ( $\mathbb{R}^{n} \backslash 0$ ) set

$$
\begin{gathered}
\mathrm{WF}^{\prime}(K) \circ \Gamma= \\
\left.\left\{(y, \eta) \in \Omega_{2} \times\left(\mathbb{R}^{m} \backslash 0\right) ; \exists((y, \eta)),(x, \xi)\right) \in \mathrm{WF}^{\prime}(K) \text { and }(x, \xi) \in \Gamma\right\}
\end{gathered}
$$

Indeed with this definition

$$
\begin{equation*}
\mathrm{WF}\left(F_{*} u\right) \subset \mathrm{WF}^{\prime}(K) \circ \mathrm{WF}(u), \quad K=\text { kernel of } F_{*} \tag{4.222}
\end{equation*}
$$

### 4.23. Wavefront relation

One serious application of our results to date is:
12.15 THEOREM 4.1. Suppose $\Omega_{1} \subset \mathbb{R}^{n}, \Omega_{2} \subset \mathbb{R}^{m}$ are open and $A \in \mathcal{C}^{-\infty}\left(\Omega_{1} \times \Omega_{2}\right)$ has proper support, in the sense that the two projections

are proper, then $A$ defines a linear map

$$
\begin{equation*}
A: \mathcal{C}_{c}^{\infty}\left(\Omega_{2}\right) \longrightarrow \mathcal{C}_{c}^{-\infty}\left(\Omega_{1}\right) \tag{4.224}
\end{equation*}
$$

and extends by continuity to a linear map

$$
\begin{align*}
A:\left\{u \in \mathcal{C}_{c}^{-\infty}(X)\right. & ; \mathrm{WF}(u) \cap\left\{(y, \eta) \in \Omega_{2} \times\left(\mathbb{R}^{n} \backslash 0\right) ;\right.  \tag{4.225}\\
& \exists(x, 0, y,-\eta) \in \mathrm{WF}(K)\}=\emptyset\} \longrightarrow \mathcal{C}_{c}^{-\infty}\left(\Omega_{1}\right) \tag{4.226}
\end{align*}
$$

for which
12.19
(4.227)

$$
\mathrm{WF}(A u) \subset \mathrm{WF}^{\prime}(A) \circ \mathrm{WF}(u)
$$

where

> | 12.20 |
| :---: |

(4.228)

$$
\begin{aligned}
& \mathrm{WF}^{\prime}(A)=\left\{((x, \xi),(y, \eta)) \in\left(\Omega_{1} \times \mathbb{R}^{n}\right) \times\left(\Omega_{2} \times \mathbb{R}^{m}\right) ;(\xi, \eta) \neq 0\right. \\
&\text { and }(x, y, \xi,-\eta) \in \mathrm{WF}(K)\}
\end{aligned}
$$

Proof. The action of the map $A$ can be written in terms of its Schwartz kernel as
12.21

$$
\begin{equation*}
A u(x)=\int K(x, y) u(y) d y=\left(\pi_{1}\right)_{*}(K \cdot(1 \boxtimes u)) \tag{4.229}
\end{equation*}
$$

Here $1 \boxtimes u$ is always defined and

$$
\begin{equation*}
\mathrm{WF}(1 \boxtimes u) \subset\{(x, y, 0, \eta) ;(y, \eta) \in \mathrm{WF}(u)\} \tag{4.230}
\end{equation*}
$$

So the main question is, when is the product defined? Our sufficient condition for this is:

$$
\begin{equation*}
(x, y, \xi, \eta) \in \mathrm{WF}(K) \Longrightarrow(x, y,-\xi,-\eta) \notin \mathrm{WF}(1 \boxtimes u) \tag{4.231}
\end{equation*}
$$

which is

$$
\begin{gather*}
(x, y, 0, \eta) \in \mathrm{WF}(K) \Longrightarrow(x, y, 0,-\eta) \notin \mathrm{WF}(1 \boxtimes u)  \tag{4.232}\\
\text { i.e. }(y,-\eta) \notin \mathrm{WF}(u) \tag{4.233}
\end{gather*}
$$

This of course is $\left(\frac{12.18}{7.226}\right)$ :

$$
\begin{gather*}
A u \text { is defined (by continuity) if }  \tag{4.234}\\
\{(y, \eta) \in \mathrm{WF}(u) ; \exists(x, 0, y,-\eta) \in \mathrm{WF}(A)\}=\emptyset . \tag{4.235}
\end{gather*}
$$

Then from our bound on the wavefront set of a product

$$
\begin{gather*}
\mathrm{WF}(K \cdot(1 \boxtimes u)) \subset \\
\left\{(x, y, \xi, \eta) ;(\xi, \eta)=\left(\xi^{\prime}, \eta^{\prime}\right)+\left(0, \eta^{\prime \prime}\right)\right. \text { with } \\
\left.\left(x, y, \xi^{\prime}, \eta^{\prime}\right) \in \mathrm{WF}(K) \text { and }\left(x, \eta^{\prime \prime}\right) \in \mathrm{WF}(u)\right\}  \tag{4.236}\\
\cup\{(x, y, \xi, \eta) ;(x, y, \xi, \eta) \in \mathrm{WF}(K), y \in \operatorname{supp}(u)\} \\
\cup\{(x, y, 0, \eta) ;(x, y) \in \operatorname{supp}(A)(y, \eta) \in \mathrm{WF}(u)\} .
\end{gather*}
$$

This gives the bound
(4.237) $\mathrm{WF}\left(\pi_{*}(K \cdot(1 \boxtimes u))\right) \subset\{(x, \xi) ;(x, y, \xi, 0) \in \mathrm{WF}(K \cdot(1 \boxtimes u))$ for some $y\}$

$$
\begin{equation*}
\subset \mathrm{WF}^{\prime}(A) \circ \mathrm{WF}(u) \tag{4.238}
\end{equation*}
$$

### 4.24. Applications

Having proved this rather general theorem, let us note some examples and applications.

First, for pseudodifferential operators we know that

$$
\begin{equation*}
\mathrm{WF}^{\prime}(A) \subset\{(x, x, \xi, \xi)\} \tag{4.239}
\end{equation*}
$$

i.e. corresponds to the identity relation (which is a map). Then $\left(\frac{12.19}{1.227}\right)$ is the microlocality of pseudodifferential operators. The next result also applies to all pseudodifferential operators.

$$
\begin{equation*}
\mathrm{WF}^{\prime}(K) \cap\{(x, y, \xi, 0)\}=\emptyset \tag{4.240}
\end{equation*}
$$

then the operator with Scwartz kernel $K$ defines a continuous linear map

$$
\begin{equation*}
A: \mathcal{C}_{c}^{\infty}\left(\Omega_{2}\right) \longrightarrow \mathcal{C}_{c}^{\infty}\left(\Omega_{1}\right) \tag{4.241}
\end{equation*}
$$

If

$$
\begin{equation*}
\mathrm{WF}^{\prime}(K) \cap\{(x, y, 0, \eta)\}=\emptyset \tag{4.242}
\end{equation*}
$$

then $A$ extends by continuity to

$$
\begin{equation*}
A: \mathcal{C}_{c}^{-\infty}\left(\Omega_{2}\right) \longrightarrow \mathcal{C}_{c}^{-\infty}\left(\Omega_{1}\right) \tag{4.243}
\end{equation*}
$$

Proof. Immediate from $\left(\frac{12.17}{4.224}\right)-\left(\frac{12.26}{4.243}\right)$.

### 4.25. Problems

Problem 4.9. Show that the general definition ( $\frac{8.10}{4.62}$ ) reduces to

$$
\begin{equation*}
\operatorname{WF}(u)=\bigcap\left\{\Sigma_{0}(A) ; A \in \Psi_{\infty}^{0}\left(\mathbb{R}^{n}\right) \text { and } A u \in \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right)\right\}, u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \tag{4.244}
\end{equation*}
$$

and prove the basic result of 'microlocal elliptic regularity:'

$$
\begin{align*}
& \text { If } u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \text { and } A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right) \text { then } \\
& \operatorname{WF}(u) \subset \Sigma(A) \cup \operatorname{WF}(A u) . \tag{4.245}
\end{align*}
$$

Problem 4.10. Compute the wavefront set of the following distributions:

$$
\begin{gather*}
\delta(x) \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right),|x| \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) \text { and } \\
\chi_{\mathbb{B}^{n}}(x)= \begin{cases}1 & |x| \leq 1 \\
0 & |x|>1\end{cases} \tag{4.246}
\end{gather*}
$$

Problem 4.11. Let $\Gamma \subset \mathbb{R}^{n} \times\left(\mathbb{R}^{n} \backslash 0\right)$ be an open cone and define

$$
\begin{align*}
\mathcal{C}_{c, \Gamma}^{-\infty}\left(\mathbb{R}^{n}\right)= & \left\{u \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right) ; A u \in \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right)\right.  \tag{4.247}\\
& \left.\forall A \in \Psi_{\infty}^{0}\left(\mathbb{R}^{n}\right) \text { with } \mathrm{WF}^{\prime}(A) \cap \Gamma=\emptyset\right\} \tag{4.248}
\end{align*}
$$

Describe a complete topology on this space with respect to which $\mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n}\right)$ is a dense subspace.
12.30 Problem 4.12. Show that, for any pseudodifferential operator $A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$, $\mathrm{WF}^{\prime}(A)=\mathrm{WF}^{\prime}\left(A^{*}\right)$.

Problem 4.13. Give an alternative proof to Lemma $\frac{9.1}{4.5}$ along the following lines (rather than using Lemma $\frac{1.15}{2.75}$ ). If $\sigma_{L}(A)$ is the left reduced symbol then for $\epsilon>0$ small enough

$$
\begin{equation*}
b_{0}=\gamma_{\epsilon} / \sigma_{L}(A) \in S_{\infty}^{-m}\left(\mathbb{R}^{n} ; \mathbb{R}^{n}\right) \tag{4.249}
\end{equation*}
$$

If we choose $B_{0} \in \Psi_{\infty}^{-m}\left(\mathbb{R}^{n}\right)$ with $\sigma_{L}\left(B_{0}\right)=b_{0}$ then
9.5

$$
\begin{equation*}
\mathrm{Id}-A \circ B_{0}=G \in \Psi_{\infty}^{0}\left(\mathbb{R}^{n}\right) \tag{4.250}
\end{equation*}
$$

has principal symbol

$$
\begin{equation*}
\sigma_{0}(G)=1-\sigma_{L}(A) \cdot b_{0} \tag{4.251}
\end{equation*}
$$

From ( $\frac{9.4}{4.7} 7$ )

$$
\begin{equation*}
\gamma_{\epsilon / 4} \sigma_{0}(G)=\gamma_{\epsilon / 4} \tag{4.252}
\end{equation*}
$$

Thus we conclude that if $\sigma_{L}(C)=\gamma_{\epsilon / 4}$ then

$$
\begin{equation*}
G=(\operatorname{Id}-C) G+C G \text { with } C G \in \Psi_{\infty}^{-1}\left(\mathbb{R}^{n}\right) \tag{4.253}
\end{equation*}
$$

Thus ( $\frac{19.5}{1.250)}$ becomes
$9.6 \quad$ Id $-A B_{0}=C G+R_{1} \quad \mathrm{WF}^{\prime}\left(R_{1}\right) \nexists z$.
Let $B_{1} \sim \sum_{j \geq 1}(C G)^{j}, B_{1} \in \Psi^{-1}$ and set

$$
\begin{equation*}
B=B_{0}\left(\operatorname{Id}+B_{1}\right) \in \Psi_{\infty}^{-m}\left(\mathbb{R}^{n}\right) \tag{4.255}
\end{equation*}
$$

From ( $\frac{9.6}{4.254)}$

$$
\begin{align*}
A B & =A B_{0}\left(I+B_{1}\right)  \tag{4.256}\\
& =(\operatorname{Id}-C G)\left(I+B_{1}\right)-R_{1}\left(\operatorname{Id}+B_{1}\right)  \tag{4.257}\\
& =\operatorname{Id}+R_{2}, \quad \mathrm{WF}^{\prime}\left(R_{2}\right) \not \supset z . \tag{4.258}
\end{align*}
$$

Thus $B$ is a right microlocal parametrix as desired. Write out the construction of a left parametrix using the same method, or by finding aright parametrix for the adjoint of $A$ and then taking adjoints using Problem 4.12.
12.32 Problem 4.14. Essential uniqueness of left and right parametrices.
12.33 PROBLEM 4.15. If $(\bar{x}, \bar{\xi}) \in \mathbb{R}^{n} \times\left(\mathbb{R}^{n} \backslash 0\right)$ is a given point, construct a distribution $u \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right)$ which has

$$
\begin{equation*}
\mathrm{WF}(u)=\{(\bar{x}, t \bar{\xi}) ; t>0\} \subset \mathbb{R}^{n} \times\left(\mathbb{R}^{n} \backslash 0\right) \tag{4.259}
\end{equation*}
$$

Problem 4.16. Suppose that $A \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ has Schwartz kernel of compact support. If $u \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right)$ use the four 'elementary operations' (and an earlier result on the wavefront set of kernels) to investigate under what conditions

$$
\begin{equation*}
\kappa(x, y)=K_{A}(x, y) u(y) \text { and then } \gamma(x)=\left(\pi_{1}\right)_{*} \kappa \tag{4.260}
\end{equation*}
$$

make sense. What can you say about $\mathrm{WF}(\gamma)$ ?
Problem 4.17. Consider the projection operation under $\pi_{1}: \mathbb{R}^{p} \times \mathbb{R}^{k} \longrightarrow \mathbb{R}^{p}$. Show that $\left(\pi_{1}\right)_{*}$ can be extended to some distributions which do not have compact support, for example

$$
\begin{equation*}
\left\{u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right) ; \operatorname{supp}(u) \cap K \times \mathbb{R}^{k} \text { is compact for each } K \subset \subset \mathbb{R}^{n}\right\} \tag{4.261}
\end{equation*}
$$

## CHAPTER 5

## Pseudodifferential operators on manifolds

In this chapter the notion of a pseudodifferential on a manifold is discussed. Some preliminary material on manifolds is therefore necessary. However the discussion of the basic properties of differentiable manifolds is kept to a bare minimum. For a more leisurely treatement the reader might well consult XX or YY. Our main aims here are first, to be able to prove the Hodge theorem (given the deRham theorem). Then we describe some global object which are very useful in applications, namely a global quantization map, the structure of complex powers and the zeta function.

## 5.1. $\mathcal{C}^{\infty}$ structures

Let $X$ be a paracompact Hausdorff topological space. A $\mathcal{C}^{\infty}$ structure on $X$ is a subspace

$$
\begin{equation*}
\mathcal{F} \subset C^{0}(X)=\{u: X \longrightarrow \mathbb{R} \text { continuous }\} \tag{5.1}
\end{equation*}
$$

with the following property:
For each $\bar{x} \in X$ there exists elements $f_{1}, \ldots, f_{n} \in \mathcal{F}$ such that for some open neighbourhood $\Omega \ni \bar{x}$

$$
\begin{equation*}
F: \Omega \ni x \longmapsto\left(f_{1}(x), \ldots, f_{n}(x)\right) \in \mathbb{R}^{n} \tag{5.2}
\end{equation*}
$$

is a homeomorphism onto an open subset of $\mathbb{R}^{n}$ and every $f \in \mathcal{F}$ satisfies

$$
\begin{equation*}
f \upharpoonright \Omega=g \circ F \quad \text { for some } g \in \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right) \tag{5.3}
\end{equation*}
$$

The map $\left(\frac{13.2}{5.2}\right)$ is a coordinate system near $\bar{x}$. Two $\mathcal{C}^{\infty}$ structures $\mathcal{F}_{1}$ and $\mathcal{F}_{2}$ are 'compatible' if $\mathcal{F}_{1} \cup \mathcal{F}_{2}$ is also a $\mathcal{C}^{\infty}$ structure. Compatibility in this sense is an equivalence relation on $\mathcal{C}^{\infty}$ structures. It therefore makes sense to say that:

DEFINITION 5.1. A $\mathcal{C}^{\infty}$ manifold is a (connected) paracompact Hausdorff topological space with a maximal $\mathcal{C}^{\infty}$ structure.

The maximal $\mathcal{C}^{\infty}$ structure is conventionally denoted

$$
\begin{equation*}
\mathcal{C}^{\infty}(X) \subset \mathcal{C}^{0}(X) \tag{5.4}
\end{equation*}
$$

It is necessarily an algebra. If we let $\mathcal{C}_{c}^{\infty}(W) \subset \mathcal{C}^{\infty}(X)$ denote the subspace of functions which vanish outside a compact subset of $W$ then any local coordinates ( 5.2 ) have the property

$$
\begin{equation*}
F^{*}: \mathcal{C}_{c}^{\infty}(F(\Omega)) \longleftrightarrow\left\{u \in \mathcal{C}^{\infty}(X) ; u=0 \text { on } X \backslash K, K \subset \subset \Omega\right\} \tag{5.5}
\end{equation*}
$$

Futhermore $\mathcal{C}^{\infty}(X)$ is local:

$$
u: X \longrightarrow \mathbb{R} \text { and } \forall \bar{x} \in X \exists \Omega_{\bar{x}} \text { open, } \Omega_{\bar{x}} \ni \bar{x}
$$

$$
\begin{equation*}
\text { s.t. } u-f_{\bar{x}}=0 \text { on } \Omega_{\bar{x}} \text { for some } f_{\bar{x}} \in \mathcal{C}^{\infty}(X) \Longrightarrow u \in \mathcal{C}^{\infty}(X) \tag{5.6}
\end{equation*}
$$

A map $G: X \longrightarrow Y$ between $\mathcal{C}^{\infty}$ manifolds $X$ and $Y$ is $\mathcal{C}^{\infty}$ if

$$
\begin{equation*}
G^{*}: \mathcal{C}^{\infty}(Y) \longrightarrow \mathcal{C}^{\infty}(X) \tag{5.7}
\end{equation*}
$$

i.e. $G \circ u \in \mathcal{C}^{\infty}(X)$ for all $u \in \mathcal{C}^{\infty}(Y)$.

### 5.2. Form bundles

A vector bundle is a triple $\pi: V \longrightarrow X$ consisting of two manifolds, $X$ and $V$, and a surjective $\mathcal{C}^{\infty}$ map $\pi$ with each

$$
\begin{equation*}
V_{x}=\pi^{-1}(x) \tag{5.8}
\end{equation*}
$$

having a linear structure such that

$$
\begin{equation*}
\mathcal{F}=\left\{u: V \longrightarrow \mathbb{R}, u \text { is linear on each } V_{x}\right\} \tag{5.9}
\end{equation*}
$$

is a $\mathcal{C}^{\infty}$ structure on $V$ compatible with $\mathcal{C}^{\infty}(V)$ (i.e. contained in it, since it is maximal).

The basic example is the cotangent bundle which we defined before for open sets in $\mathbb{R}^{n}$. The same definition works here. Namely for each $\bar{x} \in X$ set

$$
\begin{gather*}
\mathcal{I}_{\bar{x}}=\left\{u \in \mathcal{C}^{\infty}(X) ; u(\bar{x})=0\right\} \\
\mathcal{I}_{\bar{x}}^{2}=\left\{u=\sum_{\text {finite }} u_{i} u_{i}^{\prime} ; u_{i}, u_{i}^{\prime} \in \mathcal{I}_{\bar{x}}\right\}  \tag{5.10}\\
T_{\bar{x}}^{*} X=\mathcal{I}_{\bar{x}} / \mathcal{I}_{\bar{x}}^{2}, T^{*} X=\bigcup_{\bar{x} \in X} T_{\bar{x}}^{*} X
\end{gather*}
$$

So $\pi: T^{*} X \longrightarrow X$ just maps each $T_{\bar{x}}^{*} X$ to $\bar{x}$. We need to give $T^{*} X$ a $\mathcal{C}^{\infty}$ structure so that "it" (meaning $\pi: T^{*} X \longrightarrow X$ ) becomes a vector bundle. To do so note that the differential of any $f \in \mathcal{C}^{\infty}(X)$

$$
\begin{equation*}
d f: X \longrightarrow T^{*} X \quad d f(\bar{x})=[f-f(\bar{x})] \in T_{\bar{x}}^{*} X \tag{5.11}
\end{equation*}
$$

is a section $(\pi \circ d f=\mathrm{Id})$. Put

$$
\begin{equation*}
\mathcal{F}=\left\{u: T^{*} X \longrightarrow \mathbb{R} ; u \circ d f: X \longrightarrow \mathbb{R} \text { is } \mathcal{C}^{\infty} \forall f \in \mathcal{C}^{\infty}(X)\right\} \tag{5.12}
\end{equation*}
$$

Then $\mathcal{F}=\mathcal{C}^{\infty}\left(T^{*} X\right)$ is a maximal $\mathcal{C}^{\infty}$ structure on $T^{*} X$ and

$$
\mathcal{F}_{\mathrm{lin}}=\left\{u: T^{*} X \longrightarrow \mathbb{R}, \text { linear on each } T_{\bar{x}}^{*} X ; u \in \mathcal{F}\right\}
$$

is therefore compatible with it. Clearly $d f$ is $\mathcal{C}^{\infty}$.
Any (functorial) operation on finite dimensional vector spaces can be easily seen to generate new vectors bundles from old. Thus duality, tensor product, exterior powers all lead to new vector bundles:

$$
\begin{equation*}
T_{x} X=\left(T_{x}^{*} X\right)^{*}, T X=\bigcup_{x \in X} T_{x} X \tag{5.13}
\end{equation*}
$$

is the tangent bundle

$$
\Lambda_{x}^{k} X=\left\{u: T_{x} X \stackrel{k \text { factors }}{\left.\times \cdots \times T_{x} X \longrightarrow \mathbb{R} ; u \text { is multilinear and antisymmetric }\right\}}\right.
$$

leads to the $k$-form bundle

$$
\Lambda^{k} X=\bigcup_{x \in X} \Lambda_{x}^{k} X, \Lambda^{1} X \simeq T^{*} X
$$

where equivalence means there exists (in this case a natural) $\mathcal{C}^{\infty}$ diffeomorphism mapping fibres to fibres linearly (and in this case projecting to the identity on $X$ ).

A similar construction leads to the density bundles
$\Omega_{x}^{\alpha} X=\left\{u: T_{x}^{n=\operatorname{dim} X} X \wedge \wedge \wedge_{x}^{\text {factors }} X \longrightarrow \mathbb{R} ;\right.$ absolutely homogeneous of degree $\left.\alpha\right\}$
that is

$$
u\left(t v_{1} \wedge \ldots v_{n}\right)=|t|^{\alpha} u\left(v_{1} \wedge \cdots \wedge v_{n}\right)
$$

These are important because of integration. In general if $\pi: V \longrightarrow X$ is a vector bundle then

$$
\mathcal{C}^{\infty}(X ; V)=\{u: X \longrightarrow V ; \pi \circ u=\mathrm{Id}\}
$$

is the space of sections. It has a natural linear structure. Suppose $W \subset X$ is a coordinate neighbourhood and $u \in \mathcal{C}^{\infty}(X ; \Omega), \Omega=\Omega^{1} X$, has compact support in $W$. Then the coordinate map gives an identification

$$
\Omega_{x}^{*} X \longleftrightarrow \Omega_{F(x)}^{*} \mathbb{R}^{n} \quad \forall \alpha
$$

and

$$
\begin{equation*}
\int u=\int_{\mathbb{R}^{n}} g_{u}(x), \quad u=g_{u}(x)|d x| \tag{5.14}
\end{equation*}
$$

is defined independent of coordiantes. That is the integral

$$
\begin{equation*}
\int: \mathcal{C}_{c}^{\infty}(X ; \Omega) \longrightarrow \mathbb{R} \tag{5.15}
\end{equation*}
$$

is well-defined.

### 5.3. Pseudodifferential operators

Let $X$ be a $\mathcal{C}^{\infty}$ manifold, and let $\mathcal{C}_{c}^{\infty}(X) \subset \mathcal{C}^{\infty}(X)$ be the space of $\mathcal{C}^{\infty}$ functions of compact support. Then, for any $m \in \mathbb{R}, \Psi^{m}(X)$ is the space of linear operators

$$
\begin{equation*}
A: \mathcal{C}_{c}^{\infty}(X) \longrightarrow \mathcal{C}^{\infty}(X) \tag{5.16}
\end{equation*}
$$

with the following properties. First,

$$
\text { if } \phi, \psi \in \mathcal{C}^{\infty}(X) \text { have disjoint supports then } \exists K \in \mathcal{C}^{\infty}\left(X \times X ; \Omega_{R}\right)
$$

$$
\begin{equation*}
\text { such that } \forall u \in \mathcal{C}_{c}^{\infty}(X) \phi A \psi u=\int_{X} K(x, y) u(y) \tag{5.17}
\end{equation*}
$$

and secondly if $F: W \longrightarrow \mathbb{R}^{n}$ is a coordinate system in $X$ and $\psi \in \mathcal{C}_{c}^{\infty}(X)$ has support in $W$ then

$$
\begin{gathered}
\exists B \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right), \operatorname{supp}(B) \subset W \times W \text { s.t. } \\
\psi A \psi u \upharpoonright W=F^{*}\left(B\left(\left(F^{-1}\right)^{*}(\psi u)\right)\right) \forall u \in \mathcal{C}_{c}^{\infty}(X)
\end{gathered}
$$

This is a pretty horrible definition, since it requires us to check every coordinate system, at least in principle. In practice the coordinate-invariance we proved earlier means that this is not necessary and also that there are plenty of examples!

Any open cover of a $\mathcal{C}^{\infty}$ manifold has a partition of unity subordinate to it, i.e. if $A_{r} \subset X$ are open sets for $r \in R$ and

$$
\begin{equation*}
X=\bigcup_{r \in R} A_{r} \tag{5.18}
\end{equation*}
$$

there exists $\phi_{i} \in \mathcal{C}_{c}^{\infty}(X)$, all non-negative with locally finite support:

$$
\begin{equation*}
\forall i \operatorname{supp}\left(\phi_{i}\right) \cap \operatorname{supp}\left(\phi_{j}\right) \neq \emptyset \text { for a finite set of indices } j \tag{5.19}
\end{equation*}
$$

where each $\operatorname{supp}\left(\phi_{i}\right) \subset A_{r}$ for some $r=r(i)$ and

$$
\begin{equation*}
\sum_{i} \phi_{i}(x)=1 \quad \forall x . \tag{5.20}
\end{equation*}
$$

In fact one can do slightly better than this, for a covering by coordinate neighbourhoods.
13.9 Lemma 5.1. There exists a partition of unity on $X$ ( $\mathcal{C}^{\infty}$ manifold) $\phi_{i}$ s.t. for every $i$

$$
\begin{equation*}
\bigcup_{j}\left\{\operatorname{supp}\left(\phi_{j}\right) ; \operatorname{supp}\left(\phi_{j}\right) \cap \operatorname{supp}\left(\phi_{i}\right) \neq 0\right\} \tag{5.21}
\end{equation*}
$$

is contained in a coordinate neighbourhood!
Using such a partition of unity we see that every element of $\Psi^{m}(X)$ can be written in the form

$$
A=\sum_{i} \sum_{j} \phi_{j} F^{*} K_{A} \phi_{i}
$$

where the terms have smooth kernels if the supports of $\phi_{i}$ and $\phi_{j}$ do not meet, or else are pseudodifferential operators in any local coordinates in a patch containing both supports. Below we shall use this to prove:
13.10 Theorem 5.1. Let $X$ be a compact $\mathcal{C}^{\infty}$ manifold then the pseudodifferential operators $\Psi^{*}(X)$ form a symbol-filtered ring.

If $X$ is a $\mathcal{C}^{\infty}$ manifold we have defined the space $\Psi^{m}(X)$ as consisting of those linear operators

$$
\begin{equation*}
A: \mathcal{C}_{c}^{\infty}(X) \longrightarrow \mathcal{C}^{\infty}(X) \tag{5.22}
\end{equation*}
$$

which are given locally by pseudodifferential operators of order $m$ on $X$, in a precise sense. Let us recast this definition in terms of the Schwartz kernel theorem. Over the product, $X^{2}$, consider the right density bundle, $\Omega_{R}=\pi_{R}^{*} \Omega$. Here we use the pull-back operation on vector bundles:

THEOREM 5.2. If $W \longrightarrow Y$ is a $\mathcal{C}^{\infty}$ vector bundle and $F: X \longrightarrow Y$ is a $\mathcal{C}^{\infty}$ map then $F^{*} W \longrightarrow X$ is a well-defined $\mathcal{C}^{\infty}$ vector bundle over $X$ with total space

$$
\begin{equation*}
F^{*} W=\bigcup_{x \in X} W_{F(x)} \tag{5.23}
\end{equation*}
$$

if $\phi \in \mathcal{C}^{\infty}(Y ; W)$ then $F^{*} \phi$ is a section of $F^{*} W$ and $\mathcal{C}^{\infty}\left(X ; F^{*} W\right)$ is spanned by $\mathcal{C}^{\infty}(X) \cdot F^{*} \mathcal{C}^{\infty}(Y ; W)$.

Distributional sections of any $\mathcal{C}^{\infty}$ vector bundle can be defined in two equivalent ways:

$$
\begin{equation*}
\text { "Algebraically" } \mathcal{C}^{-\infty}(X ; W)=\mathcal{C}^{-\infty}(X) \bigotimes_{\mathcal{C}^{\infty}(X)} \mathcal{C}^{\infty}(X ; W) \tag{5.24}
\end{equation*}
$$

or

$$
\begin{equation*}
\text { "Analytically" } \mathcal{C}^{-\infty}(X ; W)=\left[\mathcal{C}_{c}^{\infty}\left(X ; \Omega \otimes W^{\prime}\right)\right] \tag{5.25}
\end{equation*}
$$

where $W^{\prime}$ is the dual bundle and $\Omega$ the density bundle over $X$. In order to use $\left(\frac{14.5}{5.25}\right)$ we need to define a topology on $\mathcal{C}_{c}^{\infty}(X ; U)$ for any vector bundle $U$ over $X$. One can do this by reference to local coordinates.

If $W \longrightarrow X$ is a vector bundle the spaces $S^{m}(W)$ of symbols on $W$ is welldefined for each $m \in \mathbb{R}$.
14.6 Proposition 5.1. A pseudodifferential operator $A \in \Psi^{m}(X)$ can be written in terms of its kernel

$$
\begin{equation*}
K_{A} \in \mathcal{C}^{\infty}\left(X^{2} ; \Omega_{R}\right) \tag{5.26}
\end{equation*}
$$

where
14.7

$$
\begin{equation*}
K_{A} \text { is } \mathcal{C}^{\infty} \text { in } X^{2} \backslash \Delta \tag{5.27}
\end{equation*}
$$

and if $X_{i} \subset X$ is a coordinate partition, $\rho_{i} \in \mathcal{C}_{c}^{\infty}(X)$ has support in $X_{i}$ then in terms of the same coordinates $x_{i}=F_{i}^{*}\left(x_{i}\right)$ and $y_{i}=F^{*}\left(x_{i}\right)$ in the two factors of $X$

$$
\begin{equation*}
\rho_{i}(x) \rho(y) K_{A}=K_{i}(x, y)|d y|, K_{i} \in \Psi^{m}\left(\mathbb{R}^{n}\right) \tag{tabular}
\end{equation*}
$$

Suppose $\rho_{i}^{2}$ is a partition of ${ }_{14.8}$ nity of $X$, subordinate to a coordinate covering. For each $i$ the symbol of $K_{i}$ in ( 5.28$)$ is an equivalence class on $\mathbb{R}^{n} \times\left(\mathbb{R}^{n} \backslash 0\right)$, with support in $\operatorname{supp}\left(\rho_{i}\right) *\left(\mathbb{R}^{2} \backslash 0\right)$. Set

$$
\begin{equation*}
\sigma_{m}(\tau)=\sum_{\left\{i, \pi(\tau) \in \operatorname{supp}\left(\rho_{i}\right)\right\}} a_{i}\left(x^{(i)}, \xi^{(i)}\right) \tag{5.29}
\end{equation*}
$$

where

$$
\begin{equation*}
\tau=F_{i}^{*}\left(\sum_{j} \xi_{j}^{(i)} \cdot d x_{j}\right) \quad \xi^{(i)} \cdot d x \in T_{x^{(i)}}^{*} \mathbb{R}^{n}, x^{(i)}=F_{i}(\pi(\tau)) \tag{5.30}
\end{equation*}
$$

and the $a_{i}$ are representations of the symbols of the $K_{i}$. This defines a function on $T^{*} X \backslash 0$, in fact the equivalence class

$$
\begin{equation*}
\sigma_{m}(A) \in S^{m-[1]}\left(T^{*} X\right) \tag{5.31}
\end{equation*}
$$

is well-defined.
Proposition 5.2. The principal symbol map in (14.10), defined as in $\left(\frac{14.9}{5.29}\right)$, gives a short exact sequence:

$$
\begin{equation*}
0 \hookrightarrow \Psi^{m-1}(X) \hookrightarrow \Psi^{m}(X) \xrightarrow{\sigma_{m}} S^{m-[1]}\left(T^{*} X\right) \longrightarrow 0 \tag{5.32}
\end{equation*}
$$

Proof. First we need to check that $\sigma_{m}(A)$ is indeed well-defined. This involves checking what happens under a change of coordinate covering and a change of partition of unity subordinate to it. First, under a change of partition of unity, subordinate to a fixed covering note that

$$
\begin{align*}
& \rho_{j}^{\prime}(x) \rho_{j}^{\prime}(y) K_{A}=\sum_{i} \rho_{j}^{\prime}(x) \rho_{i}^{2}(x) \rho_{j}^{\prime}(y) K_{A} \\
& \equiv \sum_{i} \rho_{j}^{\prime}(x) \rho_{i}(x) \rho_{i}(y) \rho_{j}^{\prime}(y) K_{A}  \tag{5.33}\\
& \equiv\left(\rho_{j}^{\prime}\right)^{2} \sum_{i} \rho_{i}(x) K_{A} \rho_{j}(y)
\end{align*}
$$

where equality is madule $\Psi^{m-1}$, since $\left[\phi, \Psi^{m}\right] \subset \Psi^{m-1}$ for any $\mathcal{C}_{14.9}^{\infty}$ function $\phi$.
It follows from ( $\frac{14.32}{5.32}$ ) that the principal symbols, defined by $\left(\frac{14.9}{5.29}\right)$, for the two partitions of unity agree.

For a change of coordinate covering it suffices to use the transformation law for the principal symbol under a diffeomorphism of $\mathbb{R}^{n}$ and the freedom, just established, to choose the partition of unity to be subordinates to both coordinate coverings. Thus $\sigma_{m}$ is well-defined.

Certainly if $A \in \Psi^{m-1}(X)$ then $\sigma_{m}(N) \equiv 0$. Moreover if $A \in \Psi^{m}(X)$ and $\sigma_{m}(A) \equiv 0$ then all the operators $\rho_{i}(x) K \rho_{i}(y)$ are actually of order $m-1$. Since

$$
\begin{equation*}
A \equiv \sum_{i} \sum_{i} \rho_{i}^{\prime} \rho_{j}(X) A \rho_{i}^{\prime} \rho_{j} \quad \bmod \Psi^{m-1}(X) \tag{5.34}
\end{equation*}
$$

for any two partitions of units $\rho_{i}^{\prime 2},\left(\rho_{i}^{\prime}\right)^{2}$ we can choose the $\rho_{i}^{\prime}$ to each have support in a region where $\rho_{j} \neq 0$ for some $j$. Then (5.34) shows that $A \in \Psi^{m-1}(X)$.

Thus it only remains to show that the map $\sigma_{m}$ is surjective. If $a \in S^{m}\left(T^{*} X\right)$ choose $A_{i} \in \Psi_{\infty}^{m}\left(\mathbb{R}^{n}\right)$ by

$$
\begin{equation*}
\sigma_{L}\left(A_{i}\right)=\rho_{i}(x)\left(F^{*}\right)^{-1} a_{i} \rho_{i}(y) \in S_{\infty}^{m}\left(\mathbb{R}^{n} \times \mathbb{R}^{n}\right) \tag{5.35}
\end{equation*}
$$

and set

$$
\begin{equation*}
A=\sum_{i} F_{i}^{*} A_{i} G_{i}^{*} \quad G_{i}=F_{i}^{-1} \tag{5.36}
\end{equation*}
$$

Then, from $\left(\frac{14.9}{5.29}\right) \sigma_{m}(A) \equiv a$ by invariance of the principal symbol.
The other basic properties of the calculus are easily established. For example

$$
\begin{equation*}
\sigma_{m+m^{\prime}}(A \cdot B)=\sigma_{m}(A) \cdot \sigma_{m}(B) \tag{5.37}
\end{equation*}
$$

if $A \in \Psi^{m}(X), B \in \Psi^{m^{\prime}}(X), X$ compact. Similarly note that

$$
\begin{equation*}
A B=\sum_{i, j} \rho_{i}^{2} A \rho_{j}^{2} B=\sum_{i, j} \rho_{i} A \rho_{i} \cdot \rho_{j} B \rho_{i} \quad \bmod \Psi^{m_{m}^{\prime}-1} \tag{5.38}
\end{equation*}
$$

which gives $\left(\frac{14.16}{\frac{3}{4}}\right.$
In § 4.9 we used the symbol calculus to construct a left and right parametrix for an elliptic element of $\Psi^{m}(X)$, where $X$ is compact, i.e. an element $B \in \Psi^{-m}(X)$, such that

$$
\begin{equation*}
A B-\mathrm{Id}, B A-\mathrm{Id} \in \Psi^{-\infty}(X) \tag{5.39}
\end{equation*}
$$

As a consequence of this construction note that:
14.18 Proposition 5.3. If $A \in \Psi^{m}(X)$ is elliptic, and $X$ is compact, then

$$
\begin{equation*}
A: \mathcal{C}^{\infty}(X) \longrightarrow \mathcal{C}^{\infty}(X) \tag{5.40}
\end{equation*}
$$

is Fredholm, i.e. has finite dimensional null space and closed range with finite dimensional complement. If $\nu$ is a non-vanishing $\mathcal{C}^{\infty}$ measure on $X$ and a generalized inverse of $A$ is defined by

$$
\begin{align*}
& G u=f \text { if } u \in \operatorname{Ran}(A), A f=u, f \perp_{\nu} \operatorname{Nul}(A) \\
& G u=0 \text { if } u \perp_{\nu} \operatorname{Ran}(A) \tag{5.41}
\end{align*}
$$

then $G \in \Psi^{-m}(X)$ satisfies
14.21

$$
\begin{align*}
G A & =\mathrm{Id}-\pi_{N} \\
A G & =\mathrm{Id}-\pi_{R} \tag{5.42}
\end{align*}
$$

where $\pi_{N}$ and $\pi_{R}$ are $\nu$-orthogonal projections onto the null space of $A$ and the $\nu$-orthocomplement of the range of $A$ respectively.

Proof. The main point to note is that $E \in \Psi^{-\infty}(X)$ is smoothing,

$$
\begin{equation*}
E: \mathcal{C}^{-\infty}(X) \longrightarrow \mathcal{C}^{\infty}(X) \quad \forall E \in \Psi^{-\infty}(X) \tag{5.43}
\end{equation*}
$$

Such a map is compact on $L^{2}(X)$, i.e. maps bounded sets into precompact sets by the theorem of Ascoli and Arzela. The second thing to recall is that a Hilbert space with a compact unit ball is finite dimensional. Then

$$
\begin{equation*}
\operatorname{Nul}(A)=\left\{u \in \mathcal{C}^{\infty}(X) ; A u=0\right\}=\left\{u \in L^{2}(X) ; A u=0\right\} \tag{5.44}
\end{equation*}
$$

since, from $\left(\frac{14.23}{5.44}\right) A u=0 \Longrightarrow(B A-\mathrm{Id}) u=-E u, E \in \Psi^{\infty}(X)$, so $A u=0$, $u \in \mathcal{C}^{-\infty}(X) \Longrightarrow u \in \mathcal{C}^{\infty}(X)$. Then

$$
\begin{equation*}
\operatorname{Nul}(A)=\left\{u \in L^{2}(X) ; A u=0 \int|u|^{2} \nu=1\right\} \subset L^{2}(X) \tag{5.45}
\end{equation*}
$$

is compact since it is closed ( $A$ is continuous) and so $\operatorname{Nul}(A)=E(\operatorname{Nul}(A))$ is precompact. Thus $\operatorname{Nul}(A)$ is finite dimensional.

Next let us show that $\operatorname{Ran}(A)$ is closed. Suppose $f_{j}=A u_{j} \longrightarrow f$ in $\mathcal{C}^{\infty}(X)$, $u_{j} \in \mathcal{C}^{\infty}(X)$. By what we have just shown we can assume that $u_{j} \perp_{\nu} \operatorname{Nul}(A)$. Now if $B$ is the parametrix

$$
\begin{equation*}
u_{j}=B f_{j}+E u_{j}, E \in \Psi^{-\infty}(X) \tag{5.46}
\end{equation*}
$$

Suppose, along some subsequence, $\left\|u_{j}\right\|_{\nu} \longrightarrow \infty$. Then

$$
\begin{equation*}
\frac{u_{j}}{\left\|u_{j}\right\|_{\nu}}=B\left(\frac{f_{j}}{\left\|u_{j}\right\|_{\nu}}\right)+E\left(\frac{u_{j}}{\left\|u_{j}\right\|_{\nu}}\right) \tag{5.47}
\end{equation*}
$$

shows that $\frac{u_{j}}{\left\|u_{j}\right\|_{\nu}}$ lies in a precompact subset of $L^{2}, \frac{u_{j}}{\left\|u_{j}\right\|_{\nu}} \longrightarrow u$. This is a contradiction, since $A u=0$ but $\|u\|=1$ and $u \perp_{\nu} \operatorname{Nul}(A)$. Thus the norm sequence $\left\|u_{j}\right\|$ is bounded and therefore the sequence has a weakly convergent subsequence, which we can relabel as $u_{j}$. The parametrix shows that $u=B f_{j}+E u_{j}$ is strongly convergent with limit $u$, which satisfies $A u=f$.

Finally we have to show that $\operatorname{Ran}(A)$ has a finite dimensional complement, If $\pi_{R}$ is orthogonal projection off $\operatorname{Ran}(A)$ then from the second part of ( 5.39$)$ $f=\pi_{R} E^{\prime} f$ for some smoothing operator $E$. This shows that the orthocomplement has compact unit ball, hence is finite dimensional.

Notice that it follows that the two projections in ( $\frac{14.21}{5.42)}$ are both smoothing operators of finite rank.

### 5.4. Pseudodifferential operators on vector bundles

We have just shown that any elliptic pseudodifferential operator, $A \in \Psi^{m}(X)$ on a compact manifold $X$ has a generalized inverse $B \in \Psi^{-m}(X)$, meaning

$$
\begin{align*}
& B A=\mathrm{Id}-\pi_{N} \\
& A B=\mathrm{Id}-\pi_{R} \tag{5.48}
\end{align*}
$$

where $\pi_{N}$ and $\pi_{R}$ are the orthogonal projections onto the null space of $A$ and the orthocomplement of the range of $A$ with respect to a prescribed $\mathcal{C}^{\infty}$ positive density $\nu$, both are elements of $\Psi^{-\infty}(X)$ and have finite rank. To use this theorem in geometric situations we need first to make the "trivial" extension to operators on sections of vector bundles.

As usual there are two ways (at least) to approach this extension; the high road and the low road. The "low" road is to go back to the definition of $\Psi^{m}(X)$ and
to generalize to $\Psi^{m}(X ; V, W)$. This just requires to take the definition, following (5.16), but using a covering with respect to which the bundles $V, W$ are both locally trivial. The local coordinate representatives of the pseudodifferential operator are then matrices of pseudodifferential opertors. The symbol mapping becomes

$$
\begin{equation*}
\Psi^{m}(X ; V, W) \longrightarrow S^{m-[1]}\left(T^{*} X ; \operatorname{Hom}(V, W)\right) \tag{5.49}
\end{equation*}
$$

where $\operatorname{Hom}(V, W) \simeq V \otimes W^{\prime}$ is the bundle of homomorphisms from $V$ to $W$ and the symbol space consists of symbolic sections of the lift of this bundle to $T^{*} X$. We leave the detailed description and proof of these results to the enthusiasts.

So what is the "high" road. This involves only a little sheaf-theoretic thought. Namely we want to define the space $\Psi^{m}(X ; V, W)$ using $\Psi^{m}(X)$ by:

$$
\begin{equation*}
\Psi^{m}(X ; V, W)=\Psi^{m}(X) \bigotimes_{\mathcal{C}^{\infty}\left(X^{2}\right)} \mathcal{C}^{\infty}\left(X^{2} ; V \boxtimes W^{\prime}\right) \tag{5.50}
\end{equation*}
$$

To make sense of this we first note that $\Psi^{m}(X)$ is a $\mathcal{C}^{\infty}\left(X^{2}\right)$-module as is the space $\mathcal{C}^{\infty}\left(X^{2} ; V \boxtimes W^{\prime}\right)$ where $V \boxtimes W^{\prime}$ is the "exterior" product:

$$
\begin{equation*}
\left(V \boxtimes W^{\prime}\right)_{(x, y)}=V_{x} \otimes W_{y}^{\prime} \tag{5.51}
\end{equation*}
$$

The tensor product in $\left(\frac{15.3}{5.50}\right)$ means that

$$
\begin{equation*}
A \in \Psi^{m}(X ; V, W) \text { is of the form } A=\sum_{i} A_{i} \cdot G_{i} \tag{5.52}
\end{equation*}
$$

where $A_{i} \in \Psi^{m}(X), G_{i} \in \mathcal{C}^{\infty}\left(X^{2} ; V \boxtimes W^{\prime}\right)$ and equality is fixed by the relation

$$
\begin{equation*}
\phi A \cdot G-A \cdot \phi G \equiv 0 \tag{5.53}
\end{equation*}
$$

Now what we really need to note is:
$15.7 \quad$ Proposition 5.4. For any compact $\mathcal{C}^{\infty}$ manifold $Y$ and any vector bundle $U$ over $Y$

$$
\begin{equation*}
\mathcal{C}^{-\infty}(Y ; U) \equiv \mathcal{C}^{-\infty}(Y) \bigotimes_{\mathcal{C}^{\infty}(Y)} \mathcal{C}^{\infty}(Y ; U) \tag{5.54}
\end{equation*}
$$

Proof. $\mathcal{C}^{-\infty}(Y ; U)=\left(\mathcal{C}^{\infty}\left(Y ; \Omega \otimes U^{\prime}\right)\right)^{\prime}$ is the definition. Clearly we have a mapping

$$
\begin{equation*}
\mathcal{C}^{-\infty}(Y) \bigotimes_{\mathcal{C}^{\infty}(Y)} \mathcal{C}^{\infty}(Y ; U) \ni \sum_{i} A_{i} \cdot g_{i} \longrightarrow \mathcal{C}^{-\infty}(Y ; U) \tag{5.55}
\end{equation*}
$$

given by
15.9

$$
\begin{equation*}
\sum_{i} u_{i} \cdot g_{i}(\psi)=\sum_{i} u_{i}\left(g_{i} \cdot \psi\right) \tag{5.56}
\end{equation*}
$$

since $g_{i} \psi \in \mathcal{C}^{\infty}(Y ; \Omega)$ and linearity shows that the map descends to the tensor product. To prove that the map is an isomorphism we construct an inverse. Since $Y$ is compact we can find a finite number of sections $g_{i} \in \mathcal{C}^{\infty}(Y ; U)$ such that any $u \in \mathcal{C}^{\infty}(Y ; U)$ can be written

$$
\begin{equation*}
u=\sum_{i} h_{i} g_{i} \quad h_{i} \in \mathcal{C}^{\infty}(Y) \tag{5.57}
\end{equation*}
$$

By reference to local coordinates the same is true of distributional sections with

$$
\begin{equation*}
h_{i}=u \cdot q_{i} \quad q_{i} \in \mathcal{C}^{\infty}\left(Y ; U^{\prime}\right) \tag{5.58}
\end{equation*}
$$

This gives a left and right inverse.
15.10 THEOREM 5.3. The calculus extends to operators sections of vector bundles over any compact $\mathcal{C}^{\infty}$ manifolds.

### 5.5. Laplacian on forms

Since this is more differential geometry than differential analysis I will be brief. We already defined the exterior derivative

$$
\begin{equation*}
d: \mathcal{C}^{\infty}(X) \longrightarrow \mathcal{C}^{\infty}\left(X ; T^{*} X\right) \tag{5.59}
\end{equation*}
$$

as part of the definition of $T^{*} X$, i.e.

$$
\begin{equation*}
d f(\bar{x})=\left[f(x-f(\bar{x})] \in T_{\bar{x}}^{*} X \forall \bar{x} \in X\right. \tag{5.60}
\end{equation*}
$$

The importance of the form bundles is tha they give a resolution of $\left(\frac{15.11}{5.59}\right)$.
[Actually I didn't mean to be this brief, but was interupted while writing the lecture!]

### 5.6. Hodge theorem

S.pseudo.proj
5.6.1998.228
4.6 .1998 .227 there exists $\Pi \in \Psi^{0}(M ; E)$ such that $\Pi^{2}=\Pi$ and $\Pi-P \in \Psi^{-\infty}(M ; E)$.

Proposition 5.6. If $P \in \Psi^{0}(M ; E)$ is such that $P^{2}-P \in \Psi^{-\infty}(M ; E)$ and $F \subset H^{s}(M ; E)$ is a closed subspace corresponding to which there are smoothing operators $A, B \in \Psi^{-\infty}(M ; E)$ with $\mathrm{Id}-P=A$ on $F$ and $(P+B) L^{2}(M ; E) \subset F$ then there is a smoothing operator $B^{\prime} \in \Psi^{-\infty}(M: E)$ such that $F=\operatorname{Ran}\left(P+B^{\prime}\right)$ and $\left(P+B^{\prime}\right)^{2}=P+B^{\prime}$.

Proof. Assume first that $s=0$, so $F$ is a closed subspace of $L^{2}(X ; E)$. Applying Proposition 5.5 to $P$ we may assume that it is a projection, without affecting the other conditions. Consider the intersection $E=F \cap \operatorname{Ran}(\operatorname{Id}-\Pi)$. This is a closed subspace of $L^{2}(M ; E)$. With $A$ as in the statement of the proposition, $E \subset \operatorname{Nul}(\operatorname{Id}-A)$. Indeed $P$ vanishes on $\operatorname{Ran}(\operatorname{Id}-P)$ and hence on $E$ and by hypothesis Id $-P-A$ vanishes on $F$ and hence on $E$. From the properties of smoothing operators, $E$ is contained in a finite dimensional subspace of $\mathcal{C}^{\infty}(M ; E)$, so is itself such a space. We may modify $P$ by adding a smoothing projection onto $E$ to it, and so assume that $F \cap \operatorname{Ran}(\operatorname{Id}-P)=\{0\}$.

Consider the sum $G=F+\operatorname{Ran}(\operatorname{Id}-P)$. Consider the operator $\operatorname{Id}+B=(P+$ $B)+(\operatorname{Id}-P)$, with $B$ as in the statement of the Proposition. The range of $\operatorname{Id}+B$ is contained in $G$. Thus $G$ must be a closed subspace of $L^{2}(M ; E)$ with a finite dimensional complement in $\mathcal{C}^{\infty}(M ; E)$. Adding a smoothing projection onto such a complement we can, again by altering $P$ by smoothing term, arrange that

$$
\begin{equation*}
L^{2}(M ; E)=F \oplus \operatorname{Ran}(\operatorname{Id}-P) \tag{5.61}
\end{equation*}
$$

is a (possibly non-orthogonal) direct sum. Since $P$ has only been altered by a smoothing operator the hypotheses of the Proposition continue to hold. Let $\Pi$ be the projection with range $F$ and null space equal to the range of Id $-P$. It follows that $P^{\prime}=P+(\operatorname{Id}-P) R P$ for some bounded operator $R$ (namely $R=$ $(\operatorname{Id}-P)\left(P^{\prime}-P\right) P$.) Then restricted to $F, P^{\prime}=\operatorname{Id}$ and $P=\operatorname{Id}+A$ so $R=-A$ on $F$. In fact $R=A P \in \Psi^{-\infty}(M ; E)$, since they are equal on $F$ and both vanish on $\operatorname{Ran}(\operatorname{Id}-P)$. Thus $P^{\prime}$ differs from $P$ by a smoothing operator.

The case of general $s$ follows by conjugating with a pseudodifferential isomorphism of $H^{s}(M ; E)$ to $L^{2}(M ; E)$ since this preserves both the assumptions and the conclusions.

### 5.8. Heat kernel

### 5.9. Resolvent

### 5.10. Complex powers

### 5.11. Problems

6.3 .1998 .156

## t

 tion $\frac{13.4}{5.1}$ is an equivalence relation on $\mathcal{C}^{\infty}$ structures. Conclude that there is a unique maximal $\mathcal{C}^{\infty}$ structure containing any give $\mathcal{C}^{\infty}$ structure.6.3 .1998 .158

Problem 5.2. Let $\mathcal{F}$ be a $\mathcal{C}^{\infty}$ structure on $X$ and let $O_{a}, a_{13} \in A$, be a covering of $X$ by coordinate neighbourhoods, in the sense of $\left(\frac{13.2}{5.2}\right)$ and $\left(\frac{13.3}{5.3}\right)$. Show that the maximal $\mathcal{C}_{113.3}^{\infty}$ structure containing $\mathcal{F}$ consists of ALL functions on $X$ which are of the form $\left(\frac{5.3}{}\right)$ on each of these coordinate patches. Conclude that the maximal $\mathcal{C}^{\infty}$ structure is an algebra.
6.3 .1998 .159

Problem 5.3 (Partitions of unity). Show that any $\mathcal{C}^{\infty}$ manifold admits partitions of unity. That is, if $O_{a}, a \in A$, is an open cover of $X$ then there exist elements $\rho_{a, i} \in \mathcal{C}^{\infty}(X), a \in A, i \in \mathbb{N}$, with $0 \leq \rho_{a, i} \leq 1$, with each $\rho_{a, i}$ vanishing outside a compact subset $K_{a, i} \subset O_{a}$ such that only finite collections of the $\left\{K_{a, i}\right\}$ have non-trivial intersection and for which

$$
\sum_{a \in A, i \in \mathbb{N}} \rho_{a, i}=1
$$

## CHAPTER 6

C.Elliptic.boundary

Elliptic boundary problems


#### Abstract

Summary Elliptic boundary problems are discussed, especially for operators of Dirac type. We start with a discussion of manifolds with boundary, including functions spaces and distributions. This leads to the 'jumps formula' for the relationship of the action of a differential operator to the operation of cutting off at the boundary; this is really Green's formula. The idea behind Calderòn's approach to boundary problems is introduced in the restricted context of a dividing hypersurface in a manifold without boundary. This includes the fundamental result on the boundary behaviour of a pseudodifferential operator with a rational symbol. These ideas are then extended to the case of an operator of Dirac type on a compact manifold with boundary with the use of left and right parametrices to define the Calderòn projector. General boundary problems defined by pseudodifferential projections are discussed by reference to the 'Calderòn realization' of the operator. Local boundary conditions, and the corresponding ellipticity conditions, are then discussed and the special case of Hodge theory on a compact manifold with boundary is analysed in detail for absolute and relative boundary conditions.


## Introduction

Elliptic boundary problems arise from the fact that elliptic differential operators on compact manifolds with boundary have infinite dimensional null spaces. The main task we carry out below is the parameterization of this null space, in terms of boundary values, of an elliptic differential operator on a manifold with boundary. For simplicity of presentation the discussion of elliptic boundary problems here will be largely confined to the case of first order systems of differential operators of Dirac type. This has the virtue that the principal results can be readily stated.

## Status as of 4 August, 1998

 sections are still missing.

### 6.1. Manifolds with boundary

Smooth manifolds with boundary can be defined in very much the same was as manifolds without boundary. Thus we start with a paracompact Hausdorff space $X$ and assume that it is covered by 'appropriate' coordinate patches with corresponding transition maps. In this case the 'model space' is $\mathbb{R}^{n, 1}=[0, \infty) \times \mathbb{R}^{n-1}$, a Euclidean half-space of fixed dimension, $n$. As usual it is more convenient to use
as models all open subsets of $\mathbb{R}^{n, 1}$; of course this means relatively open, not open as subsets of $\mathbb{R}^{n}$. Thus we allow any

$$
O=O^{\prime} \cap \mathbb{R}^{n, 1}, \quad O^{\prime} \subset \mathbb{R}^{n} \text { open }
$$

as local models.
By a smooth map between open sets in this sense we mean a map with a smooth extension. Thus if $O_{i}$ for $i=1,2$ are open in $\mathbb{R}^{n, 1}$ then smoothness of a map $F$ means that

$$
\begin{align*}
& F: O_{1} \rightarrow O_{2}, \exists O_{i}^{\prime} \subset \mathbb{R}^{n}, i=1,2, \text { open and } \tilde{F}: O_{1}^{\prime} \rightarrow O_{2}^{\prime}  \tag{6.1}\\
& \quad \text { which is } \mathcal{C}^{\infty} \text { with } O_{i}=O_{i}^{\prime} \cap \mathbb{R}^{n, 1} \text { and } F=F^{\prime} \mid O_{1}
\end{align*}
$$

It is important to note that the smoothness condition is much stronger than just smoothness of $F$ on $O \cap(0, \infty) \times \mathbb{R}^{n-1}$.

By a diffeomorphism between such open sets we mean an invertible smooth map with a smooth inverse. Various ways of restating the condition that a map be a diffeomorphism are discussed below.

With this notion of local model we define a coordinate system (in the sense of manifolds with boundary) as a homeomorphism of open sets

$$
X \supset U \xrightarrow{\Phi} O \subset \mathbb{R}^{n, 1}, \quad O, U \text { open. }
$$

Thus $\Phi^{-1}$ is assumed to exist and both $\Phi$ and $\Phi^{-1}$ are assumed to be continuous. The compatibility of two such coordinate systems $\left(U_{1}, \Phi_{1}, O_{1}\right)$ and $\left(U_{2}, \Phi_{2}, O_{2}\right)$ is the requirement that either $U_{1} \cap U_{2}=\phi$ or if $U_{1} \cap U_{2} \neq \phi$ then

$$
\Phi_{1,2}=\Phi_{2} \circ \Phi_{1}^{-1}: \Phi_{1}\left(U_{1} \cap U_{2}\right) \rightarrow \Phi_{2}\left(U_{1} \cap U_{2}\right)
$$

is a diffeomorphism in the sense described above. Notice that both $\Phi_{1}\left(U_{1} \cap U_{2}\right)$ and $\Phi_{2}\left(U_{1} \cap U_{2}\right)$ are open in $\mathbb{R}^{n, 1}$. The inverse $\Phi_{1,2}$ is defined analogously.

A $\mathcal{C}^{\infty}$ manifold with boundary can then be formally defined as a paracompact Hausdorff topological space which has a maximal covering by mutually compatible coordinate systems.

An alternative definition, i.e. characterization, of a manifold with boundary is that there exists a $\mathcal{C}^{\infty}$ manifold $\tilde{X}$ without boundary and a function $f \in \mathcal{C}^{\infty}(\tilde{X})$ such that $d f \neq 0$ on $\{f=0\} \subset \tilde{X}$ and

$$
X=\{p \in \tilde{X} ; f(p) \geq 0\}
$$

with coordinate systems obtained by restriction from $\tilde{X}$. The doubling construction described below shows that this is in fact an equivalent notion.

### 6.2. Smooth functions

As in the boundaryless case, the space of functions on a compact manifold with boundary is the primary object of interest. There are two basic approaches to defining local smoothness, the one intrinsic and the other extrinsic, in the style of the two definitions of a manifold with boundary above. Thus if $O \subset \mathbb{R}^{n, 1}$ is open we can simply set

$$
\begin{aligned}
\mathcal{C}^{\infty}(O)=\left\{u: O \rightarrow \mathbb{C} ; \exists \tilde{u} \in \mathcal{C}^{\infty}\left(O^{\prime}\right)\right. & \\
& \left.O^{\prime} \subset \mathbb{R}^{n} \text { open, } \quad O=O^{\prime} \cap \mathbb{R}^{n, 1}, u=\left.\tilde{u}\right|_{O}\right\} .
\end{aligned}
$$

Here the open set in the definition might depend on $u$. The derivatives of $\tilde{u} \in$ $\mathcal{C}^{\infty}\left(O^{\prime}\right)$ are bounded on all compact subsets, $K \Subset 0$. Thus

$$
\begin{equation*}
\sup _{K \cap O^{\circ}}\left|D^{\alpha} u\right|<\infty, \quad O^{\circ}=O \cap\left((0, \infty) \times \mathbb{R}^{n-1}\right) \tag{6.2}
\end{equation*}
$$

The second approach is to use ( $\frac{(\mathrm{eq}: \mathrm{F} \cdot \mathrm{F})}{}$ as a definition, i.e. to set

$$
\begin{equation*}
\mathcal{C}^{\infty}(O)=\left\{u: O^{\circ} \rightarrow \mathbb{C} ;\left(\frac{\mathrm{eq}: \mathrm{F}: \mathrm{F} 1}{6.2)} \text { holds } \forall K \Subset O \text { and all } \alpha\right\}\right. \tag{6.3}
\end{equation*}
$$

In particular this implies the continuity of $u \in \mathcal{C}^{\infty}(O)$ up to any point $p \in O \cap$ $\left(\{0\} \times \mathbb{R}^{n-1}\right)$, the boundary of $O$ as a manifold with boundary.

As the notation here asserts, these two approaches are equivalent. This follows (as does much more) from a result of Seeley:

Proposition 6.1. If $\mathcal{C}^{\infty}(O)$ is defined by ( $\frac{(\mathrm{eq}: \mathrm{F} 2}{\mathrm{F} .3}$ ) and $O^{\prime} \subset \mathbb{R}^{n}$ is open with $O=O^{\prime} \cap \mathbb{R}^{n, 1}$ then there is a linear extension map

$$
E: \mathcal{C}^{\infty}(O) \rightarrow \mathcal{C}^{\infty}\left(O^{\prime}\right),\left.\quad E u\right|_{O^{\prime}}=u
$$

which is continuous in the sense that for each $K^{\prime} \Subset O^{\prime}$, compact, there is some $K \Subset O$ such that for each $\alpha$

$$
\sup _{K^{\prime}}\left|D^{\alpha} E u\right| \leq C_{\alpha, K^{\prime}} \sup _{K \cap O}\left|D^{\alpha} u\right|
$$

The existence of such an extension map shows that the definition of a diffeomorphism of open sets $O_{1}, O_{2}$, given above, is equivalent to the condition that the map be invertible and that it, and its inverse, have components which are in $\mathcal{C}^{\infty}\left(O_{1}\right)$ and $\mathcal{C}^{\infty}\left(O_{2}\right)$ respectively.

Given the local definition of smoothness, the global definition should be evident. Namely, if $X$ is a $\mathcal{C}^{\infty}$ manifold with boundary then

$$
\mathcal{C}^{\infty}(X)=\left\{u: X \rightarrow \mathbb{C} ;\left(\Phi^{-1}\right)^{*}\left(\left.u\right|_{U}\right) \in \mathcal{C}^{\infty}(O) \forall \text { coordinate systems }\right\}
$$

This is also equivalent to demanding that local regularity, i.e. the regularity of $\left(\Phi^{-1}\right)^{*}\left(\left.u\right|_{O}\right)$, hold for any one covering by admissible coordinate systems.
 In fact the proof of Lemma $\frac{13}{5} .1$ applies verbatim; see also Problem b.3.

The topology of $\mathcal{C}^{\infty}(X)$ is given by the supremum norms of the derivatives in local coordinates. Thus a seminorm

$$
\sup _{K \Subset O}\left|D^{\alpha}\left(\Phi^{-1}\right)^{*}\left(\left.u\right|_{U}\right)\right|
$$

arises for each compact subset of each coordinate patch. In fact there is a countable set of norms giving the same topology. If $X$ is compact, $\mathcal{C}^{\infty}(X)$ is a Fréchet space, if it is not compact it is an inductive limit of Fréchet spaces (an LF space).

The boundary of $X, \partial X$, is the union of the $\Phi^{-1}\left(O \cap\left(\{0\} \times \mathbb{R}^{n-1}\right)\right)$ over coordinate systems. It is a manifold without boundary. It is compact if $X$ is compact. Furthermore, $\partial X$ has a global defining function $\rho \in \mathcal{C}^{\infty}(X)$; that is $\rho \geq 0, \partial X=\{\rho=0\}$ and $d \rho \neq 0$ at $\partial X$. Moreover if $\partial X$ is compact then any such boundary defining function can be extended to a product decomposition of $X$ near $\partial X$ :
(6.4) $\exists C \supset \partial X$, open in $X \epsilon>0$ and a diffeomorphism $\varphi: C \simeq[0, \epsilon)_{\rho} \times \partial X$.

If $\partial X$ is not compact this in still 19.1998 .251 for an appropriate choice of $\rho$. For an outline of proofs see Problem $\frac{11.6}{6.1}$

Problem 6.1.
The existence of such a product decomposition near the boundary (which might have several components) allows the doubling construction mentioned above to be carried through. Namely, let

$$
\begin{equation*}
\tilde{X}=(X \cup X) / \partial X \tag{6.5}
\end{equation*}
$$

be the disjoint union of two copies of $X$ with boundary points identified. Then consider

$$
\begin{align*}
\mathcal{C}^{\infty}(\tilde{X})= & \left\{\left(u_{1}, u_{2}\right) \in \mathcal{C}^{\infty}(X) \oplus \mathcal{C}^{\infty}(X) ;\right.  \tag{6.6}\\
& \left(\varphi^{-1}\right)^{*}\left(\left.u_{1}\right|_{C}\right)=f(\rho, \cdot),\left(\varphi^{-1}\right)^{*}\left(\left.u_{2}\right|_{C}\right)=f(-\rho, \cdot) \\
& \left.f \in \mathcal{C}^{\infty}((-1,1) \times \partial X)\right\}
\end{align*}
$$

This is a $\mathcal{C}^{\infty}$ structure on $\tilde{X}$ such that $X \hookrightarrow \tilde{X}$, as the first term in $\left(\frac{\mathrm{eq}: \mathrm{F4}}{6.5)}\right.$, is an embedding as a submanifold with boundary, so

$$
\mathcal{C}^{\infty}(X)=\left.\mathcal{C}^{\infty}(\tilde{X})\right|_{X}
$$

In view of this possibility of extending $X$ to $\tilde{X}$, we shall not pause to discuss all the usual 'natural' constructions of tensor bundles, density bundles, bundles of differential operators, etc. They can simply be realized by restriction from $\tilde{X}$. In practice it is probably preferable to use intrinsic definitions.

The definition of $\mathcal{C}^{\infty}(X)$ implies that there is a well-defined restriction map

$$
\left.\mathcal{C}^{\infty}(X) \ni u \longmapsto u\right|_{\partial X} \in \mathcal{C}^{\infty}(\partial X)
$$

It is always surjective. Indeed the existence of a product decomposition shows that any smooth function on $\partial X$ can be extended locally to be independent of the chosen normal variable, and then cut off near the boundary.

There are important points to observe in the description of functions near the boundary. We may think of $\mathcal{C}^{\infty}(X) \subset \mathcal{C}^{\infty}\left(X^{\circ}\right)$ as a subspace of the smooth functions on the interior of $X$ which describes the 'completion' (compactification if $X$ is compact!) of the interior to a manifold with boundary. It is in this sense that the action of a differential operator $P \in \operatorname{Diff}^{m}(X)$

$$
P: \mathcal{C}^{\infty}(X) \rightarrow \mathcal{C}^{\infty}(X)
$$

should be understood. Thus $P$ is just a differential operator on the interior of $X$ with 'coefficients smooth up to the boundary.'

Once this action is understood, there is an obvious definition of the space of $\mathcal{C}^{\infty}$ functions which vanish to all orders at the boundary,

$$
\dot{\mathcal{C}}^{\infty}(X)=\left\{u \in \mathcal{C}^{\infty}(X) ;\left.P u\right|_{\partial X}=0 \forall P \in \operatorname{Diff}^{*}(X)\right\}
$$

Having chosen a product decomposition near the boundary, Taylor's theorem gives us an isomorphism

$$
\mathcal{C}^{\infty}(X) / \dot{\mathcal{C}}^{\infty}(X) \cong \bigoplus_{k \geq 0} \mathcal{C}^{\infty}(\partial X) \cdot\left[\left.d \rho\right|_{\partial X}\right]^{k}
$$

### 6.3. Distributions

It is somewhat confusing that there are three (though really only two) spaces of distributions immediately apparent on a compact manifold with boundary. Understanding the relationship between them is important to the approach to boundary problems used here.

The coarsest (as it is a little dangerous to say largest) space is $\mathcal{C}^{-\infty}\left(X^{\circ}\right)$, the dual of $\mathcal{C}_{c}^{\infty}\left(X^{\circ} ; \Omega\right)$, just the space of distributions on the interior of $X$. The elements of $\mathcal{C}^{-\infty}\left(X^{\circ}\right)$ may have unconstrained growth, and unconstrained order of singularity, approaching the boundary. They are not of much practical value here and appear for conceptual reasons.

Probably the most natural space of distributions to consider is the dual of $\mathcal{C}^{\infty}(X ; \Omega)$ since this is arguably the direct analogue of the boundaryless case. We shall denote this space

$$
\begin{equation*}
\dot{\mathcal{C}}^{-\infty}(X)=\left(\mathcal{C}^{\infty}(X ; \Omega)\right)^{\prime} \tag{6.7}
\end{equation*}
$$

and call it the space of supported distributions. The 'dot' is supposed to indicate this support property, which we proceed to describe.

If $\tilde{X}$ is any compact extension of $X$ (for example the double) then, as already noted, the restriction map $\mathcal{C}^{\infty}(\tilde{X} ; \Omega) \rightarrow \mathcal{C}^{\infty}(X ; \Omega)$ is continuous and surjective. Thus, by duality, we get an injective 'extension' map

$$
\begin{equation*}
\dot{\mathcal{C}}^{-\infty}(X) \ni u \mapsto \tilde{u} \in \mathcal{C}^{-\infty}(\tilde{X}), \tilde{u}(\varphi)=u\left(\left.\varphi\right|_{X}\right) \tag{6.8}
\end{equation*}
$$

We shall regard this injection as an identification $\dot{\mathcal{C}}^{-\infty}(X) \hookrightarrow \mathcal{C}^{-\infty}(\tilde{X})$; its range is easily characterized.
 those $\tilde{u} \in \mathcal{C}^{-\infty}(\tilde{X})$ with $\operatorname{supp} \tilde{u} \subset X$.

The proof is given below. This proposition is the justification for calling $\dot{\mathcal{C}}^{-\infty}(X)$ the space of supported distributions; the dot is support to indicate that this is the subspace of the 'same' space for $\tilde{X}$, i.e. $\mathcal{C}^{-\infty}(\tilde{X})$, of elements with support in $X$.

This notation is consistent with $\dot{\mathcal{C}}^{\infty}(X) \subset \mathcal{C}^{\infty}(\tilde{X})$ being the subspace (by extension as zero) of elements with support in $X$. The same observation applies to sections of any vector bundle, so

$$
\dot{\mathcal{C}}^{\infty}(X ; \Omega) \subset \mathcal{C}^{\infty}(\tilde{X} ; \Omega)
$$

is a well-defined closed subspace. We set

$$
\begin{equation*}
\mathcal{C}^{-\infty}(X)=\left(\dot{\mathcal{C}}^{\infty}(X ; \Omega)\right)^{\prime} \tag{6.9}
\end{equation*}
$$

and call this the space of extendible distributions on $X$. The inclusion map for the test functions gives by duality a restriction map:
(6.10) $\quad R_{X}: \mathcal{C}^{-\infty}(\tilde{X}) \rightarrow \mathcal{C}^{-\infty}(X)$,

$$
R_{X} u(\varphi)=u(\varphi) \forall \varphi \in \dot{\mathcal{C}}^{\infty}(X ; \Omega) \hookrightarrow \mathcal{C}^{\infty}(\tilde{X} ; \Omega)
$$

We write, at least sometimes, $R_{X}$ for the map since it has a large null space so should not be regarded as an identification. In fact

$$
\begin{equation*}
\operatorname{Nul}\left(R_{X}\right)=\left\{v \in \mathcal{C}^{-\infty}(\tilde{X}) ; \operatorname{supp}(v) \cap X^{\circ}=\phi\right\}=\dot{\mathcal{C}}^{-\infty}\left(\tilde{X} \backslash X^{\circ}\right) \tag{6.11}
\end{equation*}
$$

is just the space of distributions supported 'on the other side of the boundary'. The primary justification for calling $\mathcal{C}^{-\infty}(X)$ the space of extendible distributions is:

Proposition 6.3. If $X$ is a compact manifold with boundary, then the space $\mathcal{C}_{c}^{\infty}\left(X^{\circ} ; \Omega\right)$ is dense in $\dot{\mathcal{C}}^{\infty}(X ; \Omega)$ and hence the restriction map

$$
\begin{equation*}
\mathcal{C}^{-\infty}(X) \hookrightarrow \mathcal{C}^{-\infty}\left(X^{\circ}\right) \tag{6.12}
\end{equation*}
$$

is injective, whereas the restriction map from (皆:ID), $R_{X}: \dot{\mathcal{C}}^{-\infty}(X) \longrightarrow \mathcal{C}^{-\infty}(X)$, is surjective.

Proof. If $V$ is a real vector field on $\tilde{X}$ which is inward-pointing across the boundary then

$$
\exp (s V): \tilde{X} \rightarrow \tilde{X}
$$

is a diffeomorphism with $F_{s}(X) \subset X^{\circ}$ for $s>0$. Furthermore if $\varphi \in \mathcal{C}^{\infty}(\tilde{X})$ then $F_{s}^{*} \varphi \rightarrow \varphi$ in $\mathcal{C}^{\infty}(\tilde{X})$ as $s \rightarrow 0$. The support property shows that $F_{s}^{*} \varphi \in \mathcal{C}_{c}^{\infty}\left(X^{\circ}\right)$ if $s<0$ and $\varphi \in \dot{\mathcal{C}}^{\infty}(X)$. This shows the density of $\mathcal{C}_{c}^{\infty}\left(X^{\circ}\right)$ in $\dot{\mathcal{C}}^{\infty}(X)$. Since all topologies are uniform convergence of all derivatiyes in open sets. The same argument applies to densities. The injectivity of ( 16.12 ) follows by duality.

On the other hand the surjectivity of ( $\overline{\mathrm{F} .10)}$ follows directly from the HahnBanach theorem.

Proof of Proposition $\frac{\text { prop: D } 4}{6.2 \text {. For }} \tilde{u} \in \mathcal{C}^{-\infty}(\tilde{X})$ the condition that supp $\tilde{u} \subset X$ is just

$$
\begin{equation*}
\tilde{u}(\varphi)=0 \forall \varphi \in \mathcal{C}_{c}^{\infty} \subset(\tilde{X} \backslash X ; \Omega) \subset \mathcal{C}^{\infty}(\tilde{X} ; \Omega) \tag{6.13}
\end{equation*}
$$

 holds, then by continuity and the density of $\mathcal{C}_{c}^{\infty}(\tilde{X} \backslash X ; \Omega)$ in $\mathcal{C}^{\infty}\left(\tilde{X} \backslash X^{\circ} ; \Omega\right)$, what follows from Proposition prop:D8

It is sometimes useful to consider topologies on the spaces of distributions $\mathcal{C}^{-\infty}(X)$ and $\dot{\mathcal{C}}^{-\infty}(X)$. For example we may consider the weak topology. This is given by all the seminorms $u \mapsto\|\langle u, \phi\rangle\|$, where $\phi$ is a test function.

Lemma 6.2. With respect to the weak topology, the subspace $\mathcal{C}_{c}^{\infty}\left(X^{\circ}\right)$ is dense in both $\mathcal{C}^{-\infty}(X)$ and $\mathcal{C}^{-\infty}(X)$.

### 6.4. Boundary Terms

To examine the precise relationship between the supported and extendible distributions consider the space of 'boundary terms'.

$$
\begin{equation*}
\dot{\mathcal{C}}_{\partial X}^{-\infty}(X)=\left\{u \in \dot{\mathcal{C}}^{-\infty}(X) ; \operatorname{supp}(u) \subset \partial X\right\} \tag{6.14}
\end{equation*}
$$

Here the support may be computed with respect to any extension, or intrinsically on $X$. We may also define a map 'cutting off' at the boundary:

$$
\begin{equation*}
\mathcal{C}^{\infty}(X) \ni u \mapsto u_{c} \in \dot{\mathcal{C}}^{-\infty}(X), u_{c}(\varphi)=\int_{X} u \varphi \forall \varphi \in \mathcal{C}^{\infty}(X ; \Omega) \tag{6.15}
\end{equation*}
$$

Proposition 6.4. If $X$ is a compact manifold with boundary then there is a commutative diagram

with the horizontal sequence exact.
Proof. The commutativity of the triangle follows directly from the definitions. The exactness of the horizontal sequence follows from the density of $\mathcal{C}_{c}^{\infty}\left(X^{\circ} ; \Omega\right)$ in $\dot{\mathcal{C}}^{\infty}(X ; \Omega)$. Indeed, this shows that $v \in \dot{\mathcal{C}}_{\partial X}^{-\infty}(X)$ maps to 0 in $\mathcal{C}^{-\infty}(X)$ since $v(\varphi)=0 \forall \varphi \in \mathcal{C}_{c}^{\infty}\left(X^{\circ} ; \Omega\right)$. Similarly, if $u \in \dot{\mathcal{C}}^{-\infty}(X)$ maps to zero in $\mathcal{C}^{-\infty}(X)$ then $u(\varphi)=0$ for all $\varphi \in \mathcal{C}_{c}^{\infty}\left(X^{\circ} ; \Omega\right)$, so $\operatorname{supp}(u) \cap X^{\circ}=\emptyset$.

Note that both maps in ( $\frac{{ }^{\frac{B T}{6} 4} 6}{6.16}$ ) from $\mathcal{C}^{\infty}(X)$ into supported and extendible distributions are injective. We regard the map into $\mathcal{C}^{-\infty}(X)$ as an identification. In particular this is consistent with the action of differential operators. Thus $P \in$ Diff ${ }^{m}(X)$ acts on $\mathcal{C}^{\infty}(X)$ and then the smoothness of the coefficients of $P$ amount to the fact that it preserves $\mathcal{C}^{\infty}(X)$, as a subspace. The formal adjoint $P^{*}$ with respect to the sesquilinear pairing for some smooth positive density, $\nu$

$$
\begin{equation*}
\langle\varphi, \psi\rangle=\int_{X} \varphi \bar{\psi} \nu \quad \forall \varphi, \psi \in \mathcal{C}^{\infty}(X) \tag{6.17}
\end{equation*}
$$

acts on $\dot{\mathcal{C}}^{\infty}(X)$ :

$$
\begin{equation*}
\left\langle P^{*} \varphi, \psi\right\rangle=\langle\varphi P \psi\rangle \quad \forall \varphi \in \dot{\mathcal{C}}^{\infty}(X), \psi \in \mathcal{C}^{\infty}(X), P^{*}: \dot{\mathcal{C}}^{\infty}(X) \longrightarrow \dot{\mathcal{C}}^{\infty}(X) \tag{6.18}
\end{equation*}
$$

However, $P^{*} \in \operatorname{Diff}^{m}(X)$ is fixed by its action over $X^{\circ}$. Thus we do have

$$
\begin{equation*}
\left\langle P^{*} \varphi, \psi\right\rangle=\langle\varphi, P \psi\rangle \quad \forall \varphi \in \mathcal{C}^{\infty}(X), \psi \in \dot{\mathcal{C}}^{\infty}(X) \tag{6.19}
\end{equation*}
$$

We define the action of $P$ by duality. In view of the possibility of confusion, we denote $P$ the action on $\mathcal{C}^{-\infty}(X)$ and by $\dot{P}$ the action on $\dot{\mathcal{C}}^{\infty}(X)$.

$$
\begin{align*}
& \langle P u, \varphi\rangle=\left\langle u, P^{*} \varphi\right\rangle \quad \forall u \in \mathcal{C}^{-\infty}(X), \varphi \in \dot{\mathcal{C}}^{\infty}(X), P: \mathcal{C}^{-\infty}(X) \longrightarrow \mathcal{C}^{-\infty}(X)  \tag{6.20}\\
& \langle\dot{P} u, \varphi\rangle=\left\langle u, P^{*} \varphi\right\rangle \quad \forall u \in \dot{\mathcal{C}}^{-\infty}(X), \varphi \in \mathcal{C}^{\infty}(X), \dot{P}: \dot{\mathcal{C}}^{-\infty}(X) \longrightarrow \dot{\mathcal{C}}^{-\infty}(X)
\end{align*}
$$

It is of fundamental importance that $\frac{(\mathbb{B T 7} 7}{5.19)}$ does not hold for all $\varphi, \psi \in \mathcal{C}^{\infty}(X)$. This failure is reflected in Green's formula for the boundary terms, which appears below as the 'Jump formula'. This is a distributional formula for the difference

$$
\begin{equation*}
\dot{P} u_{c}-(P u)_{c} \in \dot{\mathcal{C}}_{\partial X}^{-\infty}, u \in \mathcal{C}^{\infty}(X) P \in \operatorname{Diff}^{m}(X) \tag{6.21}
\end{equation*}
$$

Recall that a product decomposition of $C \subset X$ near $\partial X$ is fixed by an inward pointing vector field $V$. Let $x \in \mathcal{C}^{\infty}(X)$ be a corresponding boundary defining function, with $V x=0$ near $\partial X$, with $\chi_{V}: C \rightarrow \partial X$ the projection onto the
boundary from the product neighborhood $C$. Then Taylor's formula for $u \in \mathcal{C}^{\infty}(X)$ becomes

$$
\begin{equation*}
u \sim \sum_{k} \frac{1}{k!} \chi_{V}^{*}\left(\left.V^{k} u\right|_{\partial x}\right) x^{k} \tag{6.22}
\end{equation*}
$$

It has the property that a finite sum

$$
u_{N}=\varphi u-\varphi \sum_{k=0}^{N} \frac{1}{k!} \chi_{V}^{*}\left(\left.V^{k} u\right|_{\partial X}\right) x^{k}
$$

where $\varphi \equiv 1$ near $\partial X, \operatorname{supp} \varphi \subset C$, satisfies
BT11
(6.23) $\dot{P}\left(u_{N}\right)_{c}=\left(P u_{N}\right)_{c}, P \in \operatorname{Diff}^{m}(X), m<N$.
 course) only depend on the $\left.V^{k} u\right|_{\partial X}$ for $k \leq m$, in fact only for $k<m$.

Consider the Heaviside function $1_{c} \in \dot{\mathcal{C}}^{-\infty}(X)$, detained by cutting off the identity function of the boundary. We define distributions

$$
\begin{equation*}
\delta^{(j)}(x)=V^{j+1} 1_{c} \in \dot{\mathcal{C}}_{\partial X}^{-\infty}, j \geq 0 \tag{6.24}
\end{equation*}
$$

Thus, $\delta^{(0)}(x)=\delta(x)$ is a 'Dirac delta function' at the boundary. Clearly supp $\delta(x) \subset$ $\partial X$, so the same is true of $\delta^{(j)}(x)$ for every $j$. If $\psi \in \mathcal{C}^{\infty}(\partial X)$ we define

$$
\begin{equation*}
\psi \cdot \delta^{(j)}(x)=\varphi\left(X_{V}^{*} \psi\right) \cdot \delta^{(j)}(x) \in \dot{\mathcal{C}}_{\partial X}^{-\infty}(X) \tag{6.25}
\end{equation*}
$$

This, by the support property of $\delta^{(j)}$, is independent of the cut off $\varphi$ used to define it.

Proposition 6.5. For each $P \in \operatorname{Diff}^{m}(X)$ there are differential operators on the boundary $P_{i j} \in \operatorname{Diff}^{m-i-j-1}(\partial X), i+j<m, i, j \geq 0$, such that

$$
\begin{equation*}
\dot{P} u_{c}-(P u)_{c}=\sum_{i, j}\left(P_{i j}\left(\left.V_{u}^{j}\right|_{\partial x}\right) \cdot \delta^{(j)}(x), \forall u \in \mathcal{C}^{\infty}(X),\right. \tag{6.26}
\end{equation*}
$$

and $P_{0 m-1}=i^{-m} \sigma(P, d x) \in \mathcal{C}^{\infty}(\partial X)$.
Proof. In the local product neighborhood $C$,

$$
\begin{equation*}
P=\sum_{0 \leq l \leq m} P_{l} V^{l} \tag{6.27}
\end{equation*}
$$

where $P_{l}$ is a differential operator of the order at most $m-l$, on $X$ be depending on $x$ as a parameter. Thus the basic cases we need to analyze arise from the application of $V$ to powers of $x$ :
(6.28) $x^{l}\left(V^{j+1}\left(x^{p}\right)_{c}-\left(V^{j+1} x^{p}\right)_{c}\right)$

$$
=\left\{\begin{array}{cc}
0 & l+p>j \\
\frac{p!(j-p)!}{(j-p-l)!} & (-1)^{l \delta^{(j-p-l)}}
\end{array} \quad l+p \leq j .\right.
$$

Taking the Taylor sense of the $P_{l}$,

$$
P_{l} \sim \sum_{r} x^{r} P_{l, r}
$$

and applying $P$ to $\left(\frac{B T 10}{6.22}\right)$ gives

$$
\begin{equation*}
P u_{c}-(P u)_{c}=\sum_{r+k<l}(-1)^{r}\left(P_{l, r}\left(\left.V^{k} u\right|_{\partial x}\right)\right) \quad \delta^{(l-1-r-k)}(x) \tag{6.29}
\end{equation*}
$$

This is of the form $\left(\frac{\mathrm{BT} 15}{6.26}\right)$. The only term with $l-1-r-k=m-1$ arises from $l-m, k=r=0$ so is the operator $P_{m}$ at $x=0$. This is just $i^{-m} \sigma(P, d x)$.

### 6.5. Sobolev spaces

As with $\mathcal{C}^{\infty}$ functions we may define the standard (extendible) Sobolev spaces by restriction or intrinsically. Thus, if $\tilde{X}$ is an extension of a compact manifold with boundary, $X$, the we can define

$$
\begin{equation*}
H^{m}(X)=H_{\mathrm{c}}^{m}(\tilde{X}) \mid X, \forall m \in \mathbb{R} ; H^{m}(X) \subset \mathcal{C}^{-\infty}(X) \tag{6.30}
\end{equation*}
$$

That this is independent of the choice of $\tilde{X}$ follows from the standard properties of the Sobolev spaces, particularly their localizability and invariance under diffeomorphisms. The norm in $H^{m}(X)$ can be taken to be
11.6.1998.254

$$
\begin{equation*}
\|u\|_{m}=\inf \left\{\|\tilde{u}\|_{H^{m}(\tilde{X})} ; \tilde{u} \in H^{m}(\tilde{X}), u=\tilde{u}_{X}\right\} \tag{6.31}
\end{equation*}
$$

A more intrinsic defintion of these spaces is discussed in the problems.
There are also supported Sobolev spaces,

$$
\begin{equation*}
\dot{H}^{m}(X)=\left\{u \in H^{m}(\tilde{X}) ; \operatorname{supp}(u) \subset X\right\} \subset \dot{\mathcal{C}}^{-\infty}(X) \tag{6.32}
\end{equation*}
$$

Sobolev space of sections of any vector bundle can be defined similarly.
Proposition 6.6. For any $m \in \mathbb{R}$ and any compact manifold with boundary $X, H^{m}(X)$ is the dual of $\dot{H}^{-m}(X ; \Omega)$ with respect to the continuous extension of the densely defined bilinear pairing

$$
(u, v)=\int_{X} u v, u \in \mathcal{C}^{\infty}(X), v \in \dot{\mathcal{C}}^{\infty}(X ; \Omega)
$$

Both $H^{m}(X)$ and $\dot{H}^{m}(X)$ are $\mathcal{C}^{\infty}(X)$-modules and for any vector bundle over $X, H^{m}(X ; E) \equiv H^{m}(X) \otimes_{\mathcal{C}^{\infty}(X)} \mathcal{C}^{\infty}(X ; E)$ and $\dot{H}^{m}(X ; E) \equiv \dot{H}^{m}(X) \otimes_{\mathcal{C}^{\infty}(X)}$ $\mathcal{C}^{\infty}(X ; E)$.

Essentially from the definition of the Sobolev spaces, any $P \in \operatorname{Diff}^{k}\left(X ; E_{1}, E_{2}\right)$ defines a continuous linear map

$$
\begin{equation*}
P: H^{m}\left(X ; E_{1}\right) \longrightarrow H^{m-k}\left(X ; E_{2}\right) . \tag{6.33}
\end{equation*}
$$

We write the dual (to $P^{*}$ of course) action
11.6 .1998 .258

$$
\begin{equation*}
\dot{P}: \dot{H}^{m}\left(X ; E_{1}\right) \longrightarrow \dot{H}^{m-k}\left(X ; E_{2}\right) \tag{6.34}
\end{equation*}
$$

These actions on Sobolev spaces are consistent with the corresponding actions on distributions. Thus

$$
\begin{aligned}
\mathcal{C}^{-\infty}(X ; E) & =\bigcup_{m} H^{m}(X), \mathcal{C}^{\infty}(X ; E)=\bigcap_{m} H^{m}(X), \\
\dot{\mathcal{C}}^{-\infty}(X ; E) & =\bigcup_{m} \dot{H}^{m}(X), \dot{\mathcal{C}}^{\infty}(X ; E)=\bigcap_{m} \dot{H}^{m}(X) .
\end{aligned}
$$

### 6.6. Dividing hypersurfaces

As already noted, the point of view we adopt for boundary problems is that they provide a parametrization of the space of solutions of a differential operator on a space with boundary. In order to clearly indicate the method pioneered by Calderòn, we shall initially consider the restrictive context of an operator of Dirac type on a compact manifold without boundary with an embedded separating hypersurface.

Thus, suppose initially that $D$ is an elliptic first order differential operator acting between sections of two (complex) vector bundles $V_{1}$ and $V_{2}$ over a compact manifold without boundary, $M$. Suppose further that $H \subset M$ is a dividing hypersurface. That is, $H$ is an embedded hypersuface with oriented (i.e. trivial) normal bundle and that $M=M_{+} \cup M_{-}$where $M_{ \pm}$are compact manifolds with boundary which intersect in their common boundary, $H$. The convention here is that $M_{+}$is on the positive side of $H$ with respect to the orientation.

In fact we shall make a further analytic assumption, that

$$
\begin{equation*}
D: \mathcal{C}^{\infty}\left(M ; V_{1}\right) \longrightarrow \mathcal{C}^{\infty}\left(M ; V_{2}\right) \text { is an isomorphism. } \tag{6.35}
\end{equation*}
$$

As we already know, $D$ is always Fredholm, so this implies the topological condition that the index vanish. However we only assume ( 6.35 ) to simplify the initial discussion.

Our objective is to study the space of solutions on $M_{+}$. Thus consider the map

$$
\begin{equation*}
\left\{u \in \mathcal{C}^{\infty}\left(M_{+} ; V_{1}\right) ; D u=0 \text { in } M_{+}^{\circ}\right\} \xrightarrow{b_{H}} \mathcal{C}^{\infty}\left(H ; V_{1}\right), b_{H} u=u_{\mid \partial M_{+}} \tag{6.36}
\end{equation*}
$$

The idea is to use the boundary values to parameterize the solutions and we can see immediately that this is possible.

Lemma 6.3. The assumption ( $\frac{11.4 .1998 .195}{6.35) ~ i m p l e s ~ t h a t ~ m a p ~} b_{H}$ in $\left(\frac{11.4 .1998 .196}{6.36)}\right.$ is injective.
Proof. Consider the form of $D$ in local coordinates near a point of $H$. Let the coordinates be $x, y_{1}, \ldots, y_{n-1}$ where $x$ is a local defining function for $H$ and assume that the coordinate patch is so small that $V_{1}$ and $V_{2}$ are trivial over it. Then

$$
D=A_{0} D_{x}+\sum_{j=1}^{n-1} A_{j} D_{y_{j}}+A^{\prime}
$$

where the $A_{j}$ and $A^{\prime}$ are local smooth bundle maps from $V_{1}$ to $V_{2}$. In fact the ellipticity of $D$ implies that each of the $A_{j}$ 's is invertible. Thus the equation can be written locally

$$
D_{x} u=B u, B=-\sum_{j=1}^{n-1} A_{0}^{-1} D_{y_{j}}-A_{0}^{-1} A^{\prime}
$$

The differential operator $B$ is tangent to $H$. By assumption $u$ vanishes when restricted to $H$ so it follows that $D_{x} u$ also vanishes at $H$. Differentiating the equation with respect to $x$, it follows that all derivatives of $u$ vanish at $H$. This in turn implies that the global section of $V_{1}$ over $M$

$$
\tilde{u}= \begin{cases}u & \text { in } M_{+} \\ 0 & \text { in } M_{-}\end{cases}
$$

is smooth and satisfies $D \tilde{u}=0$. Then assumption $\left(\frac{11.4 .1998 .195}{6.35)}\right.$ implies that $\tilde{u}=0$, so $u=0$ in $M_{+}$and $b_{H}$ is injective.

In the proof of this Lemma we have used the strong assumption ( $\frac{11.4 .1998 .195}{6.35) \cdot A s}$ we show below, if it is assumed instead that $D$ is of Dirac type then the Lemma remains true without assuming (6.35). Now we can state the basic result in this setting.
29.3.1998.187
29.3.1998.193
11.4.1998.198
11.4.1998.202
11.4.1998.203

Lemma 6.4. Let $D$ be an elliptic differential operator of first order on $M$, acting between vector bundles $V_{1}$ and $V_{2}$. If $u \in \mathcal{C}^{\infty}\left(M_{+} ; V_{1}\right)$ satisfies $D u=0$ in $M_{+}^{\circ}$ then

$$
\begin{equation*}
D u_{c}=\frac{1}{i} \sigma_{1}(D)(d x)\left(b_{H} u\right) \cdot \delta(x) \in \mathcal{C}^{-\infty}\left(M ; V_{2}\right) \tag{6.39}
\end{equation*}
$$

Since the same result is true for $M_{-}$, with an obvious change of sign, $D$ defines a linear operator

$$
\begin{array}{r}
D:\left\{u \in L^{1}\left(M ; V_{1}\right) ; u_{ \pm}=u \mid M_{ \pm} \in \mathcal{C}^{\infty}\left(M_{ \pm} ; V_{1}\right), D u_{ \pm}=0 \text { in } M_{ \pm}^{\circ}\right\} \longrightarrow  \tag{6.40}\\
\frac{1}{i} \sigma(D)(d x)\left(b_{H} u_{+}-b_{H} u_{-}\right) \cdot \delta(x) \in \mathcal{C}^{\infty}\left(H ; V_{2}\right) \cdot \delta(x) .
\end{array}
$$

To define the Calderòn projection we shall use the 'inverse' of this result.

$$
\begin{equation*}
\Pi_{C} v=b_{H}\left(I_{D} \frac{1}{i} \sigma(D)(d x) v \cdot \delta(x)\right)_{+}, v \in \mathcal{C}^{\infty}\left(H ; V_{1}\right) \tag{6.41}
\end{equation*}
$$

satisfies the conditions of Theorem $\frac{29.3}{6.1 .}$
Proof. Observe that the map $\left(\frac{11.4 .1998 .205}{6.40)}\right.$ is injective, since its null space consists of solutions of $D u=0$ globally on $M$; such a solution must be smooth by elliptic regularity and hence must yanish 8 s , the assumed invertibility of $D$. Thus the main task is to show that $D$ in ( $\overline{\mathbf{W} .40})$ is surjective.

Since $D$ is elliptic and, by assumption, an isomorphism on $\mathcal{C}^{\infty}$ sections ${ }_{11}$ Qver $_{1}{ }_{1998.205}$ $M$ it is also an isomorphism on distributional sections. Thus the inverse of ( $\overline{6} . \dot{4} \dot{0})$ must be given by $D^{-1}$. To prove the surjectivity it is enough to show that

ThEOREM 6.1. If $M=M_{+} \cup M_{-}$is a compact manifold without boundary with separating hypersurface $H$ as described above and $D \in \operatorname{Diff}^{1}\left(M ; V_{1}, V_{2}\right)$ is a generalized Dirac operator then there is an element $\Pi_{C} \in \Psi^{0}(H ; V), V=V_{1} \mid H$, satisfying $\Pi_{C}^{2}=\Pi_{C}$ and such that

$$
\begin{equation*}
b_{H}:\left\{u \in \mathcal{C}^{\infty}\left(M_{+} ; V_{1}\right) ; D u=0\right\} \longrightarrow \Pi_{C} \mathcal{C}^{\infty}(H ; V) \tag{6.37}
\end{equation*}
$$

is an isomorphism. The projection $\Pi_{C}$ can be chosen so that
(6.38) $\quad b_{H}:\left\{u \in \mathcal{C}^{\infty}\left(M_{-} ; V_{1}\right) ; D u=0\right\} \longrightarrow\left(\operatorname{Id}-\Pi_{C}\right) \mathcal{C}^{\infty}(H ; V)$
then $\Pi_{C}$ is uniquely determined and is called the Calderòn projection.
This result remains true for a general elliptic operator of first order if ${ }^{(11.4 .4}\left(\frac{11}{6.35}\right)$ is assumed, and even in a slightly weakened form without ( 6.35 ). Appropriate modifications to the proofs below are consigned to problems.

For first order operators the jump formula discussed above takes the following form.

Prgegestion 6.7. If $D \in \operatorname{Diff}^{1}\left(M ; V_{1}, V_{2}\right)$ is elliptic and satisfies $\left(\frac{11.4 .1998 .195}{6.35)}\right.$ then ( 16.40 ) is an 2 somorphism, with inverse $I_{D}$, and

$$
\begin{equation*}
D^{-1}(w \cdot \delta(x)) \mid M_{ \pm} \in \mathcal{C}^{\infty}\left(M_{ \pm} ; V_{1}\right) \forall w \in \mathcal{C}^{\infty}\left(H ; V_{2}\right) \tag{6.42}
\end{equation*}
$$

There can be no singular terms supported on $H$ since $w \cdot \delta(x) \in H^{-1}\left(M ; V_{2}\right)$ implies that $u=D^{-1}(w \cdot \delta(x)) \in L^{2}\left(M ; V_{1}\right)$.

Now, recalling that $D^{-1} \in \Psi^{-1}\left(M ; V_{2}, V_{1}\right)$, certainly $u$ is $\mathcal{C}^{\infty}$ away from $H$. At any point of $H$ outside the support of $w, u$ is also smooth. Since we may decompose $w$ using a partition of unity, it suffices to suppose that $w$ has support in a small coordinate patch, over which both $V_{1}$ and $V_{2}$ are trivial and to show that $u$ is smooth 'up to $H$ from both sides' in the local coordinate patch. Discarding smoothing terms from $D^{-1}$ we may therefore replace $D^{-1}$ by any local parametrix $Q$ for $D$ and work in local coordinates and with components:
$Q_{i j}\left(w_{j}(y) \cdot \delta(x)\right)=(2 \pi)^{-n} \int e^{i\left(x-x^{\prime}\right) \xi+i\left(y-y^{\prime}\right) \cdot \eta} q_{i j}(x, y, \xi, \eta) w\left(y^{\prime}\right) \delta\left(x^{\prime}\right) d x^{\prime} d y^{\prime} d \xi d \eta$
For a general pseudodifferential operator, even of order -1 , the result we are seeking is not true. We must use special properties of the symbol of $Q$, that is $D^{-1}$.
13.4 .1998 .210

### 6.7. Rational symbols

LEMMA 6.5. The left-reduced symbol of any local parametrix for a generalized Dirac operator has an expansion of the form

$$
\begin{equation*}
q_{i j}(z, \zeta)=\sum_{l=1}^{\infty} g(z, \zeta)^{-2 l+1} p_{i j, l}(z, \zeta) \text { with } p_{i j, l} \text { a polynomial of degree } 3 l-2 \text { in } \zeta \tag{6.44}
\end{equation*}
$$

here $g(z, \zeta)$ is the metric in local coordinates; each of the terms in $\left(\frac{13.4 .1998 .210}{6.44)}\right.$ ns therefore a symbol of order $-l$.

Proof. This follows by an inductive arument, of a now familiar type. First, the assumption that $D$ is a generalized Dirac operator means that its symbol $\sigma_{1}(D)(z, \zeta)$ has inverse $g(z, \zeta)^{-1} \sigma_{1}(D)^{*}(z, \zeta)$; this is the princiapl symbol of $Q$. Using Leibniz' formula one concludes that for any polynomial $r_{l}$ of degree $j$

$$
\partial_{\zeta_{i}}\left(g(z, \zeta)^{-2 l+1} r_{j}(z, \zeta)\right)=g(z, \zeta)^{-2 l} r_{j+1}^{\prime}(z, \zeta)
$$

where $r_{j+1}$ has degree (at most) $j+1$. Using this result repeatedly, and proceeding by induction, we may suppose that $q=q_{k}^{\prime}+q_{k+1}^{\prime \prime}$ where $q_{k}^{\prime}$ has an expansion up to order $k$, and so may be taken to be such a sum, and $q_{k+1}^{\prime \prime}$ is of order at most $-k-1$. The composition formula for left-reduced symbols then shows that

$$
\sigma_{1}(D) q_{k+1}^{\prime \prime} \equiv g^{-2 k} q_{k+1} \quad \bmod S^{-k-1}
$$

where $q_{k+1}$ is a polynomial of degree at most $3 k$. Inverting $\sigma_{1}(D)(\zeta)$ as at the initial step then shows that $q_{k+1}^{\prime \prime}$ is of the desired form, $g^{-2 k-1} r_{k+1}$ with $r_{k+1}$ of degree $3 k+1=3(k+1)-2$, modulo terms of lower order. This completes the proof of the lemma.

With this form for the symbol of $Q$ we proceed to the proof of Proposition $\frac{11.4 .1998 .204}{6.7 .1}$ That is, we consider ( $\overline{\mathbf{6} .43)}$. Since we only need to consider each term, we shall drop the indicies. A term of low order in the amplitude $q_{N}$ gives an operator with kernel in $\mathcal{C}^{N-d}$. Such a kernel gives an operator

$$
\mathcal{C}^{\infty}\left(H ; V_{2}\right) \longrightarrow \mathcal{C}^{N-d}\left(M ; V_{1}\right)
$$

 term in the expansion of the symbol $q$ gives an operator as in $\left(\frac{13.4 .19}{6.42}\right)$.

To be more precise, we can assume that the amplitude $q$ is of the form

$$
q=(1-\phi) g^{-2 l} q^{\prime}
$$

where $q^{\prime}$ is a polynomial of degree $3 l-2$ and $\phi=\phi(\xi, \eta)$ is a function of compact support which is identically one near the origin. The cutoff function is to remove the singularity at $\zeta=(\xi, \eta)=0$. Using continuity in the symbol topology the integrals in $x^{\prime}$ and $y^{\prime}$ can be carried out. By assumption $w \in \mathcal{C}_{c}^{\infty}\left(\mathbb{R}^{n-1}\right)$, so the resulting integral is absolutely convergent in $\eta$. If $l>1$ it is absolutely convergent in $\xi$ as well, so becomes

$$
Q(w(y) \cdot \delta(x))=(2 \pi)^{-n} \int e^{i x \xi+i y \cdot \eta} q(x, y, \xi, \eta) \hat{w}(\eta) d \xi d \eta
$$

In $|\xi|>1$ the amplitude is a rational function of $\xi$, decaying quadratically as $\xi \rightarrow \infty$. If we assume that $x>0$ then the exponential factor is bounded in the half plane $\Im \xi \geq 0$. This means that the limit as $R \rightarrow \infty$ over the integral in $\Im \xi \geq 0$ over the semicircle $|\xi|=R$ tends to zero, and does so with uniform rapid decrease in $\eta$. Cauchy's theorem shows that, for $R>1$ the real integral in $\xi$ can be replaced by the contour integral over $\gamma(R)$, which is, for $R \gg|\eta|$ given by the real interval $[-R, R]$ together with the semicircle of radius $R$ in the upper half plane. If $|\eta|>1$ the integrand is meromorphic in the upper half plane with a possible pole at the singular point $g(x, y, \xi, \eta)=0$; this is at the point $\xi=i r^{\frac{1}{2}}(x, y, \eta)$ where $r(x, y, \eta)$ is a positive-definite quadratic form in $\eta$. Again applying Cauchy's theorem

$$
Q\left(w(y) \delta(x)=(2 \pi)^{-n+1} i \int e^{x r^{\frac{1}{2}}(x, y, \eta)+i y \cdot \eta} q^{\prime}(x, y, \eta) \hat{w}(\eta) d \eta\right.
$$

where $q^{\prime}$ is a symbol of order $-k+1$ in $\eta$.
The product $e^{x r^{\frac{1}{2}}(x, y, \eta)} q^{\prime}(x, y, \eta)$ is uniformly a symbol of order $-k+1$ in $x>1$, with $x$ derivatives of order $p$ being uniformly symbols of order $-k+1+p$. It follows from the properties of pseudodifferential operators that $Q(w \cdot \delta(x))$ is a smooth function in $x>0$ with all derivatives locally uniformly bounded as $x \downarrow 0$.

### 6.8. Proofs of Proposition $\frac{11.4 .1998 .204}{6.7 \text { and Theorem } \frac{29.3 .1998 .187}{6.1}}$

This completes the proof of $\left(\frac{13.4,1998.207}{6.42}\right)$, since a with contour deformation into the lower half plane. Thus we have shown that $\left(\frac{11.4 .1998 .205}{614.4 .1998 .204}\right.$ somorphism which is the first half of the statement of Proposition $\left(\frac{1}{6.7}\right)$. Furthermore we see that the limiting value from above is a pseudodifferential operator on $H$ :

$$
\begin{equation*}
Q_{0} w=\lim _{x \downarrow 0} D^{-1}(w \cdot \delta(x)), Q_{0} \in \Psi^{0}\left(H ; V_{2}, V_{1}\right) \tag{6.45}
\end{equation*}
$$

This in turn implies that $\Pi_{C}$, defined by $\left(\frac{11.4 .1998 .206}{(6.41)}\right.$ is an element of $\Psi^{0}\left(H ; V_{1}\right)$, since it is $Q_{0} \circ \frac{1}{i} \sigma(D)(d x)$.

Next we check that $\Pi_{C}$ is a projection, i.e. that $\Pi_{C}^{2}=\Pi_{C}$. If $w=\Pi_{C} v$, $v \in \mathcal{C}^{\infty}\left(H ; V_{1}\right)$, then $w=b_{H} u, u=\left.I_{D} \frac{1}{4} \sigma(D)(d x) v\right|_{M_{+}}$, so $u \in \mathcal{C}^{\infty}\left(M_{+} ; V_{1}\right)$ satisfies $D u=0$ in $M_{+}^{\circ}$. In particular, by $\left(\frac{1.39}{6.39}\right), P u_{c}=\frac{1}{i} \sigma_{1}(D)(d x) w \cdot \delta(x)$, which means that $w=\Pi_{C} w$ so $\Pi_{C}^{2}=\Pi_{29: 3}$ Thig ${ }_{8}$ aldso shows that the range of $\Pi_{C}$ is precisely the range of $b_{H}$ as stated ing $(1.4 .3)$. The same argument shows that this choice of the projection gives ( $\overline{\mathbf{6} .3 \overline{8}}$ ).

### 6.9. Inverses

Still for the case of a generalized Dirac operator on a compact manifold with dividing hypersurface, consider what we have shown. The operator $D$ defines a
map in (11.4.1998.203
map in (6.39) with inverse

$$
\begin{equation*}
I_{D}:\left\{v \in \mathcal{C}^{\infty}\left(H ; V_{1}\right) ; \Pi_{C} v=v\right\} \longrightarrow\left\{u \in \mathcal{C}^{\infty}\left(M_{+} ; V_{1}\right) ; D u=0 \text { in } M_{+}\right\} \tag{6.46}
\end{equation*}
$$

This operator is the 'Poisson' operator for the canonical boundary condition given by the Calderòn operator, that is $u=I_{D} v$ is the unique solution of

$$
\begin{equation*}
D u=0 \text { in } M_{+}, u \in \mathcal{C}^{\infty}\left(M_{+} ; V_{1}\right), \Pi_{C} b_{H} u=v \tag{6.47}
\end{equation*}
$$

We could discuss the regularity properties of $I_{D}$ but we shall postpone this until after we have treated the 'one-sided' case of a genuine boundary problem.

As well as $I_{D}$ we have a natural right inverse for the operator $D$ as a map from $\mathcal{C}^{\infty}\left(M_{+} ; V_{1}\right)$ to $\mathcal{C}^{\infty}\left(M_{-} ; V_{2}\right)$. Namely

Lemma 6.6. If $f \in \mathcal{C}^{\infty}\left(M_{+} ; V_{2}\right)$ then $u=\left.D^{-1}\left(f_{c}\right)\right|_{M_{+}} \in \mathcal{C}^{\infty}\left(M_{+} ; V_{1}\right)$ and the map $R_{D}: f \longmapsto u$ is a right inverse for $D$, i.e. $D \circ R_{D}=\mathrm{Id}$.

Proof. Certainly $D\left(D^{-1}\left(f_{c}\right)=f_{c}\right.$, so $u=\left.D^{-1}\left(f_{c}\right)\right|_{M_{+}} \in \mathcal{C}^{-\infty}\left(M_{+} ; V_{1}\right)$ satifies $D u=f$ in the sense of extendible distributions. Since $f \in \mathcal{C}^{\infty}\left(M_{+} ; V_{2}\right)$ we can solve the problem $D u \equiv f$ in the sense of Taylor series at $H$, with the constant term freely prescibable. Using Borel's lemma, let $u^{\prime} \in \mathcal{C}^{\infty}\left(M_{+} ; V_{1}\right)$ have the appropriate Taylor series, with $b_{H} u^{\prime}=0$.. Then $D\left(u_{c}^{\prime}\right)-f_{c}=g \in \dot{\mathcal{C}}^{\infty}\left(M+; V_{2}\right)$. Thus $u^{\prime \prime}=D^{-1} g \in \mathcal{C}^{\infty}\left(M ; V_{1}\right)$. Since $D\left(u^{\prime}-u^{\prime \prime}\right)=f_{c}$, the uniqueness of solutions implies that $u=\left.\left(u^{\prime}-u^{\prime \prime}\right)\right|_{M_{+}} \in \mathcal{C}^{\infty}\left(M_{+} ; V_{1}\right)$.

Of course $R_{D}$ cannot be a two-sided inverse to $D$ since it has a large null space, described by $I_{D}$.

Problem 6.2. Show that, for $D$ as in Theorem $\frac{29.3 .1998 .187}{6.1} \frac{1}{} f \in C^{-\infty}\left(M_{+} ; V_{2}\right)$ and $v \in$ $\mathcal{C}^{\infty}\left(H ; V_{1}\right)$ there exists a unique $u \in \mathcal{C}^{\infty}\left(M_{+} ; V_{2}\right)$ such that $D u=f$ in $\mathcal{C}^{\infty}\left(M_{+} ; V_{2}\right)$ and $b_{H} u=\Pi_{C} v$.

### 6.10. Smoothing operators

The properties of smoothing operators on a compact manifold with boundary are essentially the same as in the boundaryless case. Rather than simply point to the earlier discussion we briefly repeat it here, but in an abstract setting.

Let $\mathcal{H}$ be a separable Hilbert space. In the present case this would be $L^{2}(X)$ or $L^{2}(X ; E)$ for some vector bundle over $X$, or some space $H^{m}(X ; E)$ of Sobolev sections. Let $\mathcal{B}=\mathcal{B}(\mathcal{H})$ be the algebra of bounded operators on $\mathcal{H}$ and $\mathcal{K}=\mathcal{K}(\mathcal{H})$ the ideal of compact operators. Where necessary the norm on $\mathcal{B}$ will be written $\|\quad\|_{\mathcal{B}} ; \mathcal{K}$ is a closed subspace of $\mathcal{B}$ which is the closure of the ideal $\mathcal{F}=\mathcal{F}(\mathcal{H})$ of finite rank bounded operators.

We will consider a subspace $\mathcal{J}=\mathcal{J}(\mathcal{H}) \subset \mathcal{B}$ with a stronger topology. Thus we suppose that $\mathcal{J}$ is a Fréchet algebra. That is, it is a Fréchet space with countably many norms $\|\quad\|_{k}$ such that for each $k$ there exists $k^{\prime}$ and $C_{k}$ with

$$
\begin{equation*}
\|A B\|_{k} \leq C_{k}\|A\|_{k^{\prime}}\|B\|_{k^{\prime}} \quad \forall A, B \in \mathcal{J} \tag{6.48}
\end{equation*}
$$

In particular of course we are supposing that $\mathcal{J}$ is a subalgebra (but not an ideal) in $\mathcal{B}$. To make it a topological ${ }^{*}$-subalgebra we suppose that

$$
\begin{equation*}
\|A\|_{\mathcal{B}} \leq C\|A\|_{k} \quad \forall A \in \mathcal{J}, *: \mathcal{J} \longrightarrow \mathcal{J} \tag{6.49}
\end{equation*}
$$

In fact we may suppose that $k=0$ by renumbering the norms. The third condition we impose on $\mathcal{J}$ implies that it is a subalgebra of $\mathcal{K}$, namely we insist that
$\mathcal{F} \cap \mathcal{J}$ is dense in $\mathcal{J}$,
in the Fréchet topology. Finally, we demand, in place of the ideal property, that $\mathcal{J}$ be a bi-ideal in $\mathcal{B}$ (also called a 'corner') that is,

$$
\begin{equation*}
A_{1}, A_{2} \in \mathcal{J}, B \in \mathcal{B} \Longrightarrow A_{1} B A_{2} \in \mathcal{J} \tag{6.51}
\end{equation*}
$$

$\forall k \exists k^{\prime}$ such that $\left\|A_{1} B A_{2}\right\|_{k} \leq C\left\|A_{1}\right\|_{k^{\prime}}\|B\|_{\mathcal{B}}\left\|A_{2}\right\|_{k^{\prime}}$.
 tions of a vector bundle over a compact manifold satisfies (6.48)-(5.52) with $\mathcal{H}=$ $H^{m}(X ; E)$ for any vector bundle $E$.

Proof. The smoothing operators on sections of a bundle $E$ can be written as integral operators

$$
\begin{equation*}
A u(x)=\int_{E} A(x, y) u(y), A(x, y) \in \mathcal{C}^{\infty}\left(X^{2} ; \operatorname{Hom}(E) \otimes \Omega_{R}\right) \tag{6.52}
\end{equation*}
$$

Thus $\mathcal{J}=\mathcal{C}^{\infty}\left(X^{2} ; \operatorname{Hom}(E) \otimes \Omega_{R}\right)$ and we make this identification topological. The norms are the $C^{k}$ norms. If $P_{1}, \ldots, p_{N^{(m)}}$ is a basis, on $\mathcal{C}^{\infty}\left(X^{2}\right)$, for the differential operators of order $m$ on $\operatorname{Hom}(E) \otimes \Omega_{L}$ then we may take

$$
\begin{equation*}
\|A\|_{m}=\sup _{j}\left\|P_{j} A\right\|_{L^{\infty}} \tag{6.53}
\end{equation*}
$$

for some inner products on the bundles. In fact $\operatorname{Hom}(E)=\pi_{L}^{*} E \otimes \pi_{R}^{*} E^{*}$ from it which follows easily that this is a basis $P_{j}=P_{j, k} \otimes P_{j, R}$ decomposing as products. From this (6.48) follows easily since

$$
\begin{equation*}
\|A B\|_{m}=\sup _{j}\left\|\left(P_{j L} A\right) \cdot\left(P_{j, R} B\right)\right\|_{\infty}\|A B\|_{L^{\infty}} \leq C\|A\|_{L^{\infty}}\|B\|_{L^{\infty}} \tag{6.54}
\end{equation*}
$$

by the compactnes of $X$. From this $\left(\frac{\mathrm{SO7}}{6.53)}\right.$ ) follows with $k=0$.
The density ( 5.50 ) is the density of the finite tensor product $\mathcal{C}^{\infty}(X ; E) \otimes \mathcal{C}^{\infty}$ $\left(X ; E^{*} \otimes \Omega_{L}\right)$ in $\mathcal{C}^{\infty}\left(X^{2} ; \operatorname{Hom}(E) \otimes \Omega_{L}\right)$. This follows from the boundaryless case by doubling (or directly). Similarly the bi-ideal condition ( $\overline{6.52}$ ) can be seen from the regularity of the kernel. A more satisfying argument using distribution theory follows from the next result.

Proposition 6.9. An operator $A: \dot{\mathcal{C}}^{\infty}(X ; E) \rightarrow \mathcal{C}^{-\infty}(X ; F)$ is a smoothing operator if and only if it extends by continuity to $\dot{\mathcal{C}}^{-\infty}(X ; E)$ and then has range in $\mathcal{C}^{\infty}(X ; F) \hookrightarrow \mathcal{C}^{-\infty}(X ; F)$.

Proof. If $A$ has the stated mapping property then compose with a Seeley extension operator, then $E A=\tilde{A}$ is a continuous linear map

$$
\tilde{A}: \dot{\mathcal{C}}^{-\infty}(X ; E) \rightarrow \mathcal{C}^{\infty}(\tilde{X} ; \tilde{F})
$$

for an extension of $F$ to $\tilde{F}$ over the double $\tilde{X}$. Localizing in the domain to trivialize $E$ and testing with a moving delta function we recover the kernel of $\tilde{A}$ as

$$
\tilde{A}(x, y)=\tilde{A} \cdot \delta_{y} \in \mathcal{C}^{\infty}(\tilde{X} ; \tilde{F})
$$

Thus it follows that $\tilde{A} \in \mathcal{C}^{\infty}\left(\tilde{X} \times X ; \operatorname{Hom}(E, \tilde{F}) \otimes \Omega_{R}\right)$. The converse is more obvious.

Returning to the general case of a bi-ideal as in ( $\left.{ }^{5} \mathrm{SO1} 48\right)-\left(\frac{\mathrm{SO}}{6.52} 2\right)$ we may consider the invertibility of $\operatorname{Id}+A, A \in \mathcal{J}$.

Proposition 6.10. If $A \in \mathcal{J}$, satisfying $\left(\frac{\sqrt{501}}{5.48)}-\left(\frac{\sqrt{504}}{5.5} 2\right)\right.$, then $\mathrm{Id}+A$ has a generalized inverse of the form $\operatorname{Id}+B, B \in \mathcal{J}$, with

$$
A B=\operatorname{Id}-\pi_{R}, B A=\operatorname{Id}-\pi_{L} \in \mathcal{J} \cap \mathcal{F}
$$

both finite rank self-adjoint projections.
Proof. Suppose first that $A \in \mathcal{J}$ and $\|A\|_{\mathcal{B}}<1$. Then $\operatorname{Id}+A$ is invertible in $\mathcal{B}$ with inverse $I d+B \in \mathcal{B}$,

$$
\begin{equation*}
B=\sum_{j \geq 1}(-1)^{j} A^{j} \tag{6.55}
\end{equation*}
$$

Not only does this Neumann series converge in $\mathcal{B}$ but also in $\mathcal{J}$ since for each $k$

$$
\begin{equation*}
\left\|A^{j}\right\|_{k} \leq C_{k}\|A\|_{k^{\prime}}\left\|A^{j-2}\right\|_{\mathcal{B}}\|A\|_{k^{\prime}} \leq C_{k}^{\prime}\|A\|_{\mathcal{B}}^{j-2}, j \geq 2 \tag{6.56}
\end{equation*}
$$

Thus $B \in \mathcal{J}$, since by assumption $\mathcal{J}$ is complete (being a Fréchet space). In this case $\operatorname{Id}+B \in \mathcal{B}$ is the unique two-sided inverse.

For general $A \in \mathcal{J}$ we use the assumed approximability in ( ${ }_{(503}^{5050)}$. Then $A=$ $A^{\prime}+A^{\prime \prime}$ when $A^{\prime} \in \mathcal{F} \cap \mathcal{J}$ and $\left\|A^{\prime \prime}\right\|_{\mathcal{B}} \leq C\left\|A^{\prime \prime}\right\|_{k}<1$ by appropriate choice. It follows that $\operatorname{Id}+B^{\prime \prime}=\left(\operatorname{Id}+A^{\prime \prime}\right)^{-1}$ is the inverse for $\operatorname{Id}+A^{\prime \prime}$ and hence a parameterix for $\mathrm{Id}+A$ :

$$
\begin{align*}
\left(\operatorname{Id}+B^{\prime \prime}\right)(\operatorname{Id}+A) & =\operatorname{Id}+A^{\prime}+B^{\prime \prime} A^{\prime}  \tag{6.57}\\
(\operatorname{Id}+A)\left(\operatorname{Id}+B^{\prime \prime}\right) & =\operatorname{Id}+A^{\prime}+A^{\prime} B^{\prime \prime}
\end{align*}
$$

Unfinished.
Lemma on subprojections.
with both 'error' terms in $\mathcal{F} \cap \mathcal{J}$.


### 6.11. Left and right parametrices

Suppose that $\mathcal{H}_{1}$ and $\mathcal{H}_{2}$ are Hilbert spaces and $A: \mathcal{H}_{1} \longrightarrow \mathcal{H}_{2}$ is a bounded linear operator between them. Let $\mathcal{J}_{1} \subset \mathcal{B}\left(\mathcal{H}_{1}\right)$ and $\mathcal{J}_{2} \subset \mathcal{B}\left(\mathcal{H}_{2}\right)$ be bi-ideals as in the previous section. A left parametrix for $A$, modulo $\mathcal{J}_{1}$, is a bounded linear map $B_{L}: \mathcal{H}_{2} \longrightarrow \mathcal{H}_{1}$ such that

$$
\begin{equation*}
B_{L} \circ A=\operatorname{Id}+J_{L}, \quad J_{L} \in \mathcal{J}_{1} \tag{6.58}
\end{equation*}
$$

Similarly a right parametrix for $A$, modulo $\mathcal{J}_{2}$ is a bounded linear map $B_{R}: \mathcal{H}_{2} \longrightarrow$ $\mathcal{H}_{1}$ such that

$$
\begin{equation*}
A \circ B_{R}=\mathrm{Id}+J_{R}, \quad J_{R} \in \mathcal{J}_{2} \tag{6.59}
\end{equation*}
$$

Proposition 6.11. If a bounded linear operator $A_{\dot{\mathrm{Son}} \mathrm{So}_{1}}^{\mathcal{H}_{1} \longrightarrow \mathcal{H}_{2} \text { has a left }}$ parametrix $B_{L}$ modulo a bi-ideal $\mathcal{J}_{1}$, satisfying ( $\overline{\boldsymbol{6} .48)-(\sqrt{6} .52), ~ t h e n ~} A$ has closed range, null space of finite dimension and there is a generalized left inverse, differing from the original left parametrix by a term in $\mathcal{J}_{1}$, such that

$$
\begin{equation*}
B_{L} \circ A=\operatorname{Id}-\pi_{L}, \pi_{L} \in \mathcal{J}_{1} \cap \mathcal{F} \tag{6.60}
\end{equation*}
$$

with $\pi_{L}$ the self-adjoint projection onto the null space of $A$.

Proof. Applying Proposition $\frac{\stackrel{5}{6} 010}{6.10}, \mathrm{Id}+J_{L}$ has a generalized inverse $\mathrm{Id}+J$, $J \in \mathcal{J}_{1}$, such that $(\operatorname{Id}+J)\left(\operatorname{Id}+J_{L}\right)=\left(\operatorname{Id}-\pi_{L}^{\prime}\right), \pi_{L}^{\prime} \in \mathcal{J}_{1} \cap \mathcal{F}$. Replacing $B_{L}$ by $\tilde{B}_{L}=(\operatorname{Id}+J) B_{L}$ gives a new left parametrix with error term $\pi_{L}^{\prime} \in \mathcal{J}_{1} \cap \mathcal{F}$. The null space of $A$ is contained in the null space of $B_{L}^{\prime} \circ A$ and hence in the range of $F_{L}$; thus it is finite dimensional. Furthermore the self-dajoint projection $\pi_{L}$ onto the null space is a subprojection of $\pi_{L}^{\prime}$, so is also an element of $\mathcal{J}_{1} \cap \mathcal{F}$. The range of $A$ is closed since it has finite codimension in $\operatorname{Ran}\left(A\left(\operatorname{Id}-\pi_{L}\right)\right)$ and if $f_{n} \in \operatorname{Ran}\left(A\left(\operatorname{Id}-\pi_{L}\right)\right)=A u_{n}, u_{n}=\left(\operatorname{Id}-\pi_{L}\right) u_{n}$, converges to $f \in \mathcal{H}_{2}$, then $u_{n}=B_{L} f_{n}$ converges to $u \in \mathcal{H}_{1}$ with $A\left(\operatorname{Id}-\pi_{L}\right) u=f$.

Proposition 6.12. If a bounded linear operator $A$ : $\mathcal{H}_{1} \longrightarrow \mathcal{H}_{2}$ has a right parametrix $B_{R}$ modulo a bi-ideal $\mathcal{J}_{2}$, satisfying ( $\mathbf{6 . 4 8 ) - ( \sqrt { 6 } . 5 2 ) \text { , then it has closed }}$ range of finite codimension and there is a generalized right inverse, differing from the original right parametrix by a term in $\mathcal{J}_{2}$, such that

$$
\begin{equation*}
A \circ B_{R}=\operatorname{Id}-\pi_{R}, \pi_{R} \in \mathcal{J}_{2} \cap \mathcal{F}, \tag{6.61}
\end{equation*}
$$

with $\operatorname{Id}-\pi_{R}$ the self-adjoint projection onto the range space of $A$.
Proof. The operator Id $+J_{R}$ has, by Proposition $\frac{\text { S010 }}{510}$, a generalized inverse $\operatorname{Id}+J$ with $J \in \mathcal{J}_{1}$. Thus $B_{R}^{\prime}=B_{R} \circ(\operatorname{Id}+J)$ is a right parametrix with error term Id $-\pi_{R}^{\prime}, \pi_{R}^{\prime} \in \mathcal{J}_{1} \cap \mathcal{F}$ being a self-adjoint projection. Thus the range of $A$ contains the range of $\mathrm{Id}-\pi_{R}^{\prime}$ and is therefore closed with a finite-dimensional complement. Furthemore the self-adjoint projection onto the range of $A$ is of the form $\operatorname{Id}-\pi_{R}$ where $\pi_{R}$ is a subprojection of $\pi_{R}^{\prime}$, so also in $\mathcal{J}_{1} \cap \mathcal{F}$.

The two cases, of an operator with a right or a left parametrix are sometimes combined in the term 'semi-Fredholm.' Thus an operator $A: \mathcal{H}_{1} \longrightarrow \mathcal{H}_{2}$ is semiFredholm if it has closed range and either the null space or the orthocomplement to the range is finite dimensional. The existence of a right or left parametrix, modulo the ideal of compact operators, is a necessary and sufficient condition for an operator to be semi-Fredholm.

### 6.12. Right inverse

In treating the 'general' case of an elliptic operator on compact manifold with boundar. ${ }_{8}$ we shall start by constructing an analogue of the right inverse in Lemma 5.6 . So now we assume that $D \in \operatorname{Diff}^{1}\left(X ; V_{1}, V_{2}\right)$ is an operator of Dirac type on a compact manifold with boundary.

To construct a right inverse for $D$ we follow the procedure in the boundaryless case. That is we use the construction of a pseudodifferential parametrix. In order to make this possible we need to extend $M$ and $D$ 'across the boundary.' This is certainly possible for $X$, since we may double it to a compact manifold without boundary, $2 X$. Then there is not obstruction to extending $D$ 'a little way' across the boundary. We shall denote by $M$ an open extension of $X$ (of the same dimension) so $X \subset M$ is a compact subset and by $\tilde{D}$ an extension of Dirac type to $M$.

The extension of $D$ to $\tilde{D}$, being elliptic, has a parametrix $\tilde{Q}$. Consider the map
28.4.1998.214

$$
\begin{equation*}
\tilde{Q}^{\prime}: L^{2}\left(X ; V_{2}\right) \longrightarrow H^{1}\left(X ; V_{1}\right), \tilde{Q}^{\prime} f=\tilde{Q} f_{c} \mid X \tag{6.62}
\end{equation*}
$$

where $f_{c}$ is the extension of $f$ to be zero outside $X$. Then $\tilde{Q}^{\prime}$ is a right parametrix, $D \tilde{Q}^{\prime}=\operatorname{Id}+E$ where $E$ is an operator on $L^{2}\left(X ; V_{2}\right)$ with smooth kernel on $X^{2}$.

Following Proposition $\frac{8.6 .1998 .246}{6.12, D}$ has a generalized right inverse $\tilde{Q}^{\prime \prime}=\tilde{Q}^{\prime}\left(\operatorname{Id}+E^{\prime}\right)$ up to finite rank smoothing and

$$
\begin{equation*}
D: H^{1}\left(X ; V_{1}\right) \longleftrightarrow L^{2}\left(X ; V_{2}\right) \tag{6.63}
\end{equation*}
$$

has closed range with a finite dimensional complement in $\mathcal{C}^{\infty}\left(X ; V_{2}\right)$.
Proposition 6.13. The map ( (28.63. $^{28.1998 .215}$ maps $^{\infty}\left(X ; V_{2}\right)$ to $\mathcal{C}^{\infty}\left(X ; V_{1}\right)$, it is surjective if and only if the only solution of $D^{*} u=0, u \in \dot{\mathcal{C}}^{\infty}\left(X ; V_{2}\right)$ is the trivial solution.

Proof. The regularity statement, that $Q^{\prime} \mathcal{C}^{\infty}(X ; V) \subset \mathcal{C}^{\infty}\left(X ; V_{1}\right)$ follows as in the proof of Lemma $\tilde{\sim}^{\text {6.6. Thus }} Q^{\prime}$ maps $\mathcal{C}^{\infty}\left(X ; V_{1}\right)$ to $\mathcal{C}^{\infty}\left(X ; V_{2}\right)$ if and only if any paramatrix $\tilde{Q}^{\prime}$ does so. Given $f \in \mathcal{C}^{\infty}\left(X ; V_{2}\right)$ we may solve $D u^{\prime} \equiv f$ in Taylor series at the boundary, with $u^{\prime} \in \mathcal{C}^{\infty}\left(X ; V_{1}\right)$ satisfying $b_{H} u^{\prime}=0$. Then $D\left(u^{\prime}\right)_{c}-f \in \dot{\mathcal{C}}^{\infty}\left(X ; V_{2}\right)$ so it follows that $\left.Q^{\prime}\left(f_{c}\right)\right|_{X} \in \mathcal{C}^{\infty}\left(X ; V_{1}\right)$.

Certainly ${ }_{28}$ any solution of $D^{*} u=0$ with $u \in \dot{\mathcal{C}}^{\infty}\left(X ; V_{2}\right)$ is orthogonal ${ }_{\text {28.4.1998.215 }}$ to the range of (6.63) so the condition is necessary. So, suppose that (6.63) is not surjective. Let $f \in L^{2}\left(X ; V_{2}\right)$ be in the orthocomplement to the range. Then Green's formula gives the pairing with any smooth section

$$
(D v, f)_{X}=\left(D \tilde{v}, f_{c}\right)_{\tilde{X}}=\left(\tilde{v}, D^{*} f_{c}\right)_{\tilde{X}}=0
$$

This means that $D^{*} f_{c}=0$ in $\tilde{X}$, that is as a supported distribution. Thus, $f \in$ $\dot{\mathcal{C}}^{\infty}\left(X ; V_{2}\right)$ satisfies $D^{*} f=0$.

As noted above we will proceed under the assumption that $D^{*} f$ has no such non-trivial solutions in $\dot{\mathcal{C}}^{\infty}\left(X ; V_{2}\right)$. This condition is discussed in the next section.

ThEOREM 6.2. If unique continuation holds for $D^{*}$ then $D$ has a right inverse

$$
\begin{equation*}
Q: \mathcal{C}^{\infty}\left(X: V_{2}\right) \longrightarrow \mathcal{C}^{\infty}\left(X ; V_{1}\right), D Q=\mathrm{Id} \tag{6.64}
\end{equation*}
$$

where $Q=\tilde{Q}^{\prime}+E, \tilde{Q}^{\prime} f=\tilde{Q} f \mid X$ where $\tilde{Q}$ is a parametrix for an extension of $D$ across the boundary and $E$ is a smoothing operator on $X$.

Proof. As just noted, unique continuation for $D^{*}$ implies that $D$ in ( $\frac{28.63 .1998 .215}{}$ is surjective. Since the parametrix maps $\mathcal{C}^{\infty}\left(X ; V_{2}\right)$ to $\mathcal{C}^{\infty}\left(X ; V_{1}\right), D$ must be surjective as a map from $\mathcal{C}^{\infty}\left(X ; V_{1}\right)$ to $\mathcal{C}^{\infty}\left(X ; V_{2}\right)$. The parametrix modulo finite rank operators can therefore be corrected to a right inverse for $D$ by the addition of a smoothing operator of finite rank.

### 6.13. Boundary map

The map $b$ from $\mathcal{C}^{\infty}(X ; E)$ to $\mathcal{C}^{\infty}(\partial X ; E)$ is well defined, and hence is well defined on the space of smooth solutions of $D$. We wish to show that it has closed range. To do so we shall extend the defintion to the space of square-integrable solutions. For any $s \in \mathbb{R}$ set

$$
\begin{equation*}
\mathcal{N}^{s}(D)=\left\{u \in H^{s}(X ; E) ; D u=0\right\} \tag{6.65}
\end{equation*}
$$

Of course the equation $D u=0$ is to hold in the sense of extendible distributions, which just means in the interior of $X$. Thus $\mathcal{N}^{\infty}(D)$ is the space of solutions of $D$ smooth up to the boundary.

$$
\begin{equation*}
\dot{D} u_{c}=v \cdot \delta(x), v \in H^{-\frac{1}{2}}(\partial X ; E) \tag{6.66}
\end{equation*}
$$

defines an injective bounded map $\tilde{b}: \mathcal{N}^{0}(D) \longrightarrow H^{-\frac{1}{2}}(\partial X ; E)$ by $\tilde{b}(u)=i \sigma(D)(d x) v$ which is an extension of $b: \mathcal{N}^{\infty}(D) \longrightarrow \mathcal{C}^{\infty}(\partial X ; E)$ defined by restriction to the boundary.

Proof. Certainly $\dot{D} u_{c} \subset \dot{\mathcal{C}}_{\partial X}^{\infty}(X ; E)$ has support in the boundary, so is a sum of products in any product decomposition of $X$ near $\partial X$,

$$
D\left(u_{c}\right)=\sum_{j} v_{j} \cdot \delta^{(j)}(x)
$$

Since $D$ is a first order operator and $u_{c} \in L^{2}(\tilde{X} ; E)$, for any local extension, $\dot{D} u_{c} \in \dot{H}^{-1}(X ; E)$. Localizing so that $E$ is trivial and the localized $v_{j}$ have compact supports this means that

$$
\begin{equation*}
\left(1+|\eta|^{2}+|\xi|^{2}\right)^{-\frac{1}{2}} \widehat{v_{j}}(\eta) \xi^{j} \in L^{2}\left(\mathbb{R}^{n}\right) \tag{6.67}
\end{equation*}
$$

If $v_{j} \neq 0$ for some. $j_{9}$. 0 this is not true even in some region $|\eta|<C$. Thus $v_{j} \equiv 0$ for $j>0$ and ( $\overline{6.66}$ ) must hold. Furthermore integration in $\xi$ gives

$$
\begin{gather*}
\int_{\mathbb{R}}\left(1+|\eta|^{2}+|\xi|^{2}\right)^{-1} d \xi=c\left(1+|\eta|^{2}\right)^{-\frac{1}{2}}, c>0, \quad \text { so }  \tag{6.68}\\
\int_{\mathbb{R}^{n-1}}\left(1+|\eta|^{2}\right)^{-\frac{1}{2}}|\hat{v}(\eta)|^{2} d \eta<0
\end{gather*}
$$

Thus $v \in H^{-\frac{1}{2}}(\partial X ; E)$ and $\tilde{b}$ is well defined. The jumps formula shows it to be an extension of $b$. The injectivity of $\tilde{b}$ follows from the assumed uniqueness of solutions to $\dot{D} u=0$ in $X$.

Notice that $\left.\frac{(6.6 .1998 .234}{6.68}\right)^{15}$ actually reversible. That is if $v \in H^{-\frac{1}{2}}(\partial X ; E)$ then $v \cdot \delta(x) \in H^{-1}(X ; E)$. This is the basis of the construction of a left parametrix for $\tilde{b}$, which then shows its range to be closed.
6.6 .1998 .239

Lemma 6.8. The boundary map $\tilde{b}$ in Lemma $\frac{\text { 6.6.1998.231 }}{6.7 \text { has a continuous left paramet- }}$ rix $\widetilde{I_{D}}: H^{-\frac{1}{2}}(\partial X ; E) \longrightarrow \mathcal{N}^{0}(D), I_{D} \circ \tilde{b}=\mathrm{Id}+G$, where $G$ has smooth kernel on $X \times \partial X$, and the range of $\tilde{b}$ is therefore a closed subspace of $H^{-\frac{1}{2}}(\partial X ; E)$.

Proof. The parametrix $\widetilde{I_{D}}$ is given directly by the parametrix $\tilde{Q}$ for $\tilde{D}$, and extension to $\tilde{X}$. Applying $\tilde{Q}$ to $\left(\frac{\sigma .6 \bar{\delta}) \text { gives }}{}\right.$
with $R$ having smooth kernel. Since $\widetilde{I_{D}}$ is bounded from $H^{-\frac{1}{2}}(\partial X ; E)$ to $L^{2}(X ; E)$ and $R$ is smoothing it follows from Proposition $\frac{8.6 .1998 .245}{6.11}$ that the range of $\tilde{b}$ is closed.

### 6.14. Calderòn projector

Having shown that the range of $\tilde{b}$ in Lemma $\frac{6.6 .1998 .231}{6.7 \text { is closed in } H^{-\frac{1}{2}}(\partial X ; E) \text { we }{ }^{6}(\partial)}$ now deduce that there is a pseudodifferential projection onto it. The discussion above of the boundary values of the $\tilde{Q}(w \cdot \delta(x))$ is local, and so applies just as well
in the presentmore general case. Since this is just the definition of the map $\widetilde{I_{D}}$ in Lemma 5.8 , we conclude directly that
6.6 .1998 .240

$$
\begin{equation*}
P v=\lim _{X^{\circ}} \widetilde{I_{D}} v, v \in \mathcal{C}^{\infty}(\partial X ; E) \tag{6.70}
\end{equation*}
$$

defines $P \in \Psi^{0}(\partial X ; E)$.
 exist $A, B \in \Psi^{-\infty}(\partial X ; E)$ such that $P-\mathrm{Id}=A$ on $\operatorname{Ran}(\tilde{b})$ and $\operatorname{Ran}(P+B) \subset$ $\operatorname{Ran}(\tilde{b})$.
Proof needs clarification.

Proof. That $P^{2}-P \in \Psi^{-\infty}(\partial X ; E)$ follows, as above, from the fact that $\tilde{Q}$ is a two-sided parametrix on distributions supported in $X$. Similarly we may use the right inverse of $D$ to construct $B$. If $v \in H^{-\frac{1}{2}}(\partial X ; E)$ then by construction,

$$
D \widetilde{I_{D}} v=R^{\prime} v
$$

where $R^{\prime}$ has a smooth kernel on $X \times \partial X$. Applying the right inverse $Q$ it follows that $u^{\prime}=\widetilde{I_{D}} v-\left(Q \circ R^{\prime}\right) v \in \mathcal{N}^{0}(D)$, where $Q \circ R^{\prime}$ also has smooth kernel on $X \times \partial X$. Thus $\tilde{b}\left(u^{\prime}\right)=(P+B) v \in \operatorname{Ran}(\tilde{b})$ where $B$ has kernel arising from the restriction of the kernel of $A \circ R^{\prime}$ to $\partial X \times \partial X$, so $B \in \Psi^{-\infty}(\partial X ; E)$.

Now we may apply Proposition $\frac{4.6 .1998 .227}{5.6 \text { with } F}=\operatorname{Ran}(\tilde{b})$ and $s=-\frac{1}{2}$ to show the existence of a Calderòn projector.

Proposition 6.14. If $D$ is a generalized Dirac operator on $X$ then there is an element $\Pi_{C} \in \Psi^{0}(\partial X ; E)$ such that $\Pi_{C}^{2}=\Pi_{C}, \operatorname{Ran}\left(\Pi_{6} .999 \overline{\overline{8}} . \operatorname{Ran}(\tilde{b})\right.$ on $H^{-\frac{1}{2}}(\partial X ; E)$, $\Pi_{C}-P \in \Psi^{-\infty}(\partial X ; E)$ where $P$ is defined by $(\overline{6.7 U})$ and $\operatorname{Ran}\left(\Pi_{C}\right)=\operatorname{Ran}(b)$ on $\mathcal{C}^{\infty}(\partial X ; E)$.

Proof. The existence of psuedodifferential projection, $\Pi_{C}$, differing from $P$ by a smoothing operator and with range $\operatorname{Ran}(\tilde{b})$ is a direct consequence of the application of Proposition $\frac{4.6 .15}{5.6}$. It follows that $\operatorname{Ran}(\tilde{b}) \cap \mathcal{C}^{\infty}(\partial X ; E)$ is dense in $\operatorname{Ran}(\tilde{b})$ in the topology of $H^{-\frac{1}{2}}(\partial X ; E)$. Furthermore, if follows that if $v \in \operatorname{Ran}(\tilde{b}) \cap$ $\mathcal{C}^{\infty}(\partial X ; E)$ then $u \in \mathcal{N}^{0}(D)$ such that $\tilde{b} u=v$ is actually in $\mathcal{C}^{\infty}(X ; E)$, i.e. it is in $\mathcal{N}^{\infty}(D)$. Thus the range of $b$ is just $\operatorname{Ran}(\tilde{b}) \cap \mathcal{C}^{\infty}(\partial X ; E)$ so $\operatorname{Ran}(b)$ is the range of $\Pi_{C}$ acting on $\mathcal{C}^{\infty}(\partial X ; E)$.

In particular $\tilde{b}$ is just the continuous extension of $b$ from $\mathcal{N}^{\infty}(D)$ to $\mathcal{N}^{0}(D)$, of which it is a dense subset. Thus we no longer distinguish between these two maps and set $\tilde{b}=b$.

### 6.15. Poisson operator

### 6.16. Unique continuation

### 6.17. Boundary regularity

### 6.18. Pseudodifferential boundary conditions

The discussion above shows that for any operator of Dirac type the 'Calderòn realization' of $D$,

$$
\begin{equation*}
D_{\mathcal{C}}:\left\{u \in H^{s}\left(X ; E_{1}\right) ; \Pi_{\mathcal{C}} b u=0\right\} \longrightarrow H^{s-1}\left(X ; E_{2}\right), s>\frac{1}{2} \tag{6.71}
\end{equation*}
$$

is an isomorphism.
We may replace the Calderòn projector in $\left(\frac{28.4 .1998 .1}{5.1)}\right.$ by a more general projection $\Pi$, acting on $\mathcal{C}^{\infty}\left(\partial X, V_{1}\right)$, and consider the map

$$
\begin{equation*}
D_{\Pi}:\left\{u \in \mathcal{C}^{\infty}\left(X ; V_{1}\right) ; \Pi b u=0\right\} \longrightarrow \mathcal{C}^{\infty}\left(X ; V_{2}\right) \tag{6.72}
\end{equation*}
$$

In general this map will not be particularly well-behaved. We will be interested in the case that $\Pi \in \Psi^{0}\left(\partial X ; V_{1}\right)$ is a pseudodifferential projection. Then a condition for the map $D_{\Pi}$ to be Fredholm can be given purely in terms of the relationship between $\Pi$ and the (any) Calderòn projector $\Pi_{\mathcal{C}}$.
29.3.1998.188
29.3.1998.189
29.3 .1998 .190

Theorem 6.3. If $D \in \operatorname{Diff}^{1}\left(X ; E_{1}, E_{2}\right)$ is of Dirac type and $P i \in \Psi^{0}\left(\partial X ; E_{1}\right)$ is a projection then the map

$$
\begin{equation*}
D_{\Pi}:\left\{u \in \mathcal{C}^{\infty}\left(X ; E_{1}\right) ; \Pi\left(u_{\partial X}\right)=0\right\} \xrightarrow{D} \mathcal{C}^{\infty}\left(X ; E_{2}\right) \tag{6.73}
\end{equation*}
$$

is Fredholm if and only if
(6.74) $\quad \Pi \circ \Pi_{C}: \operatorname{Ran}\left(\Pi_{C}\right) \cap \mathcal{C}^{\infty}\left(\partial V_{1}\right) \longrightarrow \operatorname{Ran}(\Pi) \cap \mathcal{C}^{\infty}\left(\partial E_{1}\right)$ is Fredholm
and then $t_{29} .3 .1998 .1 f_{0} D_{\Pi}$ is equal to the relative index of $\Pi_{\mathcal{C}}$ and $\Pi$, that is the index of ( 6.74 ).

Below we give a symbolic condition equivalent which implies the Fredholm condition. 19 If en enough regularity conditions are imposed on the generalized inverse to ( 6.71 ) then this symbolic is also necessary.

Proof. The null space of $D_{\Pi}$ is easily analysed. Indeed $D u=0$ implies that $u \in \mathcal{N}$, so the null space is isomorphic to its image under the boundary map:

$$
\{u \in \mathcal{N} ; \Pi b u=0\} \xrightarrow{b}\{v \in \mathcal{C} ; \Pi v=0\} .
$$

Since $\mathcal{C}$ is the range of $\Pi_{\mathcal{C}}$ this gives the isomorphism

$$
\begin{equation*}
\operatorname{Nul}\left(D_{\Pi}\right) \simeq \operatorname{Nul}\left(\Pi \circ \Pi_{\mathcal{C}}: \mathcal{C} \longrightarrow \operatorname{Ran}(\Pi)\right) \tag{6.75}
\end{equation*}
$$

In particular, the null space is finite dimensional if and only if the null space of $\Pi \circ \Pi_{\mathcal{C}}$ is finite dimensional.

Similarly, consider the range of $D_{\Pi}$. We construct a map

$$
\begin{equation*}
\tau: \mathcal{C}^{\infty}\left(\partial X ; V_{1}\right) \longrightarrow \mathcal{C}^{\infty}\left(X ; V_{2}\right) / \operatorname{Ran}\left(D_{\Pi}\right) \tag{6.76}
\end{equation*}
$$

Indeed each $v \in \mathcal{C}^{\infty}\left(\partial X ; V_{1}\right)$ is the boundary value of some $u \in \mathcal{C}^{\infty}\left(X: V_{1}\right)$, let $\tau(v)$ be he class of $D U$. This is well-defined since any other extension $u^{\prime}$ is such that $b\left(u-u^{\prime}\right)=0$, so $D\left(u-u^{\prime}\right) \in \operatorname{Ran}\left(D_{\Pi}\right)$. Furthermore, $\tau$ is surjective, since $D_{\mathcal{C}}$ is surjective. Consider the null space of $\tau$. This certainly contains the null space of $\Pi$. Thus consider the quotient map

$$
\tilde{\tau}: \operatorname{Ran}(\Pi) \longrightarrow \mathcal{C}^{\infty}\left(X: V_{2}\right) / \operatorname{Ran}\left(D_{\Pi}\right)
$$

which is still surjective. Then $\tilde{\tau}(v)=0$ if and only if there exists $v^{\prime} \in \mathcal{C}$ such that $\Pi\left(v-v^{\prime}\right)=0$. That is, $\tilde{\tau}(v)=0$ if and only if $\Pi(v)=\Pi \circ \Pi_{\mathcal{C}}$. This shows that the finer quotient map
28.4.1998.213

$$
\begin{equation*}
\tau^{\prime}: \operatorname{Ran}(\Pi) / \operatorname{Ran}\left(\Pi \circ \Pi_{\mathcal{C}}\right) \longleftrightarrow \mathcal{C}^{\infty}\left(X ; V_{2}\right) / \operatorname{Ran}\left(D_{\Pi}\right) \tag{6.77}
\end{equation*}
$$

is an isomorphism. This shows that the range is closed and of finite codimension if $\Pi \circ \Pi_{\mathcal{C}}$ is Fredholm.

The converse follows by reversing these arguments.

### 6.19. Gluing

Returning to the case of a compact manifold without boundary, $M$, with a dividing hypersurface $H$ we can now give a gluing result for the index.
29.3.1998.191 THEOREM 6.4. If $D \in \operatorname{Diff}^{1}\left(M ; E_{1}, E_{2}\right)$ is of Dirac type and $M=M_{1} \cap M_{2}$ is the union of two manifolds with boundary intersecting in their common boundary $\partial M_{1} \cap \partial M_{2}=H$ then
(6.78) $\quad \operatorname{Ind}(D)=\operatorname{Ind}\left(\Pi_{1, \mathcal{C}}, \operatorname{Id}-\Pi_{2, \mathcal{C}}\right)=\operatorname{Ind}\left(\Pi_{2, \mathcal{C}}, \operatorname{Id}-\Pi_{1, \mathcal{C}}\right)$
where $\Pi_{i, \mathcal{C}}, i=1,2$, are the Calderòn projections for $D$ acting over $M_{i}$.
S.Local.boundary
S.Absolute.relative
S.Transmission.condition

### 6.20. Local boundary conditions

### 6.21. Absolute and relative Hodge cohomology

### 6.22. Transmission condition

CHAPTER 7

Scattering calculus

## CHAPTER 8

## The wave kernel

Let us return to the subject of "good distributions" as exemplified by Dirac delta 'functions' and the Schwartz kernels of pseudodifferential operators. In fact we shall associate a space of "conormal distributions" with any submanifold of a manifold.

Thus let $X$ be a $\mathcal{C}^{\infty}$ manifold and $Y \subset X$ a closed embedded submanifold we can easily drop the assumption that $Y$ is closed and even replace embedded by immersed, but let's treat the simplest case first! To say that $Y$ is embedded means that each $\bar{y} \in Y$ has a coordinate neighbourhood $U$, in $X$, with coordinate $x_{1}, \ldots, x_{n}$ in terms of which $\bar{y}=0$ and

$$
\begin{equation*}
Y \cap U=\left\{x,=\cdots=x_{k}=0\right\} . \tag{8.1}
\end{equation*}
$$

We want to define

$$
\begin{equation*}
I^{*}\left(X, Y ; \Omega^{\frac{1}{2}}\right) \subset \mathcal{C}^{-\infty}\left(X ; \Omega^{\frac{1}{2}}\right) \tag{8.2}
\end{equation*}
$$

to consist of distributions which are singular only at $Y$ and small "along $Y$."
So if $u \in \mathcal{C}_{c}^{-\infty}(U)$ then in local coordinates (지.i) we can identify $u$ with $u^{\prime} \in$ $\mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right)$ so $u^{\prime} \in H_{c}^{s}\left(\mathbb{R}^{n}\right)$ for some $s \in \mathbb{R}$. To say that $u$ is 'smooth along $Y$ ' means we want to have

$$
\begin{equation*}
D_{x_{k+1}}^{l_{1}} \ldots D_{x_{n}}^{l_{n-k}} u^{\prime} \in H_{c}^{s^{\prime}}\left(\mathbb{R}^{n}\right) \quad \forall l_{1}, \ldots, l_{n-k} \tag{8.3}
\end{equation*}
$$

and a fixed $s^{\prime}$, independent of $l$ (but just possibly ${ }_{116.3}$ different from the initial $s$ ); of course we can take $s=s^{\prime}$. Now conditions like ( $(\mathbf{8 . 3}$ ) do not limit the singular support of $u^{\prime}$ at all! However we can add a requirement that multiplication by a function which vanishes on $Y$ makes $u^{\prime}$ smooth, by one degree, i.e.

$$
\begin{equation*}
x_{1}^{p_{1}} \ldots x_{k}^{p_{k}} u^{\prime} \in H^{s+|p|}\left(\mathbb{R}^{n}\right),|p|=p_{1}+\cdots+p_{k} . \tag{8.4}
\end{equation*}
$$

This last condition implies

$$
\begin{equation*}
D_{1}^{q_{1}} \ldots D_{k}^{q_{k}} x_{1}^{p_{1}} \ldots x_{k}^{p_{k}} u^{\prime} \in H^{s}\left(\mathbb{R}^{n}\right) \text { if }|q| \leq|p| . \tag{8.5}
\end{equation*}
$$

Consider what happens if we rearrange the order of differentiation and multiplication in (18.5). Since we demand (18.5) for all $p, q$ with $|q| \leq|p|$ we can show in tial that

$$
\begin{equation*}
\Longrightarrow \tag{8.7}
\end{equation*}
$$

$$
\begin{equation*}
\forall|q| \leq|p| \leq L \tag{8.6}
\end{equation*}
$$

$$
\begin{equation*}
\prod_{i=1}^{L}\left(x_{j_{i}} D_{\ell_{i}}\right) u \in H^{s}\left(\mathbb{R}^{n}\right) \quad \forall \text { pairs, }\left(j_{i, \ell_{i}}\right) \in(1, \ldots, k)^{2} \tag{8.8}
\end{equation*}
$$

Of course we can combine $\left(\frac{16.3}{8.3}\right)$ and $\left(\frac{16.6}{8.8)}\right.$ and demand

$$
\begin{gather*}
\prod_{i=1}^{L_{2}} D_{p_{i}} \prod_{i=1}^{L_{1}}\left(x_{j_{i}} D_{\ell_{i}}\right) u^{\prime} \in H_{c}^{s}\left(\mathbb{R}^{n}\right)\left(j_{j}, \ell_{i}\right) \in(1, \ldots, k)^{2}  \tag{8.9}\\
\forall L_{1}, L_{2} p_{i} \in(k+1, \ldots u)
\end{gather*}
$$

Problem 8.1. Show that $\left(\frac{16.7}{8.9}\right)$ implies $\left(\frac{16.3}{8.3}\right)$ and $\left(\frac{16.4}{8.4}\right)$
The point about $\left(\frac{16.7}{8.9)}\right.$ is that it is easy to interpret in a coordinate independent way. Notice that putting $\mathcal{C}^{\infty}$ coefficients in front of all the terms makes no difference.
16.8 LEMMA 8.1. The space of all $\mathcal{C}^{\infty}$ vector fields on $\mathbb{R}^{n}$ tangent to the submanifold $\left\{x_{1}=\cdots=x_{k}=0\right\}$ is spanning over $\mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right)$ by

$$
\begin{equation*}
x_{i} D_{j}, D_{p} i, j \leq k, p>k \tag{8.10}
\end{equation*}
$$

Proof. A $\mathcal{C}^{\infty}$ vector field is just a sum

$$
\begin{equation*}
V=\sum_{j \leq k} a_{j} D_{j}+\sum_{p>k} b_{p} D_{p} \tag{8.11}
\end{equation*}
$$

Notice that the $D_{p}$, for $p>k$, are tangent to $\left\{x_{1}=\cdots=x_{k}=0\right\}$, so we can assume $b_{p}=0$. Tangency is then given by the condition

$$
\begin{equation*}
V x) i=0 \text { and }\left\{x_{1}=\cdots=x_{k}=0\right\}, i=1, \ldots, h \tag{8.12}
\end{equation*}
$$

i.e. $a_{j}=\sum_{\ell=1} a_{j \ell} x_{\ell}, 1 \leq j \leq h$. Thus

$$
\begin{equation*}
V=\sum_{\ell=1} a_{j \ell} x_{\ell} D_{j} \tag{8.13}
\end{equation*}
$$

which proves $\left(\frac{16.9}{8.10}\right)$.
This allows us to write $\left(\frac{16.7}{8.9)}\right.$ in the compact form

$$
\begin{equation*}
\mathcal{V}\left(\mathbb{R}^{n}, Y_{k}\right)^{p} u^{\prime} \subset H_{c}^{s}\left(\mathbb{R}^{n}\right) \forall p \tag{8.14}
\end{equation*}
$$

where $\mathcal{V}\left(\mathbb{R}^{n}, Y_{k}\right)$ is just the space of all $\mathcal{C}^{\infty}$ vector fields tangent to $Y_{k}=\left\{x_{1}=\right.$ $\left.\cdots=x_{k}=0\right\}$. Of course the local coordinate just reduce vector fields tangent to $Y$ to vector fields tangent to $Y_{k}$ so the invariant version of (8.14) is

$$
\begin{equation*}
\mathcal{V}(X, Y)^{p} u \subset H^{s}\left(X ; \Omega^{\frac{1}{2}}\right) \forall p \tag{8.15}
\end{equation*}
$$

To interpret $\left(\frac{16.11}{8.15}\right)$ we only need recall the (Lie) action of vector fields on halfdensities. First for densities: The formal transpose of $V$ is $-V$, so set

$$
\begin{equation*}
{ }^{L} V \phi(\psi)=\phi(-V \psi) \tag{8.16}
\end{equation*}
$$

if $\phi \in \mathcal{C}^{\infty}(X ; \Omega), \psi \in \mathcal{C}^{\infty}(X)$. On $\mathbb{R}^{n}$ then becomes

$$
\begin{aligned}
\int{ }^{L} V \phi \cdot & \psi=-\int \phi \cdot V \psi \\
= & -\int \phi(x) V \psi \cdot d x \\
= & \int\left(V \phi(x)+\delta_{V} \phi\right) \psi d x \\
\delta_{V} & =\sum_{i=1}^{n} D_{i} a_{i} \quad \text { if } V=\Sigma a_{i} D_{i}
\end{aligned}
$$

i.e.

$$
\begin{equation*}
L_{V}(\phi|d x|)=(V \phi)|d x|+\delta_{V} \phi \tag{8.18}
\end{equation*}
$$

Given the tensorial properties of density, set

$$
\begin{equation*}
L_{V}\left(\phi|d x|^{t}\right)=V \phi|d x|^{t}+t \delta_{V} \phi \tag{8.19}
\end{equation*}
$$

This corresponds to the natural trivialization in local coordinates.

$$
\begin{align*}
& I H^{s}\left(X, Y ; \Omega^{\frac{1}{2}}\right)=\left\{u \in H^{s}\left(X ; \Omega^{\frac{1}{2}}\right) \text { satisfying (11) }\right\} \\
& I^{*}\left(X, Y ; \Omega^{\frac{1}{2}}\right)=\bigcup_{s} I H^{s}\left(X, Y ; \Omega^{\frac{1}{2}}\right) \tag{8.20}
\end{align*}
$$

Clearly

$$
\begin{equation*}
u \in I^{*}\left(X, Y ; \Omega^{\frac{1}{2}}\right) \Longrightarrow u \upharpoonright X \backslash Y \in \mathcal{C}^{\infty}\left(X \backslash Y ; \Omega^{\frac{1}{2}}\right) \tag{8.21}
\end{equation*}
$$

and

$$
\begin{equation*}
\bigcap_{s} I H^{s}\left(X, Y ; \Omega^{\frac{1}{2}}\right)=\mathcal{C}^{\infty}\left(X ; \Omega^{\frac{1}{2}}\right) \tag{8.22}
\end{equation*}
$$

Let us try to understand these distributions in some detail! To do so we start with a very simple case, namely $Y=\{p\}$ is a point; so we only have one coordinate system. So construct $p=0 \in \mathbb{R}^{n}$.

$$
\begin{gather*}
u \in I_{c}^{*}\left(\mathbb{R}^{n},\{0\} ; \Omega^{\frac{1}{2}}\right) \Longrightarrow u=u^{\prime}|d x|^{\frac{1}{2}} \text { when }  \tag{8.23}\\
x^{\alpha} D_{x}^{\beta} u^{\prime} \in H_{c}^{s}\left(\mathbb{R}^{n}\right), \quad s \text { fixed } \forall|\alpha| \geq|\beta|
\end{gather*}
$$

Again by a simple commutative argument this is equivalent to

$$
\begin{equation*}
D_{x}^{\beta} x^{\alpha} u^{\prime} \in H_{c}^{s}\left(\mathbb{R}^{n}\right) \quad \forall|\alpha| \geq|\beta| \tag{8.24}
\end{equation*}
$$

We can take the Fourier transform of $\left(\frac{16.20}{8.24)}\right.$ and get

$$
\begin{equation*}
\xi^{\beta} D_{\xi}^{\alpha} \hat{u}^{\prime} \in\langle\xi\rangle^{-s} L^{2}\left(\mathbb{R}^{n}\right) \forall|\alpha| \geq|\beta| \tag{8.25}
\end{equation*}
$$

In this form we can just replace $\xi^{\beta}$ by $\langle\xi\rangle^{|\beta|}$, i.e. $\left(\frac{16.21}{8.25)}\right.$ just says

$$
\begin{equation*}
D_{\xi}^{\alpha} \hat{u}^{\prime}(\xi) \in\langle\xi\rangle^{-s-|\beta|} L^{2}\left(\mathbb{R}^{n}\right) \forall \alpha \tag{8.26}
\end{equation*}
$$

Notice that this is very similar to a symbol estimate, which would say

$$
\begin{equation*}
D_{\xi}^{\alpha} \hat{u}^{\prime}(\xi) \in\langle\xi\rangle^{m-|\alpha|} L^{\infty}\left(\mathbb{R}^{n}\right) \forall \alpha \tag{8.27}
\end{equation*}
$$

16. LEMMA 8.2 The The estimate $\left(\frac{16.22}{8.26)}\right.$ implies $\left(\frac{16.23}{8.27}\right)$ for any $m>-s-\frac{n}{2}$; conversely (16.27) implies (18.26) for any $s<-m-\frac{n}{2}$.

Proof. Let's start with the simple derivative, $\left(\frac{16.23}{8.27}\right)$ implies $\left(\frac{16.22}{8.26}\right)$. This really reduces to the case $\alpha=0$. Thus

$$
\begin{equation*}
\langle\xi\rangle^{M} L^{\infty}\left(\mathbb{R}^{n}\right) \subset L^{2}\left(\mathbb{R}^{n}\right) \Longrightarrow M<-\frac{n}{2} \tag{8.28}
\end{equation*}
$$

is the inequality

$$
\begin{equation*}
\left(\int|u|^{2} d \xi\right)^{\frac{1}{2}} \leq \sup \langle\xi\rangle^{-M}|u|\left(\int\langle\xi\rangle^{2 M} d \xi\right)^{\frac{1}{2}} \tag{8.29}
\end{equation*}
$$

and

$$
\begin{equation*}
\int\langle\xi\rangle^{2 M} d \xi=\int\left(1+|\xi|^{2}\right)^{M} d \xi<\infty \text { iff } M<-\frac{n}{2} \tag{8.30}
\end{equation*}
$$

To get $\left(\frac{16.23}{8.27}\right)$ we just show that $\left(\frac{16.23}{8.27}\right)$ implies

$$
\begin{equation*}
\langle\xi\rangle^{s+|\alpha|} D_{\xi}^{\alpha} \hat{u}^{\prime} \in\langle\xi\rangle^{m+s} L^{\infty} \subset L^{2} \text { if } m+s<-\frac{n}{2} \tag{8.31}
\end{equation*}
$$

The converse is a little trickier. To really see what is going $\rho$ n we can reduce $\left(\frac{16.22}{8.26)}\right.$ to a one dimensional version. Of course, near $\xi=0$, ( $\overline{8.26}$ ) just says $\hat{u}^{\prime}$ is $\mathcal{C}^{\infty}$, so we can assume that $|\xi|>1$ on $\operatorname{supp} \hat{u}^{\prime}$ and introduce polar coordinates:

$$
\begin{equation*}
\xi=t w, w \in S^{n-1} t>1 \tag{8.32}
\end{equation*}
$$

Then
Exercise 2. Show that $\left(\frac{16.22}{8.26}\right)$ (or maybe better, $\left(\frac{16.21}{8.25}\right)$ implies that

$$
\begin{equation*}
D_{t}^{k} P \hat{u}^{\prime}(t w) \in t^{-s-k} L^{2}\left(\mathbb{R}^{+} \times S^{n-1} ; t^{n-1} d t d w\right) \forall k \tag{8.33}
\end{equation*}
$$

for any $\mathcal{C}^{\infty}$ differential operator on $S^{n-1}$.
In particular we can take $P$ to be elliptic of any order, so ( $\left.{ }^{[16.27}\right)$ actually implies

$$
\begin{equation*}
\sup _{w} D_{t}^{k} P \hat{u}(t, w) \in t^{-s-k} L^{2}\left(\mathbb{R}^{+} ; t^{n-1} d t\right) \tag{8.34}
\end{equation*}
$$

or, changing the meaning to $d t$,

$$
\begin{equation*}
\sup _{w \in S^{n-1}}\left|D_{t}^{k} P \hat{u}(t, w)\right| \in t^{-s-k-\frac{n-1}{2}} L^{2}\left(\mathbb{R}^{+}, d t\right) \tag{8.35}
\end{equation*}
$$

So we are $\operatorname{ing}_{29}$ the one dimensional case, with $s$ replaced by $s+\frac{n-1}{2}$. Now we can rewrite ( 8.35 ) as

$$
\begin{equation*}
D_{t} t^{q} D_{t}^{k} P \hat{u} \in t^{r} L^{2}, \forall k, r-q=-s-k-\frac{n-1}{2}-1 . \tag{8.36}
\end{equation*}
$$

Now, observe the simple case:

$$
\begin{equation*}
f=0 t<1, D_{t} f \in t^{r} L^{2} \Longrightarrow f \in L^{\infty} \text { if } r<-\frac{1}{2} \tag{8.37}
\end{equation*}
$$

since

$$
\begin{equation*}
\sup |f|=\int_{-\infty}^{t} t^{r} g \leq\left(\int|g|^{2}\right)^{\frac{1}{2}} \cdot\left(\int_{-\infty}^{t} t^{2 r}\right)^{\frac{1}{2}} \tag{8.38}
\end{equation*}
$$

Thus from $\left(\frac{16.30}{8.36)}\right.$ we deduce $\leq\left(\int|g|^{2}\right)^{\frac{1}{2}}$

$$
\begin{equation*}
D_{t}^{k} P \hat{u} \in t^{-q} L^{\infty} \text { if } r<-\frac{1}{2} \text {, i.e. }-q>-s-k-\frac{n}{2} \tag{8.39}
\end{equation*}
$$

Finally this gives $\left(\frac{16.23}{8.27}\right)$ when we go back from polar coordinates, to prove the lemma.

Definition 8.2. Set, for $m \in \mathbb{R}$,

$$
\begin{equation*}
I_{c}^{m}\left(\mathbb{R}^{n}, \mid[0\}\right)=\left\{u \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right) ; \hat{u} \in S^{m-\frac{n}{4}}\left(\mathbb{R}^{n}\right)\right\} \tag{8.40}
\end{equation*}
$$

with this definition,
16.33

$$
\begin{equation*}
I H^{s}\left(\mathbb{R}^{n},\{0\}\right) \subset I_{c}^{m}\left(\mathbb{R}^{n},\{0\}\right) \subset I_{c}^{s^{\prime}}\left(\mathbb{R}^{n},\{0\}\right) \tag{8.41}
\end{equation*}
$$

provided

$$
\begin{equation*}
s>-m-\frac{n}{4}>s^{\prime} \tag{8.42}
\end{equation*}
$$

Exercise 3. Using Lemma 24, prove ( $\frac{16.33}{8.41}$ ) carefully.
So now what we want to do is to define $I_{c}^{m}\left(X,\{p\} ; \Omega^{\frac{1}{2}}\right)$ for any $p \in X$ by

$$
\begin{gather*}
u \in I_{c}^{m}\left(X,\{p\} ; \Omega^{\frac{1}{2}}\right) \Longleftrightarrow F^{*}(\phi u) \in I_{c}^{m}\left(\mathbb{R}^{n},\{0\}\right),  \tag{8.43}\\
u \upharpoonright X \backslash\{p\} \in \mathcal{C}^{\infty}(X \backslash\{p\}) .
\end{gather*}
$$

Here we have a little problem, namely we have to check that $I^{m}\left(\mathbb{R}^{n} n{ }_{33}\{0\}\right)$ is invariant under coordinate changes. Fortunately we can do this using (18.41).
17.7
17.8

Proof. Start with a simple case, that $F$ is linear. Then

$$
\begin{equation*}
u=(2 \pi)^{-n} \int e^{i x \xi} a(\xi) d \xi, a \in S^{m-\frac{n}{4}}\left(\mathbb{R}^{n}\right) \tag{8.45}
\end{equation*}
$$

so

$$
\begin{align*}
F^{*} u & =(2 \pi)^{-n} \int e^{i A x \cdot \xi} a(\xi) d \xi F x=A x \\
& =(2 \pi)^{-n} \int i^{i x \cdot A^{t} \xi} a(\xi) d \xi \\
& =(2 \pi)^{-n} \int e^{i x \cdot \eta} a\left(\left(A^{t}\right)^{-1} \eta\right)|\operatorname{det} A|^{-1} d \eta
\end{align*}
$$

Since $\left.a\left(\left(A^{t}\right)^{-1} \eta\right)|\operatorname{det} A|^{-1} \in S^{m-\frac{n}{4}} \mathbb{R}^{n}\right)$ we have proved the result for linear transformations. We can always factorize $F$ is

$$
\begin{equation*}
F=G \cdot A, \quad A=\left(F_{*}\right) \tag{8.47}
\end{equation*}
$$

so that the differential of $G$ at 0 is the identity, i.e.

$$
\begin{equation*}
G(x)=x+O\left(|x|^{2}\right) \tag{8.48}
\end{equation*}
$$

Now ( $\frac{17.11}{8.48}$ ) allows us to use an homotopy method, i.e. set
Lemma 8.3. If $F: \Omega \longrightarrow \mathbb{R}^{n}$ is a diffeomorphism of a neighbourhood of 0 onto its range, with $F(0)=0$, then

$$
\begin{equation*}
F^{*}\left\{u \in I_{c}^{m}\left(\mathbb{R}^{n},\{0\} ; \operatorname{supp}(u) \subset F(\Omega)\right\} \subset I_{c}^{m}\left(\mathbb{R}^{n},\{0\}\right)\right. \tag{8.44}
\end{equation*}
$$

$$
\begin{equation*}
G_{s}(x)=x+s(G(x)-x) \quad s \in[0,1) \tag{8.49}
\end{equation*}
$$

so that $G_{0}=\mathrm{Id}, G_{s}=G$. Such a 1-parameter family is given by integration of a vector field:

$$
\begin{align*}
& G_{s}^{*} \phi=\int_{0}^{s} \frac{d}{d s} G_{s}^{*} \phi d x \\
& =\int_{0} s \frac{d}{d s} \phi\left(G_{x}(x)\right) d s \\
& =\sum_{1} \int_{0}^{s} \frac{d}{G} d s\left(\partial x_{j} \phi\right)\left(G_{\delta}(x)\right) d s  \tag{8.50}\\
& =\int_{0}^{s} G_{s}^{*}\left(V_{s} \phi\right) d s
\end{align*}
$$

when the coefficients of $V_{s}$ are

$$
\begin{equation*}
G_{s}^{*} V_{s, j}=\frac{d}{d s} G_{s, i} \tag{8.51}
\end{equation*}
$$

${ }^{\text {Now }}$ by $\left(\frac{17.12}{17.49)} \frac{d}{d s} G_{s, i}=\Sigma x_{i} x_{j} a_{i j}^{s}(x)\right.$, so the same is true of the $V_{s, i}$, again using ( $\frac{17.49}{} .12$.

We can apply $\left(\frac{17.13}{8.50}\right)$ to compute

$$
\begin{equation*}
G^{*} u=\int_{0}^{1} G_{s}^{*}\left(V_{s} u\right) d s \tag{8.52}
\end{equation*}
$$

when $u \in I_{c}^{m}\left(\mathbb{R}^{n},\{0\}\right)$ has support near 0 . Namely, by $\left(\frac{16.33}{8.41}\right), u \in I H_{c}^{s}\left(\mathbb{R}^{n},\{0\}\right)$, with $s<-m-\frac{n}{4}$, but then

$$
\begin{equation*}
V_{s} u \in I H_{c}^{s+1}\left(\mathbb{R}^{n},\{0\}\right) \tag{8.53}
\end{equation*}
$$

since $V=\sum_{i, j=1}^{n} b_{i j}^{s}(x) x_{i} x_{j} D_{j}$. Applying $\left(\frac{16.33}{8.41)}\right.$ again gives

$$
\begin{equation*}
G_{s}^{*}\left(V_{s} u\right) \in I^{m^{\prime}}\left(\mathbb{R}^{n},\{0\}\right), \forall m^{\prime}>m-1 \tag{8.54}
\end{equation*}
$$

This proves the coordinates invariance.
Last time we defined the space of conormal distributions associated to a closed embedded submanifold $Y \subset X$ :

$$
\begin{align*}
& I H^{s}(X, Y)=\left\{u \in H^{s}(X) ; \mathcal{V}(X, Y)^{k} u \subset H^{s}(X) \forall k\right\} \\
& I H^{*}(X, Y)=I^{*}(X, Y)=\bigcup s I H^{s}(X, Y) \tag{8.55}
\end{align*}
$$

Here $\mathcal{V}(X, Y)$ is the space of $\mathcal{C}^{\infty}$ vector fields on $X$ tangent to $Y$. In the special case of a point in $\mathbb{R}^{n}$, say 0 , we showed that

$$
\begin{equation*}
\left.u \in I_{c}^{*}\left(\mathbb{R}^{n}\right),\{0\}\right) \Longleftrightarrow u \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right) \text { and } \hat{u} \in S^{M}\left(\mathbb{R}^{n}\right), M=M(u) \tag{8.56}
\end{equation*}
$$

In fact we then defined the "standard order filtration" by

$$
\begin{equation*}
u \in I_{c}^{m}\left(\mathbb{R}^{n},\{0\}\right)=\left\{u \in \mathcal{C}_{c}^{-\infty}\left(\mathbb{R}^{n}\right) ; \hat{u} \in S^{m-\frac{n}{4}}\left(\mathbb{R}^{n}\right)\right\} \tag{8.57}
\end{equation*}
$$

and found that
17.4

$$
\begin{equation*}
I H_{c}^{s}\left(\mathbb{R}^{n},\{0\}\right) \subset I_{c}^{-s-\frac{n}{4}}\left(\mathbb{R}^{n},\{0\}\right) \subset I H_{c}^{s^{\prime}}\left(\mathbb{R}^{n},\{0\}\right) \forall s^{\prime}<s \tag{8.58}
\end{equation*}
$$

Our next important task is to show that $I_{c}^{m}\left(\mathbb{R}^{n},\{0\}\right)$ is invariant under coordinate changes. That is, if $F: U_{1} \longrightarrow \mathbb{R}^{n}$ is a diffeomorphism of a neighbourhood of 0 to its range, with $F(0)=0$, then we want to show that

$$
\begin{equation*}
F^{*} u \in I_{c}^{m}\left(\mathbb{R}^{n},\{0\}\right) \forall u \in I_{c}^{m}\left(\mathbb{R}^{n},\{0\}\right), \operatorname{supp}(u) \subset F\left(U_{1}\right) \tag{8.59}
\end{equation*}
$$

Notice that we already know the coordinate independence of the Sobolev-based space, so using ( $\mathbf{8 . 5 8}$ ), we deduce that

$$
\begin{equation*}
F^{*} u \in I_{c}^{m^{\prime}}\left(\mathbb{R}^{n},\{0\}\right) \forall u \in I_{c}^{m}\left(\mathbb{R}^{n},\{0\}\right), n^{\prime}>m, \operatorname{supp}(u) \subset F\left(U_{1}\right) \tag{8.60}
\end{equation*}
$$

In fact we get quite a lot more for our efforts:

### 17.17

17.18

$$
\begin{equation*}
I^{m}\left(X,\{p\} ; \Omega^{\frac{1}{2}}\right) @>\sigma_{Y}^{m} \gg S^{m+\frac{n}{4}-[J]}\left(T_{p}^{*} \mathbb{R}^{n} ; \Omega^{\frac{1}{2}}\right) \tag{8.61}
\end{equation*}
$$

given by the local prescription

$$
\begin{equation*}
\sigma_{Y}^{m}(u)=\hat{u}(\xi)|d \xi|^{\frac{1}{2}} \tag{8.62}
\end{equation*}
$$

where $u=v|d x|^{\frac{1}{2}}$ is local coordinate based at 0 , with $\xi$ the dual coordinate in $T_{p}^{*} X$.
Proof. Our definition of $I^{m}\left(X,\{p\} ; \Omega^{\frac{1}{2}}\right)$ is just that in any local coordinate based at $p$

$$
\begin{equation*}
u \in I^{m}\left(X,\{p\} ; \Omega^{\frac{1}{2}}\right) \Longrightarrow \phi u=v|d x|^{\frac{1}{2}}, v \in I_{c}^{m}\left(\mathbb{R}^{n},\{0\}\right) \tag{8.63}
\end{equation*}
$$

and $u \in \mathcal{C}^{\infty}\left(X \backslash\{p\} ; \Omega^{\frac{1}{2}}\right)$. So the symbol map is clearly supposed to be

$$
\begin{equation*}
\sigma^{m}(u)^{(\zeta)} \equiv \equiv_{\downarrow} \hat{v}(\xi)|d \xi|^{\frac{1}{2}} \in S^{m+\frac{n}{4}-[1]}\left(\mathbb{R}^{n} ; \Omega^{\frac{1}{2}}\right) \tag{8.64}
\end{equation*}
$$

where $\zeta \in T_{10}^{*} X$ is the 1 -form $\zeta=\xi \cdot d x$ in the local coordinates. Of course we have to show that $\left(\frac{8.64)}{}\right.$ is independent of the choice of coordinates. We already know that $a_{1}$ change of coordinates changes $\hat{v}$ by a term of order $m-\frac{n}{4}-1$, which disappears in ( 8.64 ) 50 the residue class is determined by the Jacobian of the change of variables. From (18.46) we see exactly how $\hat{v}$ transforms under the Jacobian, namely as a density on

$$
\begin{aligned}
T_{0}^{*} \mathbb{R}^{n}: A \in G L(n, \mathbb{R}) \Longrightarrow \widehat{A^{*} v}(\eta)|d \eta|^{\frac{1}{2}} & \\
& =\hat{v}\left(\left(A^{t}\right)^{-1} \eta\right)|\operatorname{det} A|^{-1}|d y|
\end{aligned}
$$

so $\eta=A^{t} \xi \Longrightarrow$

$$
\begin{equation*}
\widehat{A^{*} v}(\eta)|d y|=\hat{v}(\xi)|d \xi| \tag{8.65}
\end{equation*}
$$

However recall from $\left(\frac{17.20}{8.63}\right)$ that $u$ is a half-density, so actually in the new coordinates $v^{\prime}=A^{*} v \cdot|\operatorname{det} A|^{\frac{1}{2}}$. This shows that ( $\left(\frac{17.21}{8.64)}\right.$ is well-defined.

Before going on to consider the general case let us note a few properties of $I^{m}\left(X,\{p\}, \Omega^{\frac{1}{2}}\right):$

Exercise: Prove that

$$
\text { If } P \in \operatorname{Diff}^{m}\left(X ; \Omega^{\frac{1}{2}}\right) \text { then }
$$

$$
\begin{gather*}
P: I^{m}\left(X,\{p\} ; \Omega^{\frac{1}{2}}\right) \longrightarrow I^{m+M}\left(X,\{p\} ; \Omega^{\frac{1}{2}}\right) \forall m  \tag{8.66}\\
\sigma^{m+M}(P u)=\sigma^{M}(P) \cdot \sigma^{m}(u)
\end{gather*}
$$

To pass to the general case of $Y \subset X$ we shall proceed in two steps. First let's consider a rather 'linear' case of $X=V$ a vector bundle over $Y$. Then $Y$ can be
identified with the zero section of $V$. In fact $V$ is locally trivial, i.e. each $p \in y$ has a neighbourhood $U$ s.t.

$$
\begin{equation*}
\pi^{-1}(U) \simeq \mathbb{R}_{x}^{n} \times U_{y^{\prime}}^{\prime} U^{\prime} \subset \mathbb{R}^{p} \tag{8.67}
\end{equation*}
$$

by a fibre-linear diffeomorphism projecting to a coordinate system on this base. So we want to define

$$
\begin{equation*}
I^{m}\left(V, Y ; \Omega^{\frac{1}{2}}\right)=\left\{u \in I^{*}\left(V, Y ; \Omega^{\frac{1}{2}}\right) ;\right. \tag{8.68}
\end{equation*}
$$

of $\phi \in \mathcal{C}_{c}^{\infty}(U)$ then under any trivialization $\left(\frac{17.24}{8.67}\right)$

$$
\begin{gathered}
\phi u(x, y) \equiv(2 \pi)^{-n} \int e^{i x \cdot \xi} a(y, \xi) d \xi|d x|^{\frac{1}{2}}, \quad \bmod \mathcal{C}^{\infty} \\
a \in S^{m-\frac{n}{2}-\frac{p}{4}}\left(\mathbb{R}_{y}^{p}, \mathbb{R}_{\xi}^{n}\right)
\end{gathered}
$$

Here $p=\operatorname{dim} Y, p+n=\operatorname{dim} V$. Of course we have to check that $\left(\frac{17.25}{8.69)}\right.$ is coordinateindependent. We can write the order of the symbol, corresponding to $u$ having order $m$ as

$$
\begin{equation*}
m-\frac{\operatorname{dim} V}{4}+\frac{\operatorname{dim} Y}{2}=m+\frac{\operatorname{dim} V}{4}-\frac{\operatorname{codim} Y}{2} \tag{8.70}
\end{equation*}
$$

These additional shifts in the order are only put there to confuse you! Well, actually they make life easier later.

Notice that we know that the space is invariant under any diffeomorphism of the fibres of $V$, varying smoothly with the base point, it is also obvious that ( $\frac{8.6 .25}{8}$ ) in independent the choice of coordinates is $U^{\prime}$, since that just transforms these variables. So a general change of variables preserving $Y$ is

$$
\begin{equation*}
(y, x) \longmapsto(f(y, x), X(y, x)) \quad X(y, 0)=0 \tag{8.71}
\end{equation*}
$$

In particular $f$ is a local diffeomorphism, which just changes the base variables in ( $\mathbf{B} \mathbf{8 . 6 2 5}$ ), so we can assume $f(y) \equiv y$. Then $X(y, x)=A(y) \cdot x+O\left(x^{2}\right)$. Since $x \longmapsto A(y) \cdot x$ is a fibre-by-fibre transformation it leaves the space invariant too, So we are reduced to considering

$$
\begin{equation*}
G:(y, x) \longmapsto\left(y, x+\Sigma a_{i j}(x, y) x_{i} x_{j}\right) y+\Sigma b_{i}(x, y) x_{i} . \tag{8.72}
\end{equation*}
$$

To handle these transformations we can use the same homotopy method as before i.e.
17.28

$$
\begin{equation*}
G_{s}\left(x, y=(y+s) \sum_{i} b_{i}(x, y) x_{i}, x+s \sum_{i, j} a_{i j}(x, y) x_{i} x_{j}\right) \tag{8.73}
\end{equation*}
$$

is a 1-parameter family of diffeomorphisms. Moreover

$$
\begin{equation*}
\frac{d}{d s} G_{s}^{*} u=G_{s}^{*} V_{s} k \tag{8.74}
\end{equation*}
$$

where

$$
\begin{equation*}
V_{s}=\sum_{i, \ell} \beta_{i, \ell}(s, x, y) x_{i} \partial_{y \ell}+\sum_{i, j, k} \alpha_{i, j, k}+\sum_{i, j, k} \alpha_{i j k}(\alpha, y, s) \ell_{i}, \ell_{j} \frac{\partial}{\partial x_{k}} \tag{8.75}
\end{equation*}
$$

So all we really have to show is that

$$
\begin{equation*}
V_{s}: I^{M}\left(U^{\prime} \times \mathbb{R}^{n}, U^{\prime} \times\{0\}\right) \longrightarrow I^{M-1}\left(U^{\prime} \times \mathbb{R}^{n}, U^{\prime} \times\{0\}\right) \forall M \tag{8.76}
\end{equation*}
$$

Again the spaces are $\mathcal{C}^{\infty}$-modules so we only have to check the action of $x_{i} \partial_{y_{\ell}}$ and $x_{i} x+j \partial_{x_{k}}$. These change the symbol to

$$
\begin{equation*}
D_{\xi_{i}} \partial_{y_{\ell}} a \text { and } i D_{\xi_{i}} D_{\xi_{j}} \cdot \xi_{k} a \tag{8.77}
\end{equation*}
$$

respectively, all one order lower.
This shows that the definition $\left(\frac{17.25}{8.69}\right)$ is actually a reasonable one, i.e. as usual it suffices to check it for any covering by coordinate partition.

Let us go back and see what the symbol showed before.

### 17.30 Lemma 8.5. If

$$
\begin{equation*}
u \in I^{m}\left(V, Y ; \Omega^{\frac{1}{2}}\right) u=v|d x|^{\frac{1}{2}}|d \xi|^{\frac{1}{2}} \tag{8.78}
\end{equation*}
$$

defines an element

$$
\begin{equation*}
\sigma^{m}(u) \in S^{m+\frac{n}{4}+\frac{p}{4}-[1]}\left(V^{*} ; \Omega^{\frac{1}{2}}\right) \tag{8.79}
\end{equation*}
$$

independent of choices.
Last time we discussed the invariant symbol for a conormal distribution associated to the zero section of a vector bundle. It turns out that the general case is not any more complicated thanks to the "tubular neighbourhood" or "normal fibration" theorem. This compares $Y \hookrightarrow X$, a closed embedded submanifold, to the zero section of a vector bundle.

Thus at each point $y \in Y$ consider the normal space:

$$
\begin{equation*}
N_{y} Y=N_{y}\{X, Y\}=T_{y} x / T_{y} Y \tag{8.80}
\end{equation*}
$$

That is, a normal vector is just any tangent vector to $X$ modulo tangent vectors to $Y$. These spaces define a vector bundle over $Y$ :

$$
\begin{equation*}
N Y=N\{X ; Y\}=\bigsqcup_{y \in Y} N_{y} Y \tag{tabular}
\end{equation*}
$$

where smoothness of a section is inherited from smoothness of a section of $T_{y} X$, i.e.

$$
N Y=T_{y} X / T_{y} Y
$$

Suppose $Y_{i} \subset X_{i}$ are $\mathcal{C}^{\infty}$ submanifolds for $i=1,2$ and that $F: X_{1} \longrightarrow X_{2}$ is a $\mathcal{C}^{\infty}$ map such that
18.4 (8.83)

$$
F\left(Y_{1}\right) \subset Y_{2}
$$

Then $F_{*}: T_{y} X_{1} \longrightarrow T_{F(y)} X_{2}$, must have the property

$$
\begin{equation*}
F_{*}: T_{y} Y_{1} \longrightarrow T_{F(y)} Y_{2} \forall y \in Y_{1} . \tag{8.84}
\end{equation*}
$$

This means that $F_{*}$ defines a map of the normal bundles


Notice the very special case that $W \longrightarrow Y$ is a vector bundle, and we consider $Y \hookrightarrow W$ as the zero section. Then

$$
\begin{equation*}
N_{y}\{W ; Y\} \longleftrightarrow W_{y} \quad \forall y \in Y \tag{8.86}
\end{equation*}
$$

since
18.8 (8.87)

$$
T_{y} W=T_{y} Y \oplus T_{y}\left(W_{y}\right) \quad \forall y \in W
$$

That is, the normal bundle to the zero section is naturally identified with the vector bundle itself.

So, suppose we consider $\mathcal{C}^{\infty}$ maps

$$
\begin{equation*}
f: B \longrightarrow N\{X ; Y\}=N Y \tag{8.88}
\end{equation*}
$$

where $B \subset X$ is an open neighbourhood of the submanifold $Y$. We can demand that

$$
\begin{equation*}
f(y)=(y, 0) \in N_{y} Y \quad \forall y \in Y \tag{8.89}
\end{equation*}
$$

which is to say that $f$ induces the natural identification of $Y$ with the zero section of $N Y$ and moreover we can demand

$$
\begin{equation*}
f_{*}: N Y \longrightarrow N Y \text { is the identity. } \tag{8.90}
\end{equation*}
$$

Here $f_{*}$ is the map $\left(\frac{18.6}{18.85}\right)$, so maps $N Y$ to the normal bundle to the zero section of $N Y$, which we have just observed is naturally just $N Y$ again.

Theorem 8.1. For any closed embedded submanifold $Y_{9} \subset X$ there exists a normal fibration, i.e. a diffeomorphism (onto its range) ( $\mathbf{8 . 8 8}$ ) satisfing ( $\overline{8.89}$ ) and (8.90); two such maps $f_{1}, f_{2}$ are such that $g=f_{2} \circ f_{1}^{-1}$ is a diffeomorphism neqr the zero section of $N Y$, inducing the identity on $Y$ and inducing the identity $\left(\frac{18.90}{8.90}\right)$.

Proof. Not bad, but since it uses a little Riemannian geometry I will not prove it, see [ ], [ ]. (For those who know a little Riemannian geometry, $f^{-1}$ can be taken as the exponential map near the zero section of $N Y$, identified as a subbundle of $T_{Y} X$ using the metric.) Of course the uniqueness part is obvious.

Actually we do not really need the global aspects of this theorem. Locally it is immediate by using local coordinates in which $Y=\left\{x_{1}=\cdots=x_{k}=0\right\}$.

Anyway using such a normal fibration of $X$ near $Y$ (or working locally) we can simply define

$$
\begin{gather*}
I^{m}\left(X, Y ; \Omega^{\frac{1}{2}}\right)=\left\{u \in \mathcal{C}^{-\infty}\left(X ; \Omega^{\frac{1}{2}}\right) ; u \text { is } \mathcal{C}^{\infty} \text { in } X \backslash Y\right. \text { and }  \tag{8.91}\\
\left.\left(f^{-1}\right)^{*}(\phi u) \in I^{m}\left(N Y, Y ; \Omega^{\frac{1}{2}}\right) \text { if } \phi \in \mathcal{C}^{\infty}(X), \operatorname{supp}(\phi) \subset B\right\}
\end{gather*}
$$

Naturally we should check that the definition doesn't depend on the choice of $f$. This means knowing that $I^{m}\left(N Y, Y ; \Omega^{\frac{1}{2}}\right)$ is invariant under $g$, as in the theorem, but we have already checked this. In fact notice that $g$ is exactly of the type of ( 8.72 ). Thus we actually know that

$$
\begin{equation*}
\sigma^{m}\left(g^{*} u\right)=\sigma^{m}(u) \text { in } S^{m+\frac{n}{4}+\frac{p}{4}-[1]}\left(N^{*} Y ; \Omega^{\frac{1}{2}}\right) \tag{8.92}
\end{equation*}
$$

So we have shown that there is a coordinate invariance symbol map

$$
\begin{equation*}
\sigma^{m}: I^{m}\left(X, Y ; \Omega^{\frac{1}{2}}\right) \longrightarrow S^{m+\frac{n}{4}+\frac{p}{4}-[1]}\left(N^{*} Y ; \Omega^{\frac{1}{2}}\right) \tag{8.93}
\end{equation*}
$$

giving a short exact sequence
18.15
$0 \hookrightarrow I^{m-1}\left(X, Y ; \Omega^{\frac{1}{2}}\right) \longrightarrow I^{m}\left(X, Y ; \Omega^{\frac{1}{2}}\right) @>\sigma^{m} \gg S^{m+\frac{n}{4}+\frac{p}{4}-[1]}\left(N^{*} Y ; \Omega^{\frac{1}{2}}\right) \longrightarrow 0$

$$
\begin{equation*}
\text { where } n=\operatorname{dim} X-\operatorname{dim} Y, p=\operatorname{dim} Y \tag{8.95}
\end{equation*}
$$

Asymptotic completeness carries over immediately. We also need to go back and check the extension of ( $\mathbf{8 . 6 . 6}$ ):
18.16 Proposition 8.1. If $Y \hookrightarrow X$ is a closed embedded submanifold and $A \in$ $\Psi_{c}^{m}\left(X ; \Omega^{\frac{1}{2}}\right)$ then

$$
\begin{equation*}
A: I^{M}\left(X, Y ; \Omega^{\frac{1}{2}}\right) \longrightarrow I^{M+m}\left(X, Y ; \Omega^{\frac{1}{2}}\right) \forall M \tag{8.96}
\end{equation*}
$$

and
18.18
(8.97)

$$
\sigma^{m+M}(A u)=\sigma^{m}(A) \sigma^{m}(A) \upharpoonright N^{*} Y \sigma^{M}(u)
$$

Notice that $\sigma^{m}(A) \in S^{m-[1]}\left(T^{*} X\right)$ so the product here makes perfectly good sense.
Proof. Since everything in sight is coordinate-independent we can simply work in local coordinates where

$$
\begin{equation*}
X \sim \mathbb{R}_{y}^{p} \times \mathbb{R}_{x}^{n}, Y=\{x=0\} \tag{8.98}
\end{equation*}
$$

Then $u \in I_{c}^{m}\left(X, Y ; \Omega^{\frac{1}{2}}\right)$ means just

$$
u=(2 \pi)^{-n} \int e^{i x \cdot \xi} a(y, \xi) d \xi \cdot|d x|^{\frac{1}{2}}, a \in S^{m-\frac{n}{4}+\frac{p}{4}}\left(\mathbb{R}^{p}, \mathbb{R}^{n}\right)
$$

Similarly $A$ can be written in the form

$$
\begin{equation*}
A=(2 \pi)^{-n-p} \int e^{i\left(x-x^{\prime}\right) \cdot \xi+i\left(y-y^{\prime}\right) \cdot \eta} b(x, y, \xi, \eta) d \xi d \eta \tag{8.100}
\end{equation*}
$$

Using the invariance properties of the Sobolev based space if we write

$$
\begin{equation*}
A=A_{0}+\Sigma x_{j} B_{j}, A_{0}=q_{L}(b(0, y, \xi, \eta)) \tag{8.101}
\end{equation*}
$$

we see that $A u \in I^{m+M}\left(X, Y ; \Omega^{\frac{1}{2}}\right)$ is equivalent to $A_{0} u \in I^{m+M}\left(X, Y ; \Omega^{\frac{1}{2}}\right)$. Then

$$
\begin{equation*}
A_{0} u=(2 \pi)^{-n-p} \int e^{i x \cdot \xi+i\left(y-y^{\prime}\right) \cdot \eta} b\left(0, y^{\prime}, \xi, \eta\right) b\left(y^{\prime}, \xi\right) d y^{\prime} d \eta d \xi \tag{8.102}
\end{equation*}
$$

where we have put $A_{0}$ in right-reduced form. This means

$$
\begin{equation*}
A_{0} u=(2 \pi)^{-n} \int e^{i x \cdot \xi} c(y, \xi) d \xi \tag{8.103}
\end{equation*}
$$

where

$$
\begin{equation*}
c(y, \xi)=(2 \pi)^{-p} \int e^{i\left(y-y^{\prime}\right) \cdot \eta} b\left(0, y^{\prime}, \xi, \eta\right) a\left(y^{\prime}, \xi\right) d y^{\prime} d \eta \tag{8.104}
\end{equation*}
$$

Regarding $\xi$ as a parameter, this is, before $y^{\prime}$ integration, the kernel of a pseudodifferential operator is $y$. It can therefore be written in left-reduced form, i.e.
18.22

$$
\begin{equation*}
c(y, \xi)=(2 \pi)^{-p} \int e^{i\left(y-y^{\prime}\right) \eta} e(y, \xi, \eta) d \eta d y^{\prime}=e(y, \xi, 0) \tag{8.105}
\end{equation*}
$$

where $e(y, \xi, \eta)=b\left(0, y, \xi_{18}\right)_{18}(y, \xi)$ plus terms of order at most $m+M-\frac{n}{4}+\frac{p}{4}-1$. This proves the formula ( $\overline{8.977}$ ).

Notice that if $A$ is elliptic then $A u \in \mathcal{C}^{\infty}$ implies $u \in \mathcal{C}^{\infty}$, i.e. there are no singular solutions. Suppose that $P$ is say a differential operator which is not elliptic and we look for solutions of

$$
\begin{equation*}
P u \in \mathcal{C}^{\infty}\left(X \Omega^{\frac{1}{2}}\right) \tag{8.106}
\end{equation*}
$$

How can we find them? Well suppose we try
18.24
$u \in I^{M}\left(X, Y ; \Omega^{\frac{1}{2}}\right)$
for some submanifold $Y$. To know that $u$ is singular we will want to have

$$
\begin{equation*}
\sigma(u) \text { is elliptic on } N^{*} Y \tag{8.108}
\end{equation*}
$$

(which certainly implies that $u \notin \mathcal{C}^{\infty}$ ).
18. The simplest case would be $Y$ a hypersurface. In any case from $\left(\frac{18.18}{8.97}\right)$ and (8.106) we deduce

$$
\begin{equation*}
\sigma^{m}(P) \cdot \sigma^{M}(u) \equiv 0 \tag{8.109}
\end{equation*}
$$

So if we assume $\left(\frac{18.25}{8.108}\right)$ then we must have

$$
\begin{equation*}
\sigma^{m}(P) \upharpoonright N^{*} Y=0 \tag{8.110}
\end{equation*}
$$

Definition 8.3. A submanifold is said to be characteristic for a given operator $P \in \operatorname{Diff}^{m}\left(X ; \Omega^{\frac{1}{2}}\right)$ if $\left(\frac{18.27}{8.110}\right)$ holds.
[18.23f course even if $P$ is characteristic for $y$, and so $\left(\frac{18.26}{8.109)}\right.$ holds we do not recover ( 18.106 ), just

$$
\begin{equation*}
P u \in I^{m+M-1}\left(X, Y ; \Omega^{\frac{1}{2}}\right) \tag{8.111}
\end{equation*}
$$

i.e., one order smoother than it "should be". The task might seem hopeless, but let us note that these are examples, and important ones at that!!

Consider the (flat) wave operator

$$
\begin{equation*}
P=P_{t}^{2}-\sum_{i=1}^{n} D_{i}^{2}=D_{t}^{2}-\Delta \text { on } \mathbb{R}^{n+1} \tag{8.112}
\end{equation*}
$$

A hypersurface in $\mathbb{R}^{n+1}$ looks like

$$
\begin{equation*}
H=\{h(t, x)=0\},(d h \neq 0 \text { on } H) . \tag{8.113}
\end{equation*}
$$

The symbol of $P$ is

$$
\begin{equation*}
\sigma^{2}(P)=\tau^{2}-|\xi|^{2}=\tau^{2}-\xi_{1}^{2}-\cdots-\xi_{n}^{2} \tag{8.114}
\end{equation*}
$$

where $\tau, \xi$ are the dual variables to $t, x$. So consider ( $\left.\frac{18.27}{8.110}\right)$,

$$
\begin{equation*}
N^{*} Y=\{(t, x ; \lambda d h(t, y)) ; h(t, x)=0\} \tag{8.115}
\end{equation*}
$$

Inserting this into $\left(\frac{18.32}{8.114}\right)$ we find:

$$
\begin{equation*}
\left(\lambda \frac{\partial h}{\partial t}\right)^{2}-\left(\lambda \frac{\partial h}{\partial x_{1}}\right)^{2}-\cdots-\left(\lambda \frac{\partial h}{\partial x_{n}}\right)^{2}=0 \text { on } h=0 \tag{8.116}
\end{equation*}
$$

i.e. simply:

$$
\begin{equation*}
\left(\frac{\partial h}{\partial t}\right)^{2}=\left|d_{x} h\right|^{2} \text { on } h(t, x)=0 \tag{8.117}
\end{equation*}
$$

This is the "eikonal eqguation" for $h$ (and hence $H$ ).
Solutions to ( $\overline{8.117})$ are easy to find - we shall actually find all of them (locally) next time. Examples are given by taking $h$ to be linear:

$$
\begin{equation*}
H=\{h=a t+b \cdot x=0\} \text { is characteristic for } P \Longleftrightarrow a^{2}=|b|^{2} \tag{8.118}
\end{equation*}
$$

Since $h / a$ defines the same surface, all the linear solutions correspond to planes

$$
\begin{equation*}
t=\omega \cdot x, \omega \in \mathbb{S}^{n-1} \tag{8.119}
\end{equation*}
$$

So, do solutions of $P u \in \mathcal{C}^{\infty}$ which are conormal with respect to such hypersurfaces exist? Simply take

$$
\begin{equation*}
u=v(t-\omega \cdot x) \quad v \in I^{*}\left(\mathbb{R},\{0\} ; \Omega^{\frac{1}{2}}\right) \tag{8.120}
\end{equation*}
$$

$$
\begin{equation*}
P u=0, u \in I^{*}\left(\mathbb{R}^{n+1}, H ; \Omega^{\frac{1}{2}}\right) \tag{8.121}
\end{equation*}
$$

For example $v(s)=\delta(s), u=\delta(t-\omega \cdot x)$ is a "travelling wave".

### 8.1. Hamilton-Jacobi theory

Let $X$ be a $\mathcal{C}^{\infty}$ manifold and suppose $p \in \mathcal{C}^{\infty}\left(T^{*} X \backslash 0\right)$ is homogeneous of degree $m$. We want to find characteristic hypersurfaces for $p$, namely hypersurfaces (locally) through $\bar{x} \in X$
19.1

$$
\begin{equation*}
H=\{f(x)=0\} \quad h \in \mathcal{C}^{\infty}(x) h(\bar{x})=0, d h(\bar{x}) \neq 0 \tag{8.122}
\end{equation*}
$$

such that
19.2 (8.123)

$$
p(x, d h(x))=0
$$

Here we demand that $\left(\frac{19.2}{8.12} 3\right)$ hold near $\bar{x}$, not just on $H$ itself. To solve ( $\frac{19.2}{8.12} 3$ ) we need to impose some additional conditions, we shall demand

$$
\begin{equation*}
p \text { is real-valued } \tag{8.124}
\end{equation*}
$$

and

$$
\begin{equation*}
d_{\text {fibre }} p \neq 0 \text { or } \Sigma(p)=\{p=0\} \subset T^{*} X \backslash 0 \tag{8.125}
\end{equation*}
$$

This second condition is actually stronger than really needed (as we shall see) but in any case it implies that

$$
\begin{equation*}
\Sigma(P) \subset T^{*} X \backslash 0 \text { is a } \mathcal{C}^{\infty} \text { conic hypersurface } \tag{8.126}
\end{equation*}
$$

by the implicit function theorem.
The strategy for solving ( 1.123 ) is a geometric one. Notice that

$$
\begin{equation*}
\Lambda_{h}=\left\{(x, d h(x)) \in T^{*} X \backslash 0\right\} \tag{8.127}
\end{equation*}
$$

actually determines $h$ up to an additive constant. The first question we ask is precisely which submanifold $\Lambda \subset T^{*} X \backslash 0$ corresponds to graphs of differentials of $\mathcal{C}^{\infty}$ functions? The answer to this involves the tautologous contact form.
19.7

$$
\begin{gather*}
\alpha: T^{*} X \longrightarrow T^{*}\left(T^{*} X\right) \not \subset \tilde{\pi} \circ \alpha=\mathrm{Id} \\
\alpha(x, \xi)=\tilde{\pi}^{*} \xi \in T_{(x, \xi)}^{*}\left(T^{*} X\right) . \tag{8.128}
\end{gather*}
$$

Here $\tilde{\pi}: T^{*}\left(T^{*} X\right) \longrightarrow T^{*} X$ is the projection. Notice that if $x_{1}, \ldots, x_{n}$ are local coordinates in $X$ then $x_{1}, \ldots, x_{n}, \xi_{1}, \ldots, \xi_{n}$ are local coordinates $T^{*} X$, where $\xi \in$ $T_{x}^{*} X$ is written

$$
\begin{equation*}
\xi=\sum_{i=1}^{n} \xi_{i} d x_{i} \tag{8.129}
\end{equation*}
$$

Since $x_{1}, \ldots, x_{n}, \xi_{1}, \ldots, \xi_{n}$ are local coordinates in $T^{*} X$ they together with the dual coordinates $\Xi_{1}, \ldots, \Xi_{n}, X_{1}, \ldots, X_{n}$ are local coordinates in $T^{*}\left(T^{*} X\right)$ where

$$
\begin{equation*}
\zeta \in T_{(x, \xi)}^{*}\left(T^{*} X\right) \Longrightarrow \zeta=\sum_{j=1}^{n} \Xi_{j} d x_{j}+\sum_{j=1}^{n} X_{j} d \xi_{j} \tag{8.130}
\end{equation*}
$$

In these local coordinates

$$
\begin{equation*}
\alpha=\sum_{j=1}^{n} \xi_{j} d x_{j}! \tag{8.131}
\end{equation*}
$$

The first point is. that $\alpha$ is independent of the original choice of coordinates, as is evident from (8.128).
19.11 Lemma 8.6. A submanifold $\Lambda \subset T^{*} X \backslash 0$ is, near $(\bar{x}, \bar{\xi}) \in \Lambda$, of the form ( ${ }^{\frac{19}{8} .127}$ ) for some $h \in \mathcal{C}^{\infty}(X)$, if $\pi: \Lambda \longrightarrow X$ is a local diffeomorphism
and
$\alpha$ restricted to $\Lambda$ is exact.
Proof. The first condition, ( $\left(18.12_{192}^{12}\right)$, means that $\Lambda$ is locally the image of a section of $T^{*} X$ :

$$
\begin{equation*}
\Lambda=\left\{(x, \zeta(x)), \zeta \in \mathcal{C}^{\infty}\left(X ; T^{*} X\right)\right\} \tag{8.134}
\end{equation*}
$$

Thus the section $\zeta$ gives an inverse $Z$ to $\pi$ in ( ${ }^{\frac{1}{8.1} .122}$ ). It follows from $\left(\frac{19.7}{8.12} 8\right)$ that (8.135)

$$
Z^{*} \alpha=\zeta
$$

Thus if $\alpha$ is exact on $\Lambda$ then $\zeta$ is exact on $X, \zeta=d h$ as required.
Of course if we are only working locally near some point $(\bar{x}, \bar{\xi}) \in \Lambda$ then (19.13 can be replaced by the condition

$$
\begin{equation*}
\omega=d \alpha=0 \text { on } X \tag{8.136}
\end{equation*}
$$

Here $\omega=d \alpha$ is the symplectic form on $T^{*} X$ :

$$
\begin{equation*}
\omega=\sum_{j=1}^{n} d \xi_{j} \wedge d x_{j} \tag{8.137}
\end{equation*}
$$

Definition 8.4. A submanifold $\Lambda \subset T^{*} X$ of dimension equal to that of $X$ is said to be Lagrangian if the fundamental 2-form, $\omega$, vanishes when pulled back to $\Lambda$.

By definition a symplectic manifold is a $\mathcal{C}^{\infty}$ manifold $S$ with a $\mathcal{C}^{\infty} 2$-form $\omega \in \mathcal{C}^{\infty}\left(S ; \Lambda^{2}\right)$ fixed satisfying two constraints

$$
\begin{equation*}
\omega \underset{n \text { factors }}{\wedge \ldots \wedge} \omega \neq 0 \quad \operatorname{dim} S=2 n \tag{8.138}
\end{equation*}
$$

A particularly simple example of a symplectic manifold is a real vector space, necessarily of even dimension, with a non-degenerate antisymmetric 2-form:

$$
\left\{\begin{array}{l}
\omega: E \times E \longrightarrow \mathbb{R}  \tag{8.140}\\
\tilde{\omega}: E \longleftrightarrow E^{*}
\end{array}\right.
$$

Here $\tilde{\omega}(v)(w)=\omega(v, w) \forall w \in E$. Now $\left(\frac{19.17}{8.138}\right)$ is trivially true if we think of $\omega$ as a translation-invariant 2 -form on $E$, thought of as a manifold.

Then a subspace $V \subset E$ is Lagrangian if

$$
19.20
$$

$$
\begin{align*}
& \omega(v, w)=0 \forall v, w \in V \\
& 2 \operatorname{dim} V=\operatorname{dim} E \tag{8.141}
\end{align*}
$$

Of course the point of looking at symplectic vector spaces and Lagrangian subspaces is:
19.21 Lemma 8.7. If $S$ is a symplectic manifold then $T_{z} S$ is a symplectic vector space for each $z \in S$. A submanifold $\Lambda \subset S$ is Lagrangian iff $T_{z} \Lambda \subset T_{z} S$ is a Lagrangian subspace $\forall z \in \Lambda$.

We can treat $\omega$, the antisymmetric 2-form on $E$, as though it were a Euclidean inner product, at least in some regards! Thus if $W \subset E$ is any subspace set

$$
\begin{equation*}
W^{\omega}=\{v \in E ; \omega(v, w)=0 \forall w \in W\} . \tag{8.142}
\end{equation*}
$$

LEMMA 8.8. If $W \subset E$ is a linear subspace of a symplectic vector space then $\operatorname{dim} W^{\omega}+\operatorname{dim} W=\operatorname{dim} E ; W$ is Lagrangian if and only if

$$
\begin{equation*}
W^{\omega}=W \tag{8.143}
\end{equation*}
$$

Proof. Let $W^{0} \subset E^{*}$ be the usual annihilator:

$$
\begin{equation*}
W^{0}=\left\{\alpha \in E^{*} ; \alpha(v)=0 \forall v \in W\right\} \tag{8.144}
\end{equation*}
$$

Then $\operatorname{dim} W^{0}=\operatorname{dim} E-\operatorname{dim} W$. Observe that

$$
\begin{equation*}
\tilde{\omega}: W^{\omega} \longleftrightarrow W^{0} . \tag{8.145}
\end{equation*}
$$

Indeed if $\alpha \in W^{0}$ and $\tilde{\omega}(v)=\alpha$ then

$$
\begin{equation*}
\alpha(w)=\tilde{\omega}(v)(w)=\omega(v, w)=0 \forall w \in W \tag{8.146}
\end{equation*}
$$

implies that $v \in W^{\omega}$. Conversely if $v \in W^{\omega}$ then $\alpha=\tilde{\omega}(v) \in W^{0}$. Thus $\operatorname{dim} W^{\omega}+$ $\operatorname{dim} W=\operatorname{dim} E$.

Now if $W$ is Lagrangian then $\alpha=\tilde{\omega}(w), w \in W$ implies

$$
\begin{equation*}
\alpha(v)=\tilde{\omega}(w)(v)=\omega(w, v)=0 \forall v \in w \tag{8.147}
\end{equation*}
$$

Thus $\tilde{\omega}(W) \subset W^{0} \Longrightarrow W \subset W^{\omega}$, by (19.25 $)$, and since $\operatorname{dim} W=\operatorname{dim} W^{\omega}$, (19.24 183 ) holds. The converse follows similarly.

The "lifting" isomorphism $\tilde{\omega}: E \longleftrightarrow E^{*}$ for a symplectic vector space is like the Euclidean identification of vectors and covectors, but "twisted". It is of fundamental importance, so we give it several names! Suppose that $S$ is a symplectic manifold. Then

$$
\begin{equation*}
\tilde{\omega}_{z}: T_{z} S \longleftrightarrow T_{z}^{*} S \forall z \in S \tag{8.148}
\end{equation*}
$$

This means that we can associate (by the inverse of ( ${ }^{(19.148)}$ ) a vector field with each 1-form. We write this relation as

$$
\begin{gather*}
H_{\gamma} \in \mathcal{C}^{\infty}(S ; T S) \text { if } \gamma \in \mathcal{C}^{\infty}\left(S ; T^{*} S\right) \text { and } \\
\tilde{\omega}_{z}\left(H_{\gamma}\right)=\gamma \forall z \in S . \tag{8.149}
\end{gather*}
$$

Of particular importance is the case $\gamma=d f, f \in_{11} \mathcal{G}_{.27}^{\infty}(S)$. Then $H_{d f}$ is written $H_{f}$ and called the Hamilton vector field of $f$. From (8.149)

$$
\begin{equation*}
\omega\left(H_{f}, v\right)=d f(v)=v f \forall v \in T_{z} S, \forall z \in S \tag{8.150}
\end{equation*}
$$

The identity $\left(\frac{19.28}{8.150}\right)$ implies one important thing immediately:

$$
\begin{equation*}
H_{f} f \equiv 0 \forall f \in \mathcal{C}^{\infty}(S) \tag{8.151}
\end{equation*}
$$

since

$$
\begin{equation*}
H_{f} f=d f\left(H_{f}\right)=\omega\left(H_{f}, H_{f}\right)=0 \tag{8.152}
\end{equation*}
$$

by the antisymmetry of $\omega$. We need a generalization of this:
Lemma 8.9. Suppose $L \subset S$ is a Lagrangian submanifold of a symplectic manifold then for each $f \in \mathcal{I}(S)=\left\{f \in \mathcal{C}^{\infty}(X) ; f \upharpoonright\{s=0\}, H_{f}\right.$ is tangent to $\Lambda$.

Proof. $H_{f}$ tangent to $\Lambda$ means $H_{f}(z) \in T_{z} \Lambda \forall z \in$ 19. If $f=0$ on $\Lambda$ then $d f=0$ on $T_{z} \Lambda$, i.e. $d f(z) \in\left(T_{z} \Lambda\right)^{0} \subset\left(T_{z} S\right) \forall z \in \Lambda$. By (8.143) the assumption that $\Lambda$ is Lagrangian means $\tilde{\omega}_{z}(d f(z)) \in T_{z} \Lambda$, i.e. $H_{f}(z) \in T_{\zeta} \Lambda$ as desired.

This lemma gives us a necessary condition for our construction of a Lagrangian submanifold

$$
\begin{equation*}
\Lambda \subset \Sigma(P) \tag{8.153}
\end{equation*}
$$

Namely $H_{p}$ must be tangent to $\Lambda$ ! We use this to construct $\Lambda$ as a union of integral curves of $H_{p}$. Before thinking agbout this seriously, let's look for a moment at the conditions we imposed on $p$, (18.124) and (8.125). If $p$ is real then $H_{p}$ is real (since $\omega$ is real). Notice that

If $S=T^{*} X$ then each fibre $T_{x}^{*} X \subset T^{*} X$ is Lagrangian .
Remember that on $T^{*} X, \omega=d \alpha, \alpha=\xi \cdot d x$ the canonical 1-form. Thus $T_{x}^{*} X$ is just $x_{11 \overline{9.32}}$ const, so $d x=0$, so $\alpha=0$ on $T^{*} X$ and hence in particular $\omega=0$, proving ( $\frac{10.32}{8.154)}$. This allows us to interpret $\left(\frac{19.4}{8.125)}\right.$ in terms of $H_{p}$ as
(8.155) $\quad\left(\frac{19.4}{8.12} 5\right) \longleftrightarrow H_{p}$ is everywhere transversal to the fibres $T_{x}^{*} X$.

Now we want to construct a little piece of Lagrangian manifold satisfying (19.151). Suppose $z \in \Sigma(P) \subset T^{*} X \backslash 0$ and we want to construct a piece of $\Lambda$ through $z$. Since $\pi_{*}\left(H_{p}(z)\right) \neq 0$ we can choose a local coordinate, $t \in \mathcal{C}^{\infty}(X)$, such that

$$
\begin{equation*}
\pi_{*}\left(H_{p}(z)\right) t \neq 0, \text { i.e. } H_{p}\left(\pi^{*} t\right)(z) \neq 0 \tag{8.156}
\end{equation*}
$$

Consider the hypersurface through $\pi(z) \in X$,

$$
\begin{equation*}
H=\{t=t(z)\} \Longrightarrow \pi(z) \in H \tag{8.157}
\end{equation*}
$$

Suppose $f \in \mathcal{C}^{\infty}(H), d f(\pi(z))=0$. In fact we can choose $f$ so that

$$
\begin{equation*}
f=f^{\prime} \upharpoonright H, f^{\prime} \in \mathcal{C}^{\infty}(X), d f^{\prime}(\pi(z))=z \tag{8.158}
\end{equation*}
$$

where $z \in \Xi(P)$ was our chosen base point.
19.37 THEQREM 8.2. (Hamilton-Jacobi) Suppose $p \in \mathcal{C}^{\infty}\left(T^{*} X \backslash 0\right.$ ) sqtis.3ies ( $\left(\frac{19.3}{8, i .34} 1\right)$
 and $f \in \mathcal{C}^{\infty}(H)$ satisfies (15.158), then there exists $\tilde{f} \in \mathcal{C}^{\infty}(X)$ such that

$$
\begin{gather*}
\Lambda=\operatorname{graph}(d \tilde{f}) \subset \Sigma(P) \text { near } z \\
\tilde{f} \upharpoonright H=f \text { near } \pi(z)  \tag{8.159}\\
d \tilde{f}(\pi(z))=z
\end{gather*}
$$

and any other such solution, $\tilde{f}^{\prime}$, is equal to $\tilde{f}$ in a neighbourhood of $\pi(z)$.
Proof. We need to do a bit more work to prove this important theorem, but let us start with the strategy. First notice that $\Lambda \cap \pi^{-1}(H)$ is already determined, near $\pi(z)$.

To see this we have to understand the relationship between $d f(h) \in T^{*} H$ and $d \tilde{f}(h) \in T^{*} X, h \in H, \tilde{f} \upharpoonright H=f$. Observe that $H=\{t=0\}$ lifts to $T_{H}^{*} X \subset T^{*} X$ a
hypersurface. By $\left(\frac{19.29}{8.151}\right), H_{t}$ is tangent to $T_{H}^{*} X$ and non-zero. In local coordinates $t, x, \ldots, x_{n-1}$, the $x$ 's in $H$,

$$
\begin{equation*}
H_{t}=-\frac{\partial}{\partial \tau} \tag{8.160}
\end{equation*}
$$

where $\tau, \xi_{1}, \ldots, \xi_{n}$ are the dual coordinates. Thus we see that

$$
\begin{equation*}
\pi_{H}: T_{H}^{*} X \longrightarrow T^{*} H \quad \pi_{H}(\beta)(v)=\beta(v), v \in T_{h} H \subset T_{h} X \tag{8.161}
\end{equation*}
$$

is projection along $\partial_{\tau}$. Now starting from $f \in \mathcal{C}^{\infty}(H)$ we have

$$
\begin{equation*}
\Lambda_{f} \subset T^{*} H \tag{8.162}
\end{equation*}
$$

Notice that if $\tilde{f} \in \mathcal{C}^{\infty}(X), \tilde{f} \mid H=f$ then

$$
\begin{equation*}
\Lambda_{\tilde{f}} \cap T_{H}^{*} X \text { has dimension } n-1 \tag{8.163}
\end{equation*}
$$

and

$$
\begin{equation*}
\pi_{H}\left(\Lambda_{\tilde{f}} \cap T_{H}^{*} X\right)=\Lambda_{f} \tag{8.164}
\end{equation*}
$$

The first follows from the fact that $\Lambda_{\tilde{f}}$ is a graph over $X$ and the second from the definition, ( $\overline{8.161}$ ). So we find
19.42 [19.34MMA ${ }^{2} 10_{35}$ If $z \in \Sigma(P)$ and $H$ is a hypersurface through $\pi(z)$ satisfying (8.156) and (노.157) then $\pi_{36}^{P}:\left(\Sigma(P) \cap T_{H}^{*} X\right) \longrightarrow T^{*} H$ is a local diffeomorphism in a neighbourhood $z$; if ( $\mathbf{8 . 1 5 8}$ ) is to hold then

$$
\begin{equation*}
\Lambda_{\tilde{f}} \cap T_{H}^{*} X=\left(\pi_{H}^{P}\right)^{-1}\left(\Lambda_{f}\right) \text { near } z \tag{8.165}
\end{equation*}
$$

Proof. We know that $H_{p}$ is tangent to $\Sigma(P)$ but, by assumption $\left(\frac{19.36}{8.158}\right)$ is not tangent to $T_{H}^{*} X$ at $z$. Then $\Sigma(P) \cap T_{H}^{*} X$ does have dimension $2 n-1-1=2(n-1)$. Moreover $\pi_{H}$ is projection along $\partial_{\tau}$ which cannot be tangent to $\Sigma(P) \cap T_{H}^{*} X$ (since it would be tangent to $\Sigma(P)$ ). Thus $\pi_{H}^{P}$ has injective differential, hence is a local isomorphism.

So this is our strategy:
Start with $f \in \mathcal{C}^{\infty}(H)$, look at $\Lambda_{f} \subset T^{*} H$, lift to $\Lambda \cap T_{H}^{*} X \subset \Sigma(P)$ by $\pi_{H}^{P}$. Now let

$$
\begin{equation*}
\Lambda=\bigcup\left\{H_{p}-\text { curves through }\left(\pi_{H}^{P}\right)^{-1}\left(\Lambda_{f}\right)\right\} \tag{8.166}
\end{equation*}
$$

This we will show to be Lagrangian and of the form $\Lambda_{\tilde{f}}$, it follows that

$$
\begin{equation*}
p(x, d \tilde{f})=0, \tilde{f} \upharpoonright H=f \tag{8.167}
\end{equation*}
$$

### 8.2. Riemann metrics and quantization

Metrics, geodesic flow, Riemannian normal form, Riemann-Weyl quantization.

### 8.3. Transport equation

The first thing we need to do is to finish the construction of characteristic hypersurfaces using Hamilton-Jacobi theory, i.e. prove Theorem XIX.37. We have already defined the submanifold $\Lambda$ as follows:

1) We choose $z \in \Sigma(P)$ and $t \in \mathcal{C}^{\infty}(X)$ s.t. $H_{p} \pi^{*}(t) \neq 0$ at $d z$, then selected $f \in \mathcal{C}^{\infty}(H), H=\{t=0\} \cap \Omega, \Omega \ni \pi z$ s.t.

$$
\begin{equation*}
z(v)=d f(v) \forall v \in T_{\pi z} H \tag{8.168}
\end{equation*}
$$

Then we consider

$$
20.2
$$

$$
\begin{equation*}
\Lambda_{f}=\operatorname{graph}\{d f\}=\{(x, d f(x)), x \in H\} \subset T^{*} H \tag{8.169}
\end{equation*}
$$

as our "initial data" for $\Lambda$. To move it into $\Sigma(P)$ we noted that the map
20.3 (8.170)

$$
\Sigma(P) \cap \underset{\substack{\| \\\left\{t=0 \text { in } T^{*} X\right\}}}{T_{H}^{*} X} \longrightarrow T^{*} H
$$

 $\left(\frac{20.3}{8.170)}\right.$ is therefore a submanifold $\tilde{\Lambda}_{f} \subset \Sigma(p) \cap T_{H}^{*} X$ of dimension $\operatorname{dim} X-1=$ $\operatorname{dim} H$. We define

$$
\begin{equation*}
\Lambda=\bigcup\left\{H_{p}-\text { curves of length } \epsilon \text { starting on } \tilde{\Lambda}_{f}\right\} \tag{8.171}
\end{equation*}
$$

So we already know:

$$
\begin{equation*}
\Lambda \subset \Sigma(P) \text { is a manifold of dimension } n \tag{8.172}
\end{equation*}
$$

and
(8.173) $\pi: \Lambda \longrightarrow X$ is a local diffeomorphism near $n$,

What we need to know most of all is that
$\qquad$ $\Lambda$ is Lagrangian.
That is, we need to show that the symplectic two form vanishes identically on $T_{z^{\prime}} \Lambda, \forall z^{\prime} \in \Lambda$ (at least near $\left.z\right)$. First we check this at $z$ itself! Now

$$
\begin{equation*}
T_{z} \Lambda=T_{z} \tilde{\Lambda}_{f}+\operatorname{sp}\left(H_{p}\right) \tag{tabular}
\end{equation*}
$$

Suppose $v \in T_{z} \tilde{\Lambda}_{f}$, then

$$
\begin{equation*}
\omega\left(v, H_{p}\right)=-d p(v)=0 \text { since } \tilde{\Lambda}_{f} \subset \Sigma(P) \tag{8.176}
\end{equation*}
$$

Of course $\omega\left(H_{p}, H_{p}\right)=0$ so it is enough to consider
20.10

$$
\begin{equation*}
\omega \mid\left(T_{z} \tilde{\Lambda}_{f} \times T_{z} \tilde{\Lambda}_{f}\right) \tag{8.177}
\end{equation*}
$$

Recall from our discussion of the projection $\left(\frac{20.3}{8.170} 0\right)$ that we can write it as projection along $\partial_{\tau}$. Thus

$$
\begin{gather*}
\omega_{X}(v, w)=\omega_{H}\left(v^{\prime}, w^{\prime}\right) \text { if } v, w \in T_{z}\left(T_{H} X\right) \\
\left(c_{H}^{*}\right)_{*} v=v^{\prime}\left(c_{H}^{*}\right)_{*} w=w^{\prime} \in T_{z}\left(T^{*} H\right) \tag{8.178}
\end{gather*}
$$

where $z=d f(\pi(z))$. Thus the form $\left(\frac{20.10}{8.177}\right)$ vanishes identically because $\Lambda_{f}$ is Lagrangian.

In fact the same argument applies at every point of the initial surface $\tilde{\Lambda}_{f} \subset \Lambda$ :
20.12

$$
\begin{equation*}
T_{z^{\prime}} \Lambda \text { is Lagrangian } \forall z^{\prime} \in \tilde{\Lambda}_{f} \tag{8.179}
\end{equation*}
$$

To extend this result out into $\Lambda$ we need to use a little more differential geometry. Consider the local diffeomorphisms obtained by exponentiating $H_{p}$ :

$$
\begin{equation*}
\exp \left(\epsilon H_{p}\right)(\Lambda \cap \Omega) \subset \Lambda \forall \epsilon \text { small, } \Omega \ni z \text { small. } \tag{8.180}
\end{equation*}
$$

This indeed is really the definition of $\Lambda_{j}$ more precisely,

$$
\begin{equation*}
\Lambda=\bigcup_{\epsilon \text { small }} \exp \left(\epsilon H_{p}\right)\left(\tilde{\Lambda}_{f}\right) . \tag{8.181}
\end{equation*}
$$

The main thing to observe is that, on $T^{*} H$, the local diffeomorphisms $\exp \left(\epsilon H_{p}\right)$ are symplectic:

```
exp}(\epsilon\mp@subsup{H}{p}{}\mp@subsup{)}{}{*}\mp@subsup{\omega}{X}{}=\mp@subsup{\omega}{X}{}
```

Clearly $\left(\frac{20.15}{8.182}\right),\left(\frac{20.13}{8.180}\right)$ and $\left(\frac{20.12}{8.179}\right)$ prove $\left(\frac{20.7}{8.174)}\right.$. The most elegant wary to prove ( 8.182 ) is to use Cartan's identity (valid for $H_{p}$ any vector field, $\omega$ any form)

$$
\begin{equation*}
\frac{d}{d \epsilon} \exp \left(\epsilon H_{p}\right)^{*} \omega=\exp \left(\epsilon H_{p}\right)^{*}\left(\mathcal{L}_{H_{p}} \omega\right) \tag{8.183}
\end{equation*}
$$

$$
\mathcal{L}_{V}=d \circ \iota_{V}+\iota_{V} \circ d,
$$

$c_{V}$ being contradiction with $V$ (i.e. $\left.\alpha(\cdot, \cdot, \ldots) \longrightarrow \alpha(V, \cdot, \cdot, \ldots)\right)$. Thus

$$
\begin{equation*}
\mathcal{L}_{H_{p}} \omega=d\left(\omega\left(H_{p}, \cdot\right)\right)+\iota_{V}(\underset{\|}{\|} \underset{0}{d \omega})=d(d p)=0 . \tag{8.185}
\end{equation*}
$$

Thus from $\left(\frac{20.5}{8.17} 2\right),\left(\frac{20.6}{8.17} 3\right)$ and $\left(\frac{20.7}{8.17} 4\right)$ we know that

$$
\begin{equation*}
\Lambda=\operatorname{graph}(d \tilde{f}), \tilde{f} \in \mathcal{C}^{\infty}(X), \text { near } \pi(z) \tag{8.186}
\end{equation*}
$$

must satisfy the eikonal equation
20.20

$$
\begin{equation*}
p(x, d \tilde{f}(x))=0 \text { near } \pi(z), H \tilde{f} \upharpoonright H=f \tag{8.187}
\end{equation*}
$$

where we may actually have to add a constant to $\tilde{f}$ to get the initial condition since we only have $d \tilde{f}=d f$ on $T H$.

So now we can return to the construction of travelling waves: We want to find

$$
\begin{equation*}
u \in I^{*}\left(X, G ; \Omega^{\frac{1}{2}}\right) \quad G=\{f=0\} \tag{8.188}
\end{equation*}
$$

such that $u$ is elliptic at $z \in \Sigma(p)$ and

$$
\begin{equation*}
P u \in \mathcal{C}^{\infty}(X) . \tag{8.189}
\end{equation*}
$$

So far we have noticed that
20.23

$$
\begin{equation*}
\sigma_{m+M}(P u)=\sigma_{m}(P) \upharpoonright N^{*} G \cdot \sigma(u) \tag{8.190}
\end{equation*}
$$

so that

$$
\begin{equation*}
N^{*} G \subset \Sigma(p) \Longleftrightarrow p(x, d f)=0 \text { on } f=0 \tag{8.191}
\end{equation*}
$$

implies

$$
\begin{equation*}
P u \in I^{m+M-1}\left(X, G ; \Omega^{\frac{1}{2}}\right) \text { near } \pi(z) \tag{8.192}
\end{equation*}
$$

which is one order smoother than without (80.24 $)$.
It is now clear, I hope, that we need to make the "next symbol" vanish as well, i.e. we want
20.26

$$
\begin{equation*}
\sigma_{m+M-1}(P u)=0 \tag{8.193}
\end{equation*}
$$

Of course to arrange this it helps to know what the symbol is!
Proposition 8.2. Suppose $P \in \Psi^{m}\left(X ; \Omega^{\frac{1}{2}}\right)$ and $G \subset X$ is a $\mathcal{C}^{\infty}$ hypersurface characteristic for $P$ (i.e. $\left.N^{*} G \subset \Sigma(P)\right)$ then $\forall u \in I^{M}\left(X, G ; \Omega^{\frac{1}{2}}\right)$

$$
\begin{equation*}
\sigma_{m+M-1}(P u)=\left(-i H_{p}+a\right) \sigma_{m}(u) \tag{8.194}
\end{equation*}
$$

where $a \in S^{m-1}\left(N^{*} G\right)$ and $H_{p}$ is the Hamilton vector field of $p=\sigma_{m}(P)$.
Proof. Observe first that the formula makes sense since $\Lambda=N^{*} G$ is Lagrangian, $\Lambda \subset \Sigma(p)$ implies $H_{p}$ is tangent to $\Lambda$ and if $p$ is homogeneous of degree $m$ (which we are implicitly assuming) then
20.29

$$
\begin{equation*}
\mathcal{L}_{H_{p}}: S^{r}\left(\Lambda ; \Omega^{\frac{1}{2}}\right) \longrightarrow S^{r+m-1}\left(\Lambda ; \Omega^{\frac{1}{2}}\right) \forall m \tag{8.195}
\end{equation*}
$$

where one can ignore the half-density terms. So suppose $G=\left\{x_{1}=0\right\}$ locally, which we can always arrange by choice of coordinates. Then

$$
\begin{equation*}
X=N^{*} G=\left\{\left(0, x^{\prime}, \xi_{1}, 0\right) \in T^{*} X\right\} \tag{8.196}
\end{equation*}
$$

To say $N^{*} G \subset \Sigma(p)$ means $p=0$ on $\Lambda$, i.e.

$$
\begin{equation*}
p=x_{1} q(x, \xi)+\sum_{j>1} \xi_{j} p_{j}(x, \xi) \text { near } z \tag{tabular}
\end{equation*}
$$

with $q$ homogeneous of degree $m$ and the $p_{j}$ homogeneous of degree $m-1$. Working microlocally we can choose $Q \in \Psi^{m}\left(X, \Omega^{\frac{1}{2}}\right), P_{j} \in \Psi^{m-1}\left(X, \Omega^{\frac{1}{2}}\right)$ with

$$
\begin{equation*}
\sigma_{m}(Q)=q, \sigma_{m-1}\left(P_{j}\right)=p_{j} \text { near } z \tag{8.198}
\end{equation*}
$$

Then, from $\left(\frac{20.30}{18.197}\right)$

$$
\begin{equation*}
P=x_{1} Q+D_{x_{j}} P_{j}+R+P^{\prime}, \quad R \in \Psi^{m-1}\left(X ; \Omega^{\frac{1}{2}}\right) z \notin W F^{\prime}\left(P^{\prime}\right), P^{\prime} \in \Psi^{m}\left(X, \Omega^{\frac{1}{2}}\right) \tag{8.199}
\end{equation*}
$$

Of course $P^{\prime}$ does not affect the symbol near $z$ so we only need observe that

$$
\begin{align*}
& \sigma_{r-1}(x, u)=-d_{\xi_{1}} \sigma_{r}(u) \\
& \quad \forall u \in I^{r}\left(X, G ; \Omega^{\frac{1}{2}}\right)  \tag{8.200}\\
& \sigma_{r}\left(D_{x_{j}} u\right)=D_{x_{j}} \sigma_{r}(u)
\end{align*}
$$

This follows from the local expression

$$
\begin{equation*}
u(x)=(2 \pi)^{-1} \int e^{i x_{1} \xi_{1}} a\left(x^{\prime}, \xi_{1}\right) d \xi_{1} \tag{8.201}
\end{equation*}
$$

Then from $\left(\frac{20.32}{8.199)}\right.$ we get

$$
\begin{align*}
& \sigma_{m+M-1}(P u)=-D_{\xi_{1}}\left(q \sigma_{M}(u)\right)+\sum_{j} D_{x_{j}}\left(p_{j} \sigma_{M}(u)\right)+r \cdot \sigma_{m}(u) \\
& \quad=-i\left(\sum_{j>1} p_{j} \upharpoonright \Lambda \frac{\partial}{\partial x_{j}}-q \upharpoonright \Lambda \frac{\partial}{\partial \xi_{i}}\right) \sigma_{M}(u)+a^{\prime} \sigma_{M}(u) \tag{8.202}
\end{align*}
$$

Observe from $\left(\frac{20.30}{8.197}\right)$ that the Hamilto $\eta_{28.28}$. expression in parenthesis. This proves (8.194).

So, now we can solve ( ${ }^{20.26}$ 8.193). We just set

$$
\begin{equation*}
\sigma_{M}(u)\left(\exp \left(\epsilon H_{p}\right) z^{\prime}\right)=e^{i \epsilon A} \exp \left(\epsilon H_{p}\right)^{*}[b] \forall z^{\prime} \in \tilde{\Lambda}_{f}=\Lambda \cap\{t=0\} \tag{8.203}
\end{equation*}
$$

where $A$ is the solution of

$$
\begin{equation*}
H_{p} A=a, A \upharpoonright t=0=0 \quad \text { on } \Lambda_{0} \tag{8.204}
\end{equation*}
$$

and $b \in S^{r}\left(\Lambda_{0}\right)$ is a symbol defined on $\Lambda_{0}=\Lambda \cap\{t=0\}$ near $z$.
Proposition 8.3. Suppose $P \in \Psi^{m}\left(X ; \Omega^{\frac{1}{2}}\right)$ has homogeneous principal symbol of degree $m$ satisfying

$$
\begin{equation*}
p=\sigma_{m}(P) \text { is real } \tag{8.205}
\end{equation*}
$$

and $z \in \Sigma(p)$ is fixed. Then if $H \ni \pi(z)$ is a hypersurface such that $\pi_{*}\left(H_{p}\right) \cap H$ and $G \subset H$ is an hypersurface in $H$ s.t.

$$
\begin{equation*}
\bar{z}=c_{H}^{*}(z) \in H_{\pi z}^{*} G \tag{8.207}
\end{equation*}
$$

there exist a characteristic hypersurface $\tilde{G} \subset X$ for $P$ such that $\tilde{G} \cap H=G$ near $\pi(z), z \in N_{\pi z}^{*} \tilde{G}$. For each

$$
\begin{equation*}
u_{0} \in I^{m+\frac{1}{4}}\left(H, G ; \Omega^{\frac{1}{2}}\right) \text { with } W F\left(u_{0}\right) \subset \gamma, \tag{8.208}
\end{equation*}
$$

$\gamma$ a fixed small conic neighbourhood of $\bar{z} n T^{*} H$ there exists

$$
\begin{equation*}
u \in I\left(X, \tilde{G} ; \Omega^{\frac{1}{2}}\right) \text { satisfying } \tag{8.209}
\end{equation*}
$$

$$
\begin{equation*}
u \upharpoonright G=u_{0} \text { near } \pi z \in H \tag{8.210}
\end{equation*}
$$

$$
\begin{equation*}
P u \in \mathcal{C}^{\infty} \text { near } \pi z \in X \tag{8.211}
\end{equation*}
$$

Proof. All the stuff about $G$ and $\tilde{G}{ }^{2}$ is ${ }_{3}$ just Hamilton-Jacobi theory. We can take the symbol of $u_{0}$ to be the $\frac{b}{}$ in $(\overline{8.203})_{44}$ once we think a little about halfdensities, and thereby expect ( 8.210 ) and (8.211) to hold, modulo certain singularities. Indeed, we would get

$$
\begin{gather*}
u_{1} \upharpoonright G-u_{0} \in I^{r+\frac{1}{4}-1}\left(H, G ; \Omega^{\frac{1}{2}}\right) \text { near } \pi z \in H  \tag{8.212}\\
P u \in I^{r+m-2}\left(X, \tilde{G} ; \Omega^{\frac{1}{2}}\right) \text { near } \pi z \in X .
\end{gather*}
$$

So we have to work a little to remove lower order terms. Let me do this informally, without worrying too much about ( 8.210 ) for a moment. In fact I will put ( 8.212 ) into the exercises!

All we really have to observe to improve $\left(\frac{20.46}{8.213}\right)$ to $\left(\frac{20.44}{8.211}\right)$ is that

$$
g \in I^{r}\left(X, \tilde{G} ; \Omega^{\frac{1}{2}}\right) \Longrightarrow \exists \quad u \in I^{r+m-1}\left(X ; \tilde{G} ; \Omega^{\frac{1}{2}}\right)
$$

$$
\begin{equation*}
\text { s.t. } \quad P u-g \in I^{r-1}\left(X, \tilde{G} ; \Omega^{\frac{1}{2}}\right) \tag{8.214}
\end{equation*}
$$

which we can then iterate and asymptoticall. $y_{7}$ sum. In fact we can choose the solution so $u \upharpoonright H \in \mathcal{C}^{\infty}$, near $\pi \bar{z}$. To solve ( $\overline{8.214}$ ) we just have to be able to solve

$$
\begin{equation*}
-i\left(H_{p}+a\right) \sigma(u)=\sigma(g) \tag{8.215}
\end{equation*}
$$

which we can do by integration (duHamel's principle).

The equation $\left(\frac{20.48}{8.215}\right)$ for the symbol of the solution is the transport equation. We shall use this construction next time to produce a microlocal parametrix for $P$ !

### 8.4. Problems

Problem 8.2. Let $X$ be a $\mathcal{C}^{\infty}$ manifold, $G \subset X$ on $\mathcal{C}^{\infty}$ hypersurface and $t \in \mathcal{C}^{\infty}(X)$ a real-valued function such that
$\mathrm{T}(8.216) \quad d t \neq 0$ on $T_{p} G \forall p \in L=G \cap\{t=0\}$.
Show that the transversality condition $\left(\frac{\mathbb{T}}{\mathbf{1}} .216\right)$ ensures that $H=\{t=0\}$ and $L=H \cap G$ are both $\mathcal{C}^{\infty}$ submanifolds.
20.50 Problem 8.3. Assuming ( $\left(\frac{T}{8} .216\right)$ show that $d t$ gives an isomorphism of line bundles

$$
\begin{equation*}
\Omega^{\frac{1}{2}}(H) \equiv \Omega_{H}^{\frac{1}{2}}(X) \sim \Omega_{H}^{\frac{1}{2}}(X) /|d t|^{\frac{1}{2}} \tag{8.217}
\end{equation*}
$$

and hence one can define a restriction map,

$$
\begin{equation*}
\mathcal{C}^{\infty}\left(X ; \Omega^{\frac{1}{2}}\right) \longrightarrow \mathcal{C}^{\infty}\left(H ; \Omega^{\frac{1}{2}}\right) \tag{8.218}
\end{equation*}
$$

20.51 Problem 8.4. Assuming 1 and 2, make sense of the restriction formula

$$
\begin{equation*}
\upharpoonright H: I^{m}\left(X, G ; \Omega^{\frac{1}{2}}\right) \longrightarrow I^{m+\frac{1}{4}}\left(H, L ; \Omega^{\frac{1}{2}}\right) \tag{8.219}
\end{equation*}
$$

and prove it, and the corresponding symbolic formula

$$
\begin{equation*}
\sigma_{m+\frac{1}{4}}(u \upharpoonright H)=\left(\iota_{H}^{*}\right)^{*}\left(\sigma_{m}(u) \upharpoonright N_{L}^{*} G\right) /|d \tau|^{\frac{1}{2}} \tag{8.220}
\end{equation*}
$$

$N B$. Start from local coordinates and try to understand restriction at that level before going after the symbol formula!

### 8.5. The wave equation

We shall use the construction of travelling wave solutions to produce a parametrix, and then a fundamental solution, for the wave equation. Suppose $X$ is a Riemannian manifold, e.g. $\mathbb{R}^{n}$ with a 'scattering' metrice:
21.1

$$
\begin{equation*}
g=\sum_{i, j=1}^{n} g_{i j}(x) d x^{i} d x^{j}, g_{i j}=\delta_{i j}|x| R \tag{8.221}
\end{equation*}
$$

Then the associates Laplacian, on functions, i.e.

## 21.2

$$
\begin{equation*}
\Delta u=-\sum_{i, j=1}^{n} \frac{1}{\sqrt{g}} \frac{\partial}{\partial x_{j}}\left(\delta g g^{i j}(x)\right) \frac{\partial}{\partial x_{i}} u \tag{8.222}
\end{equation*}
$$

where $g^{i j}(x)=\left(g_{i j}(x)\right)^{-1}$ and $g=\operatorname{det} g_{i j}$. We are interested in the wave equation

$$
\begin{equation*}
P u=\left(D_{t}^{2}-\Delta\right) u=f \quad \text { on } \mathbb{R} \times X \tag{8.223}
\end{equation*}
$$

For simplicity we assume $X$ is either compact, or $X=\mathbb{R}^{n}$ with a metric of the form ( 8.221 ).

The cotangent bundle of $\mathbb{R} \times X$ is

$$
\begin{equation*}
T^{*}(\mathbb{R} \times X) \simeq T^{*} \mathbb{R} \times T^{*} X \tag{8.224}
\end{equation*}
$$

with canonical coordinates $(t, x, \tau, \xi)$. In terms of this

$$
\begin{equation*}
\sigma(P)=\tau^{2}-|\xi|^{2}|\xi|=\sum_{i, j=1}^{n} g^{i j}(x) \xi_{i} \xi_{j} \tag{8.225}
\end{equation*}
$$

Thus we certainly have an operator satisfying the conditions of $\left(\frac{21.3}{8.223} 3\right)$ and $\left(\frac{21.4}{8.225} 5\right)$, since

$$
\begin{equation*}
d_{\text {fibre }} p=\left(\frac{\partial p}{\partial \tau}, \frac{\partial p}{\partial \xi}\right)=0 \Longrightarrow \tau=0 \text { and } g^{i j}(x) \xi_{i}=0 \Longrightarrow \xi=0 \tag{8.226}
\end{equation*}
$$

As initial surface we consider the obvious hypersurface $\{t=0\}$ (although it will be convenient to consider others). We are after the two theorems, one local and global, the other microlocal, although made to look global.

Theorem 8.3. If $X$ is a Riemannian manifold, as above, then for every $f \in$ $\mathcal{C}_{c}^{-\infty}(\mathbb{R} \times X) \quad \exists!u \in \mathcal{C}^{-\infty}(\mathbb{R} \times X)$ satisfying

$$
\begin{equation*}
P u=f, u=0 \text { in } t<\inf \{\bar{t} ; \quad \exists(\bar{t}, x) \in \operatorname{supp}(f)\} . \tag{8.227}
\end{equation*}
$$

THEOREM 8.4. If $X$ is a Riemannian manifold, as above, then for every $u \in$ $\mathcal{C}^{-\infty}(\mathbb{R} \times X)$,

$$
\begin{equation*}
W F(u) \backslash W F(P u) \subset \Sigma(P) \backslash W F(P u) \tag{8.228}
\end{equation*}
$$

is a union of maximally extended $H_{o}$-curves in the open subset $\Sigma(P) \backslash W F(P u)$ of $\Sigma(P)$.

Let us think about Theorem $\frac{21.5}{8.3}$ first. Suppose $\bar{x} X$ is fixed on $\delta_{\bar{x}} \in \mathcal{C}^{-\infty}(X ; \Omega)$ is the Dirac delta ( $g$ measure) at $\bar{x}$. Ignoring, for a moment, the fact that this is not quite a generalized function we can look for the "forward fundamental solution" of $P$ with pole at $(0, \bar{x})$ :

$$
\begin{gather*}
P E_{\bar{x}}(t, x)=\delta(t) \delta_{\bar{x}}(x) \\
E_{\bar{x}}=0 \text { in } t<0 . \tag{8.229}
\end{gather*}
$$

Theorem $\frac{21.5}{8.3}$ asserts its existence and uniqueness. Conversely if we can construct $E_{\bar{x}}$ for each $\bar{x}$, and get reasonable dependence on $\bar{x}$ (continuity is almost certain once we prove uniqueness) then

$$
\begin{equation*}
K(t, x ; \bar{t}, \bar{x})=E_{\bar{x}}(t-\bar{t}, x) \tag{8.230}
\end{equation*}
$$

is the kernel of the operator ${ }_{12} f . \leftrightarrows u$ solving $\left(\frac{21.6}{8.227} 7\right)$.
So, we want to solve ( 8.229 ). First we convert it (without worrying about rigour) to an initial value problem. Namely, suppose we can solve instead

$$
\begin{gather*}
P G_{\bar{x}}(t, x)=0 \text { in } \mathbb{R} \times X \\
G_{\bar{x}}(0, x)=0, D_{t} G_{\bar{x}}(0, x)=\delta_{\bar{x}}(x) \text { in } X . \tag{8.231}
\end{gather*}
$$

Note that
21.12

$$
\begin{equation*}
(g(t, x, \tau, 0) \notin \Sigma(P) \Longrightarrow(t, x ; \tau, 0) \notin W F(G) \tag{8.232}
\end{equation*}
$$

This means the restriction maps, to $t=0$, in (21.231) are well-defined. In fact so is the product map:

$$
\begin{equation*}
E_{\bar{x}}(t, x)=H(t) G_{\bar{x}}(t, x) \tag{8.233}
\end{equation*}
$$

Then if $G$ satisfied $\left(\frac{21.11}{8.231}\right)$ a simple computation shows that $E_{\bar{x}}$ satisfies $\left(\frac{21.9}{8.229} 9\right)$. Thus we want to solve ( 8.231 ).

Now ( 8.231 ) seems very promising. The initial data, $\delta_{\bar{x}}$, is certainly conormal to the point $\{\bar{x}\}$, so we might try to use our construction of travelling wave solutions. However there is a serious problem. We already noted that, for the wave equation,
there cannot be any smooth characteristic surface other than a hypersurface. The point is that if $H$ has codimension $k$ then

$$
\begin{equation*}
N_{\bar{x}}^{*} H \subset T_{\bar{x}}^{*}(\mathbb{R} \times X) \text { has dimension } k \tag{8.234}
\end{equation*}
$$

To be characteristic we must have

$$
\begin{equation*}
N_{\bar{x}}^{*} H \subset \Sigma(P) \Longrightarrow k=1 \tag{8.235}
\end{equation*}
$$

Since the only linear space contained in a (proper) cone is a line.
However we can easily 'guess' what the characteristic surface corresponding to the point $(x, \bar{x})$ is - it is the cone through that point:

This certainly takes us beyond our conormal theory. Fortunately there is a way around the problem, namely the possibility of superposition of conormal solutions.

To see where this comes from consider the representation in terms of the Fourier transform:

$$
\begin{equation*}
\delta(x)=(2 \pi)^{-n} \int e^{i x \xi} d \xi \tag{8.236}
\end{equation*}
$$

The integral of course is not quite a proper one! However introduce polar coordinates $\xi=r \omega$ to get, at least formally

$$
\begin{equation*}
\delta(x)=(2 \pi)^{-n} \int_{0}^{\infty} \int_{\mathbb{S}^{n-1}} e^{i r x \cdot \omega} r^{n-1} d r d \omega \tag{8.237}
\end{equation*}
$$

In odd dimensions $r^{n-1}$ is even so we can write

$$
\begin{equation*}
\delta(x)=\frac{1}{2(2 \pi)^{n}} \int_{\mathbb{S}^{n-1}} \int_{-\infty}^{\infty} e^{i r x \cdot \omega} r^{n-1} d r d \omega, n \text { odd } \tag{8.238}
\end{equation*}
$$

Now we can interpret the $r$ integral as a 1-dimensional inverse Fourier transform so that, always formally,

$$
\begin{gather*}
\delta(x)=\frac{1}{2(2 \pi)^{n-1}} \int_{\mathbb{S}^{n-1}} f_{n}(x \cdot \omega) d \omega \\
\quad f_{n}(s)=\frac{1}{(2 \pi)} \int e^{i r s} \gamma^{n-1} d r \tag{8.239}
\end{gather*}
$$

In even dimensions we get the same formula with

$$
\begin{equation*}
f_{n}(s)=\frac{1}{2 \pi} \int e^{i r s}|r|^{n-1} d r \tag{8.240}
\end{equation*}
$$

These formulas show that
21.20

$$
\begin{equation*}
f_{n}(s)=\left|D_{s}\right|^{n-1} \delta(s) \tag{8.241}
\end{equation*}
$$

Here $\left|S_{s}\right|^{n-1}$ is a pseudodifferential operator for $n$ even or differential operator $\left(=D_{s}^{n-1}\right)$ if $n$ is odd. In any case

$$
\begin{equation*}
f_{n} \in I^{n-1+\frac{1}{4}}(\mathbb{R},\{0\})! \tag{8.242}
\end{equation*}
$$

Now consider the map
21.22

$$
\begin{equation*}
\mathbb{R}^{n} \times \mathbb{S}^{n-1} \ni(x, \omega) \mapsto x \cdot \omega \in \mathbb{R} \tag{8.243}
\end{equation*}
$$

Thus $\mathcal{C}^{\infty}$ has different

$$
\begin{equation*}
\omega \cdot d x+x \cdot d \omega \neq 0 \text { or } x \cdot \omega=0 \tag{8.244}
\end{equation*}
$$

So the inverse image of $\{0\}$ is a smooth hypersurface $R$.

$$
\begin{equation*}
f_{n}(x, \omega)=\frac{1}{2 \pi} \int e^{i(x \cdot \omega) r}|r|^{n-1} d r \in I^{\frac{n}{4}-\frac{1}{4}}\left(\mathbb{R} \times \mathbb{S}^{n-1}, R\right) \tag{8.245}
\end{equation*}
$$

Proof. Replacing $|r|^{n-1}$ by $\rho(r)|r|^{n-1}+(1-\rho(r))|r|^{n-1}$, where $\rho(r)=0 \mathrm{n}$ $r<\frac{1}{2}, \rho(r)=1$ in $r>1$, expresses $f_{n}$ as a sum of a $\mathcal{C}^{\infty}$ term and a conormal distribution. Check the order yourself!

Proposition 8.4. (Radon inversion formula). Under pushforward corresponding to $\mathbb{R}^{n} \times \mathbb{S}^{n-1} @>\pi_{1} \gg \mathbb{R}^{n}$

$$
\begin{align*}
\left(\pi_{1}\right)_{*} f_{n}^{\prime} & =2(2 \pi)^{n-1} \delta(x)  \tag{8.246}\\
f_{n}^{\prime} & =f_{n}|d \omega||d x|
\end{align*}
$$

Proof. Pair with a test function $\phi \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ :

$$
\begin{equation*}
\left(\pi_{1}\right)_{*} f_{n}^{\prime}=\iint f_{n}(x \cdot \omega) \phi(x) d x d \omega \tag{8.247}
\end{equation*}
$$

by the Fourier inversion formula.
So now we have a superposition formula expressing $\delta(x)$ as an integral:

$$
\begin{equation*}
\delta(x)=\frac{1}{2(2 \pi)^{n-1}} \int_{\mathbb{S}^{n-1}} f_{n}(x \cdot \omega) d \omega \tag{8.248}
\end{equation*}
$$

where for each fixed $f_{1}(x \cdot \omega)$ is conormal with respect to $x \cdot \omega=0$. This gives us a strategy to solve (8.231).

Proposition 8.5. Each $\bar{x} \in X$ has a neighbourhood, $U_{\bar{x}}$, such that for $\bar{t}>0$ (independent of $\bar{x}$ ) there are two characteristic hypersurfaces for each $\omega \in \mathbb{S}^{n-1}$

$$
\begin{equation*}
H_{\bar{x}, \omega)}^{ \pm} \subset(-\bar{t}, \bar{t}) \times U_{\bar{x}} \tag{8.249}
\end{equation*}
$$

depending on $\bar{x}, \omega$, and there exists

$$
\begin{equation*}
u^{ \pm}(t, x ; \bar{x}, \omega) \in I^{*}\left(\left(-\bar{t} \bar{t} \mid \times U_{\bar{x}}, H_{(\bar{x}, \omega)}^{ \pm}\right)\right. \tag{8.250}
\end{equation*}
$$

such that

$$
\begin{gather*}
P u^{ \pm} \in \mathcal{C}^{\infty}  \tag{8.251}\\
\begin{cases}u^{+}+\bar{u} \upharpoonright t=0=\delta_{\bar{x}}(x \cdot \omega) & \text { in } U_{\bar{x}} \\
D_{t}\left(u^{+}+u^{-}\right) \upharpoonright\{t=0\}=0 & \text { in } U_{\bar{x}}\end{cases} \tag{8.252}
\end{gather*}
$$

Proof. The characteristic surfaces are constructed through Hamilton-Jacobi theory:

$$
\begin{gather*}
N^{*} H^{ \pm} \subset \Sigma(P) \\
H_{0}=H^{ \pm} \cap\{t=0\}=\{x \cdot \omega=0\} \tag{8.253}
\end{gather*}
$$

There are two or three because the conormal direction to $H_{0}$ at $0 ; \omega d x$, has two $\Sigma(P):$

$$
\begin{equation*}
\tau= \pm 1, \quad(\tau, \omega) \in T_{0}^{*}(\mathbb{R} \times X) \tag{8.254}
\end{equation*}
$$

With each of these two surfaces we can associate a microlocally unique conormal solution
21.32

$$
\begin{gather*}
P u^{ \pm}=0, \quad u^{ \pm} \upharpoonright\{t=0\}=u_{0}^{ \pm} \\
u_{0}^{ \pm} \in I^{*}\left(\mathbb{R}^{n},\{x \cdot \omega=0\}\right) \tag{8.255}
\end{gather*}
$$

Now, it is easy to see that there are unique choices
21.33

$$
\left.\begin{array}{c}
u_{\delta}^{+}+u_{0}^{-}=\delta(x \cdot \omega) \\
D_{t} u^{+}+D_{t} u^{-} \tag{8.256}
\end{array}\right)\{t=0\}=0 . ~ \$
$$

(See exercise 2.) This solves $\left(\frac{21.30}{8.252}\right)$ and proves the proposition (modulo a fair bit of hard work!).

So now we can use the superposition principle. Actually it is better to add the variables $\omega$ to the problem and see that

$$
\begin{align*}
u^{ \pm}(t, x ; \omega, \bar{x}) & \in I^{*}\left(\mathbb{R} \times \mathbb{R}^{n} \times \mathbb{S}^{n-1} \times \mathbb{R}^{n} ; H^{ \pm}\right) \\
H^{ \pm} & \subset \mathbb{R} \times \mathbb{R}^{n} \times \mathbb{S}^{n-1} \times \mathbb{R}^{n} \tag{8.257}
\end{align*}
$$

being fixed by the condition that

$$
\begin{equation*}
H^{ \pm} \cap \mathbb{R} \times \mathbb{R}^{n} \times\{\omega\} \times\{\bar{x}\}=H_{\bar{x}, \omega}^{ \pm} \tag{8.258}
\end{equation*}
$$

Then we set
21.35

$$
\begin{equation*}
G_{\bar{x}}^{\prime}(t, x)=\int_{\mathbb{S}^{n-1}}\left(u^{+}+u^{-}\right)(x, x ; \omega, \bar{x}) \tag{8.259}
\end{equation*}
$$

This satisfies $\left(\frac{21.11}{8.231}\right)$ locally near $\bar{x}$ and modulo $\mathcal{C}^{\infty}$. i.e.

$$
\left\{\begin{array}{l}
P G_{\bar{x}}^{\prime} \in \mathcal{C}^{\infty}\left((-\bar{t}(\bar{t})) \times U_{\bar{x}}\right)  \tag{8.260}\\
G_{\bar{x}}^{\prime} \upharpoonright\{t=0\}=x v, \\
\\
D_{t} G_{\bar{x}}^{\prime}=\delta_{\bar{x}}(x)+v_{2}
\end{array} \quad v_{i} \in \mathcal{C}^{\infty}\right.
$$

Let us finish off by doing a calculation. We have (more or less) shown that $u^{ \pm}$are conormal with respect to the hypersurfaces $H^{ \pm}$. A serious question then is, what is (a bound one) the wavefront set of $G_{\bar{x}}^{\prime}$ ? This is fairly easy provided we understand the geometry. First, since $u^{ \pm}$are conormal,

$$
\begin{equation*}
W F\left(u^{ \pm}\right) \subset N^{*} H^{ \pm} \tag{8.261}
\end{equation*}
$$

Then the push-forward theorem says
21.37

$$
\begin{gather*}
W F\left(G^{ \pm}\right) \subset\left\{(t, x, \tau, \xi) ; \exists \quad(t, x, \tau, \xi, \omega, w) \in W F\left(u^{ \pm}\right)\right\} \\
G^{ \pm}=\left(\pi_{1}\right)_{*} u^{ \pm}=\int_{\mathbb{S}^{n-1}} u^{ \pm}(t, s ; \omega, \bar{x}) d \omega \tag{8.262}
\end{gather*}
$$

so here

$$
\begin{equation*}
(t, x, \tau, \xi, \omega, w) \in T^{*}\left(\mathbb{R} \times \mathbb{R}^{n} \times \mathbb{S}^{n-1}\right)=T^{*}\left(\mathbb{R} \times \mathbb{R}^{n}\right) \times T^{*} \mathbb{S}^{n-1} \tag{8.263}
\end{equation*}
$$

We claim that the singularities of $G_{\bar{x}}^{\prime}$ lie on a cone:

$$
\begin{equation*}
W F\left(G_{\bar{x}}^{\prime}\right) \subset \Lambda_{\bar{x}} \subset T^{*}\left(\mathbb{R} \times \mathbb{R}^{n}\right) \tag{8.264}
\end{equation*}
$$

where $\Lambda_{\bar{x}}$ is the conormal bundle to a cone:

$$
\begin{gather*}
\Lambda_{\bar{x}}=\operatorname{cl}\{(t, x ; \tau, \xi) ; t \neq 0, D(x, \bar{x})= \pm t \\
(\tau, \xi)=\tau\left(1, \mp d_{x} D(x, \bar{x})\right) \tag{8.265}
\end{gather*}
$$

where $D(x, \bar{x})$ is the Riemannian distance from $x$ to $\bar{x}$.

### 8.6. Forward fundamental solution

Last time we constructed a local parametrix for the Cauchy problem:

$$
\begin{cases}P G_{\bar{x}}^{\prime}=f \in \mathcal{C}^{\infty}(\Omega) & (0, \bar{x}) \in \Omega \subset \mathbb{R} \times X  \tag{8.266}\\ G_{\bar{x}}^{\prime} \upharpoonright t=0=u^{\prime} & \\ D_{t} G_{\bar{x}}^{\prime} \upharpoonright\{t=0\}=\delta_{\bar{x}}(x)+u^{\prime \prime} & u^{\prime}, u^{\prime \prime} \in \mathcal{C}^{\infty}\left(\Omega_{0}\right)\end{cases}
$$

where $P=D_{t}^{2}-\Delta$ is the wave operator for a Riemann metric on $X$. We also computed the wavefront set, and hence singular support of $G_{\bar{x}}$ and deduced that

$$
\begin{equation*}
\text { sing } \cdot \operatorname{supp} .\left(G_{\bar{x}}\right) \subset\{(t, x) ; d(x, \bar{x})=|t|\} \tag{8.267}
\end{equation*}
$$

in terms of the Riemannian distance.

This allows us to improve $\left(\frac{22.1}{8.266}\right)$ in a very significant way. First we can chop $G_{\bar{x}}$ off by replacing it by

$$
\begin{equation*}
\phi\left(\frac{t^{2}-d^{2}(x, \bar{x})}{\epsilon^{2}}\right) . \tag{8.269}
\end{equation*}
$$

where $\phi \in \mathcal{C}^{\infty}(\mathbb{R})$ has supporf near 0 and is identically equal to 1 in some neighbourhood of 0 . This gives ( $\overline{8.26} 6$ ) again, with $G_{\bar{x}}^{\prime}$ now supported in say $d^{2}<t^{2}+\epsilon^{2}$.

Next we can improve $\left(\frac{22.1}{8.266}\right)$ a little bit by arranging that

$$
\begin{equation*}
u^{\prime}=u^{\prime \prime}=0,\left.D_{t}^{k} f\right|_{t=0}=0 \forall k \tag{8.271}
\end{equation*}
$$

This just requires adding to $G^{\prime}$ a $\mathcal{C}^{\infty}, v$, function, so that

$$
\begin{equation*}
\left.v\right|_{t=0}=u^{\prime},\left.D_{t} v\right|_{t=0}=-u^{\prime \prime},\left.\quad D_{t}^{k}(P u)\right|_{t=0}=-\left.D_{t}^{k} f\right|_{t=0} \quad k>0 \tag{8.272}
\end{equation*}
$$

Once we have done this we consider

$$
\begin{equation*}
E_{\bar{x}}^{\prime}=i H(t) G_{\bar{x}}^{\prime} \tag{8.273}
\end{equation*}
$$

which now satisfies

$$
\begin{gather*}
P E_{\bar{x}}^{\prime}=\delta(t) \delta_{\bar{t}}(x)+F_{\bar{x}}, F_{\bar{x}} \in \mathcal{C}^{\infty}\left(\Omega_{\bar{x}}\right) \\
\operatorname{supp}\left(E_{x}^{\prime}\right) \subset\left\{d^{2}(x, \bar{x}) \leq t^{2}+\epsilon^{2}\right\} \cap\{t \geq 0\} \tag{8.274}
\end{gather*}
$$

Here $F$ vanishes in $t<0$, so vanishes to infinite order at $t=0$.

Next we remark that we can actually do all this with smooth dependence of $\bar{x}$. This should really be examined properly, but I will not do so to save time. Thus we actually have

$$
\left\{\begin{array}{l}
E^{\prime}(t, x, \bar{x}) \in \mathcal{C}^{-\infty}(P(-\infty, \epsilon) \times X \times X)  \tag{8.275}\\
P E^{\prime}=\delta(t) \sigma_{\bar{x}}(x)+F \\
\operatorname{supp} E^{\prime} \subset\left\{d^{2}(x, \bar{x}) \geq t^{2}+\epsilon^{2}\right\} \cap\{t \geq 0\}
\end{array}\right.
$$

We can, and later shall, estimate the wavefront set of $E$. In case $X=\mathbb{R}^{n}$ we can take $E$ to be the exact forward fundamental solution where $|x|$ or $\bar{x} \geq R$, so

$$
\begin{align*}
\operatorname{supp}(F) \subset\{ & \{\geq 0\} \cap\{|x|,|\bar{x}| \leq R\} \cap\left\{d^{2} \leq t^{2}+\epsilon^{2}\right\}  \tag{8.276}\\
F & \in \mathcal{C}^{\infty}((-\infty, \epsilon) \times X \times X)
\end{align*}
$$

Of course we want to remove $F$, the error term. We can do this because it is a Valterra operator, very similar to an upper triangular metric. Observe first that the operators of the form $\left(\frac{2.210}{8.276}\right)$ form an algebra under $t$-convolution:
22.11

$$
\begin{equation*}
F=F_{1} \circ F_{1}, F(t, x, \bar{x})=\int_{0}^{t} \int F_{1}\left(t,-t^{\prime}, x, x^{\prime}\right) F_{2}\left(t^{1}, x^{1}, \bar{x}\right) d x^{\prime} d t^{\prime} \tag{8.277}
\end{equation*}
$$

In fact if one takes the iterates of a fixed operator

$$
\begin{equation*}
F^{(k)}=F^{(k-1)} \circ F \tag{8.278}
\end{equation*}
$$

One finds exponential convergence:

$$
\begin{equation*}
\left|D_{x}^{\alpha} D_{t}^{p} F^{(k)}(t, x, \bar{x})\right| \leq \frac{C^{k+1} N, \delta}{k!}|t|^{N} \quad \text { in } t<\epsilon-\delta \forall N \tag{8.279}
\end{equation*}
$$

Thus if $F$ is as in $\left(\frac{22.10}{8.276}\right)$ then $I d+F$ has inverse $I d+\tilde{F}$,

$$
\begin{equation*}
\tilde{F}=\sum_{j \geq 1}(-1)^{j} F^{(j)} \tag{8.280}
\end{equation*}
$$

again of this form.
Next note that the composition of $E^{\prime}$ with $\tilde{F}$ is again of the form $\left(\frac{22.10}{8.276}\right)$, with $R$ increased. Thus

$$
\begin{equation*}
E=E^{\prime}+E^{\prime} \circ F \tag{8.281}
\end{equation*}
$$

is a forward fundamental solution, satisfying $\left(\frac{22.9}{8.275}\right)$ with $F \equiv 0$.
In fact $E$ is also a left parametrix, in an appropriate sense:

Proposition 8.6. Suppose $u \in \mathcal{C}^{-\infty}((-\infty, \epsilon) \times X)$ is such that (8.282) $\operatorname{supp}(u) \cap[-T, \tau] \times X$ is compact $\forall T$ and for $\tau<\epsilon$ then $P u=0 \Longrightarrow u=0$.

Proof. The trick is to make sense of the formula

$$
\begin{equation*}
0=E \cdot P u=u \tag{8.283}
\end{equation*}
$$

In fact the operators $G$ with kernel $G(t, x, \bar{x})$, defined in $t<\epsilon$ and such that $G * \phi \subset \mathcal{C}^{\infty} \forall \phi \in \mathcal{C}^{\infty}$ and
act. pn the space $\left(\frac{22.16}{8.282}\right)$ as $t$-convolution operators. For this algebra $E * P=\operatorname{Id}$ so ( 8.283 ) holds!

We can use this proposition to prove that $E$ itself is unique. Actually we want to do more.

$$
\begin{equation*}
\operatorname{supp}(E) \subset\{t \geq 0\}, P^{E}=\mathrm{Id} \tag{8.285}
\end{equation*}
$$

and

$$
\begin{equation*}
\operatorname{supp}(E) \subset\{(t, x, \bar{x}) \in \mathbb{R} \times X \times X ; d(x, \bar{x}) \leq t\} \tag{8.286}
\end{equation*}
$$

and further
22.22

$$
\begin{equation*}
W F^{\prime}(E) \subset \operatorname{Id} \cup \mathcal{F}_{+} \tag{8.287}
\end{equation*}
$$

where $\mathcal{F}_{+}$is the forward bicharacteristic relation on $T^{*}(\mathbb{R} \times X)$

$$
\begin{gather*}
\zeta=(t, x, \tau, \xi) \notin \Sigma(P) \Longrightarrow \mathcal{F}_{+}(\zeta)=\emptyset \\
\zeta=(t, x, \tau, \xi) \in \Sigma(P) \Longrightarrow \mathcal{F}_{+}(\zeta)=\left\{\zeta^{\prime}=\left(t^{\prime}, x^{\prime}, \tau^{\prime}, \xi^{\prime}\right)\right.  \tag{8.288}\\
\left.t^{\prime} \geq t \times \zeta^{\prime}=\exp \left(T H_{p}\right) \zeta \text { for some } T\right\}
\end{gather*}
$$

Proof. (1) Use $E_{1}$ defined in $(-\infty, \epsilon \times X$ to continue $E$ globally.
(2) U2. ${ }^{2}$ the freedom of choice of $\{t=0\}$ and uniqueness of $E$ to show that $\left(\frac{22.21}{8.26}\right)$ can $32.2 r r a n g e d$ for small, and hence all,
(3) Then get $\left(\frac{22.288)}{8.28 y}\right.$ by checking the wavefront set of $G$.

As corollary we get proofs of $\left(\frac{22.5}{8.270} 0\right)$ and $\left(\frac{22.6}{8.271}\right)$.

## Proof of Theorem XXI.5.

$$
\begin{equation*}
u(t, x)=\int E\left(t-t^{\prime}, x, x^{\prime}\right) f\left(t^{\prime}, x^{\prime}\right) d x^{\prime} d t^{\prime} \tag{8.289}
\end{equation*}
$$

Proof of Theorem XXI.6. We have to show that if both $\mathrm{WF}(P u) \not \supset z$ and $\mathrm{WF}(u) \not \supset z$ then $\exp \left(\delta H_{p}\right) z \notin W F(u)$ for small $\delta$. The general case that follows from the (assumed) connectedness of $H_{p}$ curves. This involves microlocal uniqueness of solutions of $P u=f$. Thus if $\phi \in \mathcal{C}^{\infty}(\mathbb{R})$ has support in $t>-\delta$, for $\delta>0$ small enough, $\pi^{*} t(z)=\bar{t}$

$$
\begin{equation*}
P(\phi(t-\bar{t}) u)=g \text { has } z \notin W F(g), \tag{8.290}
\end{equation*}
$$

and vanishes in $t<\delta$. Then

$$
\begin{gather*}
\phi(t-\bar{t}) u=E \times g \\
\Longrightarrow \exp \left(\tau H_{p}\right)(z) \notin W F(u) \text { for small } \tau . \tag{8.291}
\end{gather*}
$$

### 8.7. Operations on conormal distributions

I want to review and refine the push-forward theorem, in the general case, to give rather precise results in the conormal setting. Thus, suppose we have a projection

$$
\begin{equation*}
X \times Y @>x \gg X \tag{8.292}
\end{equation*}
$$

where we can view $X \times Y$ as compact manifolds or Euclidean spaces as desired, since we actually work locally. Suppose

## 23.2

$$
\begin{equation*}
Q \subset X \times Y \text { is an embeded submanifold. } \tag{8.293}
\end{equation*}
$$

Then we know how to define and examine the conormal distribution associated to $Q$. If

$$
\begin{equation*}
u \in I^{m}(X \times Y, Q ; \Omega) \tag{8.294}
\end{equation*}
$$

when is $\pi_{*}(u) \in \mathcal{C}^{-\infty}(X ; \Omega)$ conormal? The obvious thing we ned is a submanifold with respect to what it should be conormal! From our earlier theorem we know that

$$
\begin{equation*}
W F\left(\pi_{*}(u)\right) \subset\left\{(x, \xi) ; \quad \exists \quad(x, \xi, y, 0) \in W F(u) \subset N^{*} Q\right\} \tag{8.295}
\end{equation*}
$$

So suppose $Q=\left\{q_{j}(x, y)=0, j=1, \ldots, k\right\}, k=\operatorname{codim} Q$. Then we see that

> | 23.5 |
| :--- |

$$
(\bar{x}, \bar{\xi}, \bar{y}, 0) \in N^{*} Q \Longleftrightarrow(\bar{x}, \bar{y}) \in Q, \bar{\xi}=\sum_{j=1}^{k} \tau_{j} d_{x} q_{j}, \sum_{j=1}^{k} \tau_{j} d y q_{j}=0
$$

Suppose for a moment that $Q$ has a hypersurface, i.e. $k=1$, and that
23.6 (8.297)

$$
Q \longrightarrow \pi(Q) \text { is a fibration }
$$

then we expect
23.7 Theorem 8.6. $\pi_{*}: I^{m}(X \times Y, Q, \Omega) \longrightarrow I^{m^{\prime}}(X, \pi(Q))$.

Proof. Choose local coordinates so that

$$
\begin{align*}
Q & =\left\{x_{1}=0\right\}  \tag{8.298}\\
u & =\frac{1}{2 \pi} \int e^{i x_{1} \xi_{1}} a\left(x^{\prime}, y, \xi_{1}\right) d \xi_{1}  \tag{8.299}\\
\pi^{*} u & =\frac{1}{2 \pi} \int e^{i x_{1} \xi_{1}} b\left(x^{\prime}, \xi_{1}\right) d \xi_{1}  \tag{8.300}\\
b & =\int a\left(x^{\prime}, y, \xi\right) d y \tag{8.301}
\end{align*}
$$

Next consider the case of restriction to a submanifold. Again let us suppose $Q \subset X$ is a hypersurface and $Y \subset X$ is an embedded submanifold transversal to $Q$ :

$$
\begin{aligned}
& Q \pitchfork Y=Q Y \\
& \text { i.e. } T_{q} Q+T_{q} Y=T_{q} X \quad \forall q \in Q y \\
& \Longrightarrow Q_{y} \quad \text { is a hypersurface in } X \text {. }
\end{aligned}
$$

Indeed locally we can take coordinates in which

$$
\begin{equation*}
Q=\left\{x_{1}=0\right\}, Y=\left\{x^{\prime \prime}=0\right\}, \quad x=\left(x_{1}, x^{\prime}, x^{\prime \prime}\right) \tag{8.303}
\end{equation*}
$$

$$
\begin{equation*}
C_{Y}^{*}: I^{m}(X, Q) \longrightarrow I^{m+\frac{k}{4}}\left(Y, Q_{Y}\right) k=\operatorname{codim} Y \text { in } X \tag{8.304}
\end{equation*}
$$

Proof. In local coordinates as in ( $\left.{ }^{23.309} 3\right)$

$$
\begin{align*}
& u=\frac{1}{2 \pi} \int e^{i x_{1} \xi_{1}} a\left(x\left(x^{\prime}, x^{\prime \prime}, \xi_{1}\right)\right) d \xi, \\
& c^{*} u=\frac{1}{2 \pi} \int e^{i x_{1} \xi_{1}} a\left(x^{\prime}, 0, \xi_{1}\right) d \xi_{1} . \tag{8.305}
\end{align*}
$$

Now let's apply this to the fundamental solution of the wave equation. Well rather consider the solution of the initial value problem

$$
\left\{\begin{array}{l}
P G(t, x, \bar{x})=0  \tag{8.306}\\
G(0, x, \bar{x})=0 \\
D_{t} G(0, x, \bar{x})=\delta_{\bar{x}}(x)
\end{array}\right.
$$

We know that $G$ exists for all time and that for short time it is

$$
\begin{equation*}
G-\int_{\mathbb{S}^{n-1}}\left(u_{+}(t, x, \bar{x} ; \omega)+u_{-}(t, x, \bar{x} ; \omega)\right) d \omega+\mathcal{C}^{\infty} \tag{8.307}
\end{equation*}
$$

where $u_{ \pm}$are conormal for the term characteristic hypersurfaces $H_{p}$ satisfying

$$
\begin{gather*}
N^{*} H_{ \pm} \subset \Sigma(P) \\
H_{ \pm} \cap\{t=0\}=\{(x-\bar{x}) \cdot \omega=0\} \tag{8.308}
\end{gather*}
$$

Consider the $2 \times 2$ matrix of distribution

$$
U(t)=\left(\begin{array}{cc}
D_{t} G & G  \tag{8.309}\\
D_{t}^{2} G & D_{t} G
\end{array}\right) .
$$

Since $W F U \subset \Sigma(P)$, in polar $\tau \neq 0$ we can consider this as a smooth function of $t$, with values in distribution on $X \times X$.
23.15 Theorem 8.8. For each $t \in \mathbb{R} U(t)$ is a boundary operator on $L^{2}(X) \oplus H^{\prime}(X)$ such that

$$
\begin{equation*}
U(t)\binom{u_{0}}{u_{1}}=\binom{u(t)}{D_{t} u(t)} \tag{8.310}
\end{equation*}
$$

where $u(t, x)$ is the unique solution of

$$
\begin{align*}
& \left(D_{t}^{2}-\Delta\right) u(t)=0 \\
& u(0)=u_{0}  \tag{8.311}\\
& D_{t}+u(0)=u_{1} .
\end{align*}
$$

Proof. Just check it!
Consider again the formula ( ${ }^{(23.307)}$. First notice that at $x=\bar{x}, t=0, d H^{ \pm}=$ $d t \pm d(x-\bar{x}) \omega$ ) (by construction). so

$$
\begin{equation*}
H_{ \pm} \pitchfork\{x=\bar{x}\}=\{t=0\} \subset \mathbb{R} \times X \hookrightarrow \mathbb{R} \times X \times Y \times \mathbb{S}^{n-1} \tag{8.312}
\end{equation*}
$$

Moreover the projection
23.19

$$
\begin{equation*}
\mathbb{R} \times X \times \mathbb{S}^{n-1} \longrightarrow \mathbb{R} \tag{8.313}
\end{equation*}
$$

clearly fibres $\{t=0\}$ over $\{t=0\} \in=\{0\} \subset \mathbb{R}$. Then we can apply the two theorems, on push-forward and pull-back, above to conclude that

$$
\begin{equation*}
T(t)=\int_{X} G(t, x, \bar{x}) \upharpoonright x=\bar{x} d x \in \mathcal{C}^{-\infty}(\mathbb{R}) \tag{8.314}
\end{equation*}
$$

is conormal near $t=0$ i.e. $\mathcal{C}^{\infty}$ in $(-\epsilon, \epsilon) \backslash\{0\}$ for some $\epsilon>0$ and conormal at 0 . Moreover, we can, at least in principle, work at the symbol of $T(t)$ at $t=0$. We return to this point next time.
23. For the moment let us think of a more 'fundamental analytic' interpretation of ( 8.314 ). By this I mean

## (8.315) <br> $$
T(t)=\operatorname{tr} U(t)
$$

REmARK 8.1. Trace class operators $\Delta \lambda$; Smoother operators are trace order, $t r=\int K(x, x)$

$$
\begin{gather*}
\int U(t) \phi(t) \text { is smoothing }  \tag{8.316}\\
\langle T(t), \phi(t)\rangle=\operatorname{tr}\langle U(t), \phi(t)\rangle . \tag{8.317}
\end{gather*}
$$

### 8.8. Weyl asymptotics

Let us summarize what we showed last time, and a little more, concerning the trace of the wave group

Proposition 8.7. Let $X$ be a compact Riemann manifold and $U(t)$ the wave group, so
(8.318) $U(t): \mathcal{C}^{\infty}(X) \times \mathcal{C}^{\infty}(X) \ni\left(u_{0}, u_{1}\right) \mapsto(u,(t), D+t u(t)) \in \mathcal{C}^{\infty}(X) \times \mathcal{C}^{\infty}(X)$ where $u$ is the solution to

$$
\begin{align*}
\left(D_{t}^{2}-\Delta\right) u(t) & =0 \\
u(0) & =u_{0}  \tag{8.319}\\
D_{t} u(0) & =u_{1} .
\end{align*}
$$

The trace of the wave group, $T \in \mathcal{S}^{\prime}(\mathbb{R})$, is well-defined by

$$
\begin{equation*}
T(\phi)=\operatorname{Tr} U(\phi), U(\phi)=\int U(t) \phi(t) d t \forall \phi \in \mathcal{S}(\mathbb{R}) \tag{8.320}
\end{equation*}
$$

and satisfies

$$
\begin{equation*}
T=Y\left(\left(1+\sum_{j=1}^{\infty} 2 \cos \left(t \lambda_{j}\right)\right)\right. \tag{8.321}
\end{equation*}
$$

$$
\begin{equation*}
\text { where } 0=\lambda_{0}<\lambda_{1}^{2} \leq \lambda_{2}^{2} \ldots \quad \lambda_{j} \geq 0 \tag{8.322}
\end{equation*}
$$

is the spectrum of the Laplacian repeated with multiplicity

$$
\begin{equation*}
\operatorname{sing} \cdot \operatorname{supp}(T) \subset \mathcal{L} \cup\{0\} \cup-\mathcal{L} \tag{8.323}
\end{equation*}
$$

where $\mathcal{L}$ is the set of lengthes of closed geodesics of $X$ and

$$
\begin{gather*}
\text { if } \psi \in \mathcal{C}_{c}^{\infty}(\mathbb{R}), \psi(t)=0 \text { if }|t| \geq \text { inf } \mathcal{L}-\epsilon, \epsilon>0 \\
\psi T \in I(\mathbb{R},\{0\})  \tag{tabular}\\
\sigma(\psi T)=
\end{gather*}
$$

Proof. We have already discussed $\left(\frac{24.4}{8.321} 1\right)$ and the first part of $\left(\frac{24.6}{8.324} 4\right)$ (given $\left.\left(\frac{24.5}{8.32} 3\right)\right)$. Thus we need to show ( $\left(\frac{1.323}{} 3\right)$, the Poisson relation, and compute the symbol of $T$ as a cononormal distribution at 0 .

Let us recall that if $G$ is the solution to

$$
\begin{align*}
\left(D_{t}^{2}=\Delta\right) G(t, x, \bar{x}) & =0 \\
G(0, x, \bar{x}) & =0  \tag{8.325}\\
D_{t} G(0, x, \bar{x}) & =\delta_{\bar{x}}(x)
\end{align*}
$$

then

$$
\begin{equation*}
T=\pi_{*}\left(\iota_{\Delta}^{*} 2 D_{t} G\right) \tag{8.326}
\end{equation*}
$$

where
24.9 (8.327)

$$
\iota_{\Delta}: \mathbb{R} \times X \hookrightarrow \mathbb{R} \times X \times X
$$

is the embedding of the diagonal and

$$
\begin{equation*}
\pi: \mathbb{R} \times X \longrightarrow \mathbb{R} \tag{8.328}
\end{equation*}
$$

is projective. We also know about the wavefront set of $G$. That is,

$$
\begin{align*}
& W F(G) \subset\left\{(t, x, \bar{x}, \tau, \xi, \bar{\xi}) ; \tau^{2}=|\xi|^{2}=|\bar{\xi}|^{2}\right. \\
& \left.\exp \left(s H_{p}\right)(0, \bar{x}, \tau, \bar{\xi})=(t, x, \tau, \xi), \text { some } s\right\} \tag{8.329}
\end{align*}
$$

Let us see what $\left(\frac{24.11}{8.329}\right)$ says about the wavefront set of $T$. First under the restriction map to $\mathbb{R} \times \Delta$

$$
\begin{gather*}
W F\left(\iota_{\Delta}^{*} D_{t} G\right) \subset\{(t, y, \tau, \eta) ; \quad \exists \\
\quad(t, x, y, \tau, \xi, \bar{\xi}) ; \eta=\xi-\bar{\xi}\} \tag{8.330}
\end{gather*}
$$

Then integration gives

$$
\begin{equation*}
W F(T) \subset\left\{(t, \tau) ; \quad \exists \quad(t, y, \tau, 0) \in W F\left(D_{t} G\right)\right\} \tag{8.331}
\end{equation*}
$$

$\left(\frac{24.12}{8.330}\right)$ and $\left(\frac{24.13}{8.331}\right)$ we see

$$
\begin{gathered}
t \in \operatorname{sing} \cdot \operatorname{supp}(T) \Longrightarrow \quad \exists \quad(t, \tau) \in W F(T) \\
\Longrightarrow \quad \exists \quad(t, x, x, \tau, \xi, \xi) \in W F\left(D_{t} G\right) \\
\Longrightarrow \quad \exists \mathrm{s} \text { s.t. } \quad \exp \left(s H_{p}\right)(0, x, \tau, \xi)=(t, x, \tau, \xi) .
\end{gathered}
$$

Now

$$
\begin{equation*}
p=\tau^{2}-|\xi|^{2}, \text { so } H_{p}=2 \tau \partial_{t}-H_{g}, g=|\xi|^{2} \tag{8.333}
\end{equation*}
$$

$H_{g}$ being a vector field on $T^{*} X$. Since $W F$ is conic we can take $|\xi|=1$ in the last condition in ( 8.332 ). Then it says

$$
\begin{equation*}
s=2 \tau t, \quad \exp \left(t H_{\frac{1}{2} g}\right)(x, \xi)=(x, \xi) \tag{8.334}
\end{equation*}
$$

since $\tau^{2}=1$.
The curves in $X$ with the property that their tangent vectors have unit length and the lift to $T^{*} X$ is an integral curve of $H_{\frac{1}{2} g}$ are by definition geodesic, parameterized by arclength. Thus.5 $\left(\frac{2.334)}{}\right.$ is the statement that $|t|$ is the length of a closed geodesic. This proves ( $\overline{8.323}$ ).

So now we have to compute the symbol of $T$ at 0 . We use, of course, our local representation of $G$ in terms of conormal distributions. Namely

$$
\begin{equation*}
G=\sum_{j} \phi_{j} G_{j}, \quad \phi_{j} \in \mathcal{C}^{\infty}(X) \tag{8.335}
\end{equation*}
$$

where the $\phi_{j}$ has support in coordinate particles in which

$$
\begin{align*}
G_{j}(t, x, \bar{x}) & =\int_{\mathbb{S}^{n-1}}\left(u_{+}(t, x, \bar{x} ; \omega)+u_{-}(t, x, \bar{x} ; \omega)\right) d \omega \\
u_{p} m & =\frac{1}{2 \pi} \int_{\xi} e^{i h_{ \pm}(t, x, \bar{x}, \omega) \xi} a_{ \pm}(x, \bar{x}, \xi, \omega) d \xi \tag{8.336}
\end{align*}
$$

Here $h_{ \pm}$are solutions of the eikonal equation (i.e. are characteristic for $P$ )

$$
\begin{gather*}
\left|\partial_{t} h_{ \pm}\right|^{2}=\left|h_{ \pm}\right|^{2} \\
\left.h_{ \pm}\right|_{t=0}=(x-\bar{x}) \cdot \omega  \tag{8.337}\\
\pm \partial_{t} h_{ \pm}>0
\end{gather*}
$$

which fixes them locally uniquely. The $a_{ \pm}$are chosen so that

$$
\begin{equation*}
\left(u_{+}+\left.u_{ \pm}\right|_{t=0}=0,\left.\left(D_{t} u_{+} D_{t} u_{-}\right)\right|_{t=0} \delta((x-\bar{x}) \cdot \omega) P u_{ \pm} \in \mathcal{C}^{\infty}\right. \tag{8.338}
\end{equation*}
$$

Now, from ( $\left.\frac{24.17}{8.336}\right)$

$$
\begin{equation*}
u_{+}+\left.u_{-}\right|_{t=0}=\frac{1}{2 \pi} \int e^{((x-x \bar{x}) \cdot \omega) \xi}\left(a_{+}+a_{-}\right)(x, \bar{x}, \xi, \omega) d \xi=0 \tag{8.339}
\end{equation*}
$$

so $a_{+}-a_{-}$. Similarly

$$
\begin{align*}
D_{t} u_{+}+\left.D_{t} u_{-}\right|_{t=0} & =\frac{1}{2 \pi} \int e^{i((x-\bar{x}) \cdot \omega) \xi}\left[\left(D_{t} h_{+}\right) a_{+}+\left(D_{t} h_{-}\right) a_{-}\right] d \xi \\
& =\frac{1}{2(2 \pi)^{n-1}} f_{n}((x-\bar{x}) \cdot \omega) \tag{8.340}
\end{align*}
$$

From $\left(\frac{24.18}{8.337}\right)$ we know that $D_{t} h_{ \pm}=\mp i\left|d_{x}(x-\bar{x}) \cdot \omega\right|=\mp i|\omega|$ where the length is with respect to the Riemann measure. We can compute the symbols or both sides in $(8.340)$ and consider that

$$
\begin{equation*}
-2 i|\omega| a_{+} \equiv \frac{1}{2(2 \pi)^{n-1}}|\xi|^{n-1}=D_{t} h_{+} a_{+}+\left.D_{t} h_{-} a_{-}\right|_{t=0} \tag{8.341}
\end{equation*}
$$

is necessary to get ( $\left.\frac{24.19}{8.338}\right)$. Then
24.22

$$
\begin{align*}
T(t) & =2 \pi_{*}\left(\iota_{\Delta}^{*} D_{t} G\right) \\
& =\frac{1}{2 \pi} \sum_{j, \pm} 2 \int_{X} \int_{\mathbb{S} n-1} e^{i h_{ \pm}(t, x, x, \omega) \xi}\left(D_{t} h_{ \pm} a_{ \pm}\right)(x, \bar{x}, \omega, \xi) d \xi d \omega d x \tag{8.342}
\end{align*}
$$

Here $d x$ is really the Riemann measure on $X$. From $\left(\frac{24.21}{8.341}\right)$ the leading part of this is

$$
\begin{equation*}
\frac{2}{2 \pi} \sum_{j \pm} \int_{X} \int_{\mathbb{S} n-1} e^{i h_{ \pm}(t, x, x, \omega) \xi} \frac{1}{4(2 \pi)^{n-1}}|\xi|^{n-1} d \xi d \omega d x \tag{8.343}
\end{equation*}
$$

since any term vanishes at $t$ contributes a weaker singularity. Now

$$
\begin{equation*}
h_{ \pm}= \pm|\omega| t+(x-\bar{x}) \cdot \omega+0\left(t^{2}\right) \tag{8.344}
\end{equation*}
$$

From which we deduce that

$$
\begin{gather*}
\psi(t) T(t)=\frac{1}{2 \pi} \int e^{i t \tau} a(\tau) d \tau  \tag{8.345}\\
a(\tau) \sim C_{n} \operatorname{Vol}(X)|\tau|^{n-1} C_{n}=
\end{gather*}
$$

where $C_{n}$ is a universal constant depending only on dimension. Notice that if $n$ is odd this is a "little" function.

The final thing I want to do is to show how this can be used to describe the asymptotic behaviour of the eigenvalue of $\Delta$ :

Proposition 8.8. ("Weyl estimates with optimal remainder".) If $N(\lambda)$ is the number of eigenvalues at $\Delta$ satisfying $\lambda_{1}^{2} \leq \lambda$, counted with multiplicity, the

$$
\begin{equation*}
N(\lambda)=C_{n} \operatorname{Vol}(X) \lambda^{n}+o\left(\lambda^{n-1}\right) \tag{8.346}
\end{equation*}
$$

The estimate of the remainder terms is the here - weaker estimates are easier to get.

Proof. (Tauberian theorem). Note that

$$
\begin{equation*}
T=\mathcal{F}(\mu) \text { where } N(\lambda)=\int_{0}^{\lambda} \mu(\lambda) \tag{8.347}
\end{equation*}
$$

$\mu(\lambda)$ being the measure

$$
\begin{equation*}
\mu(\lambda)=\sum_{\lambda_{j}^{2} \in \operatorname{spec}(\Delta)} \delta\left(\lambda-\lambda_{j}\right) \tag{8.348}
\end{equation*}
$$

Now suppose $\rho \in \mathcal{S}(\mathbb{R})$ is even and $\int \rho=1, \rho \geq 0$. Then $N_{\rho}(\lambda)=\int\left(\lambda^{\prime}\right) \rho\left(\lambda-\lambda^{\prime}\right)$ is a $\mathcal{C}^{\infty}$ function. Moreover

$$
\begin{equation*}
\widehat{\frac{d}{d \lambda} N_{\rho}(\lambda)}=\hat{\mu} \cdot \hat{\rho} \tag{8.349}
\end{equation*}
$$

Suppose we can choose $\rho$ so that
24.27

$$
\begin{equation*}
\rho \geq 0, \int \rho=1, \rho \in \mathcal{S}, \hat{\rho}(t)=0,|t|>\epsilon \tag{8.350}
\end{equation*}
$$

for a given $\epsilon>0$. Then we know $\hat{\mu} \hat{\rho}$ is conormal and indeed

$$
\begin{gather*}
\frac{d}{d \lambda} N \rho(\lambda) \sim C \operatorname{Vol}(X) \lambda^{n-1}+\ldots  \tag{8.351}\\
\Longrightarrow N_{\rho}(\lambda) \sim C^{\prime} \operatorname{Vol}(X) \lambda^{n}+\text { lots }
\end{gather*}
$$

So what we need to do is look at the difference

$$
\begin{equation*}
N_{\rho}(\lambda)-N(\lambda)=\int N\left(\lambda-\lambda^{\prime}\right) \rho\left(\lambda^{\prime}\right)-N(\lambda) \rho\left(\lambda^{\prime}\right) \tag{8.352}
\end{equation*}
$$

It follows that a bound for $N$

$$
\begin{equation*}
|N(\lambda+\mu)-N(\lambda)| \leq\left((1+|\lambda|+|\mu|)^{n-1}(1+|\lambda|)\right. \tag{8.353}
\end{equation*}
$$

gives

$$
\begin{equation*}
N(\lambda)-N_{\rho}(\lambda) \leq C \lambda^{n-1} \tag{8.354}
\end{equation*}
$$

which is what we want. Now $\left(\frac{24.31}{8.355)}\right.$ follows if we have
24.31

$$
\begin{equation*}
N(\lambda+1)-N(\lambda) \leq C(1+|\lambda|) \quad t / \lambda \tag{8.355}
\end{equation*}
$$

This in turn follows from the positivity of $\rho$, since

$$
\begin{equation*}
\int \rho\left(\lambda-\lambda^{\prime}\right) \mu\left(\lambda^{\prime}\right) \leq C(1+|\lambda|)^{n-1} \tag{8.356}
\end{equation*}
$$

Finally then we need to check the existence of $\rho$ as in $\left(\frac{24.27}{8.350}\right)$. If $\phi$ is real and even so is $\hat{\phi}$. Take $\phi$ with support in $\left(-\frac{\epsilon}{2}, \frac{\epsilon}{2}\right)$ and construct $\phi * \phi$, real and even with $\phi$.

### 8.9. Problems

Problem 8.5. Show that if $E$ is a symplectic vector space, with non-degenerate bilinear form $\omega$, then there is a basis $v_{1}, \ldots, v_{n}, w_{1}, \ldots, w_{n}$ of $E$ such that in terms of the dual basis of $E^{*}$

$$
\begin{equation*}
\omega=\sum_{j} v_{j}^{*} \wedge w_{j}^{*} \tag{8.357}
\end{equation*}
$$

Hint: Construct the $w_{j}, v_{j}$ successive. Choose $v_{1} \neq 0$. Then choose $w_{1}$ so that $\omega\left(v_{1}, w_{1}\right)=1$. Then choose $v_{2}$ so $\omega\left(v_{1}, v_{2}\right)=\omega\left(w_{1}, v_{2}\right)=0$ (why is this possible?) and $w_{i}$ so $\omega\left(v_{2}, w_{2}\right)=1, \omega\left(v_{1}, w_{2}\right)=\psi\left(w_{1}, w_{2}\right)=0$. Then proceed and conclude that ( 8.357 ) must hold.

Deduce that there, is a linear transformation $T: E \longrightarrow \mathbb{R}^{2 n}$ so that $\omega=T^{*} \omega_{D}$, with $\omega_{D}$ given by (ㅇ.137).
 given Lagrangian plane $V \subset E$ to

$$
\begin{equation*}
\{x=0\} \subset \mathbb{R}^{2 n} \tag{8.358}
\end{equation*}
$$

Hint: Construct the basis choosing $v_{j} \in V \forall j$ !
19.47

Problem 8.7. Suppose $S$ is a symplectic manifold. Show that the Poisson bracket

$$
\begin{equation*}
\{f, g\}=H_{f} g \tag{8.359}
\end{equation*}
$$

makes $\mathcal{C}^{\infty}(S)$ into a Lie algebra.

## CHAPTER 9

## K-theory

### 9.1. Vector bundles

9.2. The ring $K(X)$
9.3. Chern-Weil theory and the Chern character
9.4. $K_{1}(X)$ and the odd Chern character
9.5. $\mathcal{C}^{*}$ algebras
9.6. K-theory of an algeba
9.7. The norm closure of $\Psi^{0}(X)$
9.8. The index map

### 9.9. Problems

## CHAPTER 10

## Hochschild homology

### 10.1. Formal Hochschild homology

The Hochschild homology is defined, formally, for any associative algebra. Thus if $\mathcal{A}$ is the algebra then the space of formal $k$-chains, for $k \in \mathbb{N}_{0}$ is the $(k+1)$-fold tensor product

$$
\begin{equation*}
\mathcal{A}^{\otimes(k+1)}=\mathcal{A} \otimes \mathcal{A} \otimes \cdots \otimes \mathcal{A} \tag{10.1}
\end{equation*}
$$

The 'formal' here refers to the fact that for the 'large' topological algebras we shall consider it is wise to replace this tensor product by an appropriate completion, usually the 'projective' tensor product. At the formal level the differential defining the cohomolgy is given in terms of the product, $\star$, by

$$
\begin{gathered}
b\left(a_{0} \otimes a_{1} \otimes \cdots \otimes a_{k}\right)=b^{\prime}\left(a_{0} \otimes a_{1} \otimes \cdots \otimes a_{k}\right)+(-1)^{k}\left(a_{0} \star a_{k}\right) \otimes a_{1} \otimes \cdots \otimes a_{k-1}, \\
b^{\prime}\left(a_{0} \otimes a_{1} \otimes \cdots \otimes a_{k}\right)=\sum_{j=0}^{k-1}(-1)^{j} a_{0} \otimes \cdots \otimes a_{j-1} \otimes a_{j+1} \star a_{j} \otimes a_{j+2} \otimes \cdots \otimes a_{k} .
\end{gathered}
$$

Lemma 10.1. Both the partial map, $b^{\prime}$, and the full map, $b$, are differentials, that is

$$
\begin{equation*}
\left(b^{\prime}\right)^{2}=0 \text { and } b^{2}=0 \tag{10.3}
\end{equation*}
$$

Proof. This is just a direct computation. From (HHdifferential ${ }^{(10.2) \text { it follows }}$ that

$$
\begin{equation*}
\left(b^{\prime}\right)^{2}\left(a_{0} \otimes a_{1} \otimes a_{2} \otimes \cdots \otimes a_{m}\right) \tag{10.4}
\end{equation*}
$$

$$
\begin{aligned}
& =\sum_{j=2}^{m-1} \sum_{p=0}^{j-2}(-1)^{j}(-1)^{p}\left(\cdots \otimes a_{p+1} \star a_{p} \otimes \cdots \otimes a_{j-1} \otimes a_{j+1} \star a_{j} \otimes a_{j+2} \otimes \cdots \otimes a_{m}\right) \\
& \quad-\sum_{j=1}^{m-1}\left(\cdots \otimes a_{j+1} \star a_{j} \star a_{j-1} \otimes \cdots\right)-\sum_{j=0}^{m-2}\left(\cdots \otimes a_{j+21} \star a_{j+1} \star a_{j} \star \otimes \cdots\right) \\
& +\sum_{j=0}^{m-3} \sum_{p=j+2}^{m-1}(-1)^{j}(-1)^{p-1}\left(a_{0} \otimes \cdots \otimes a_{j-1} \otimes a_{j+1} \star a_{j} \otimes a_{j+2} \otimes \cdots \otimes a_{p+1} \star a_{p} \otimes \cdots\right)=0 .
\end{aligned}
$$

Similarly, direct computation shows that

$$
\begin{aligned}
& \left(b-b^{\prime}\right) b^{\prime}\left(a_{0} \otimes \cdots \otimes a_{m}\right)=(-1)^{m-1}\left(a_{1} \star a_{0} \star a_{m} \otimes \cdots a_{m-1}\right) \\
& +\sum_{i=1}^{m-2}(-1)^{i+m-1}\left(a_{0} \star a_{m} \otimes \cdots \otimes a_{i+1} \star a_{i} \otimes \cdots\right)+\left(a_{0} \star a_{m} \star a_{m-1} \otimes \cdots\right), \\
& b^{\prime}\left(b-b^{\prime}\right)\left(a_{0} \otimes \cdots \otimes a_{m}\right)=(-1)^{m}\left(a_{1} \star a_{0} \star a_{m} \otimes \cdots a_{m-1}\right) \\
& \quad+\sum_{i=1}^{m-2}(-1)^{i+m}\left(a_{0} \star a_{m} \otimes \cdots \otimes a_{i+1} \star a_{i} \otimes \cdots\right) \text { and } \\
& \quad\left(b-b^{\prime}\right)^{2}\left(a_{0} \otimes \cdots \otimes a_{m}\right)=-\left(a_{0} \star a_{m} \star a_{m-1} \otimes \cdots\right)
\end{aligned}
$$

so

$$
\begin{equation*}
\left(b-b^{\prime}\right) b^{\prime}+b^{\prime}\left(b-b^{\prime}\right)=-\left(b-b^{\prime}\right)^{2} . \tag{10.5}
\end{equation*}
$$

The difference between these two differentials is fundamental, roughly speaking $b^{\prime}$ is 'trivial'.

Lemma 10.2. For any algebra with identity the differential $b^{\prime}$ is acyclic, since it satifies

$$
\begin{gather*}
b^{\prime} s+s b^{\prime}=\mathrm{Id} \text { where }  \tag{10.6}\\
s\left(a_{0} \otimes \cdots \otimes a_{m}\right)=\operatorname{Id} \otimes a_{0} \otimes \cdots \otimes a_{m} . \tag{10.7}
\end{gather*}
$$

Proof. This follows from the observation that

$$
\begin{equation*}
b^{\prime}\left(\operatorname{Id} \otimes a_{0} \otimes \cdots \otimes a_{m}\right)=a_{0} \otimes \cdots \otimes a_{m}+\sum_{i=1}^{m}(-1)^{i}\left(\operatorname{Id} \otimes \cdots a_{i} \star a_{i-2} \otimes \cdots\right) \tag{10.8}
\end{equation*}
$$

Definition 10.1. An associative algebra is said to be H -unital if its b' complex is acyclic.

Thus the preceeding lemma just says that every unital algebra is H-unital.

### 10.2. Hochschild homology of polynomial algebras

Consider the algebra $\mathbb{C}[x]$ of polynomials in $n$ variables ${ }^{1}, x \in \mathbb{R}^{n}$ (or $x \in \mathbb{C}^{n}$ it makes little difference). This is not a finite dimensional algebra but it is filtered by the finite dimensional subspaces, $P_{m}[x]$, of polynomials of degree at most $m$;

$$
\mathbb{C}[x]=\bigcup_{m=0}^{\infty} P_{m}[x], P_{m}[x] \subset P_{m+1}[x] .
$$

Furthermore, the Hochschild differential does not increase the total degree so it is enough to consider the formal Hochschild homology.

The chain spaces, given by the tensor product, just consist of polynomials in $n(k+1)$ variables

$$
(\mathbb{C}[x])^{\hat{\otimes}(k+1)}=\mathbb{C}\left[x_{0}, x_{1}, \ldots, x_{k}\right], x_{j} \in \mathbb{R}^{n}
$$

[^14]Furthermore composition acts on the tensor product by

$$
p\left(x_{0}\right) q\left(x_{1}\right)=p \otimes q \longmapsto p\left(x_{0}\right) q\left(x_{0}\right)
$$

which is just restriction to $x_{0}=x_{1}$. Thus the Hochschild differential can be written

$$
\begin{aligned}
& b: \mathbb{C}\left[x_{0}, \ldots, x_{k}\right] \longrightarrow \mathbb{C}\left[x_{0}, \ldots, x_{k-1}\right] \\
&(b q)\left(x_{0}, x_{1}, \ldots, x_{k-1}\right)=\sum_{j=0}^{k-1}(-1)^{j} p\left(x_{0}, \ldots, x_{j-1}, x_{j}, x_{j}, x_{j+1}, \ldots, x_{k-1}\right) \\
&+(-1)^{k} q\left(x_{0}, x_{1}, \ldots, x_{k-1}, x_{0}\right)
\end{aligned}
$$

One of the fundamental results on Hochschild homology is
THEOREM 10.1. The Hochschild homology of the polynomail algebra in $n$ variables is

$$
\begin{equation*}
\mathrm{HH}_{k}(\mathbb{C}[x])=\mathbb{C}[x] \otimes \Lambda^{k}\left(\mathbb{C}^{n}\right) \tag{10.9}
\end{equation*}
$$

with the identification given by the map from the chain spaces

$$
\left.\mathbb{C}\left[x_{0}, \ldots, x_{k}\right] \ni q \longrightarrow \sum_{1 \leq j_{i} \leq n} \frac{\partial}{\partial x_{1}^{j_{1}}} \ldots \frac{\partial}{\partial x_{k}^{j_{k}}} p\right|_{x=x_{0}=x_{1}=\cdots=x_{k}} d x_{1}^{j_{1}} \wedge \cdots \wedge d x_{k}^{j_{k}} .
$$

24.91

Note that the appearance of the original algebra $\mathbb{C}[x]$ on the left in $\left(\frac{124.91}{10.9)}\right.$ is not surprising, since the differential commutes with multilplication by polynomails in the first variable, $x_{0}$

$$
b\left(r\left(x_{0}\right) q\left(x_{0}, \ldots, x_{k}\right)\right)=r\left(x_{0}\right)\left(b q\left(x_{0}, \ldots, x_{k}\right)\right)
$$

Thus the Hochschild homology is certainly a module over $\mathbb{C}[x]$.
Proof. Consider first the cases of small $k$. If $k=0$ then $b$ is identically 0 . If $k=1$ then again

$$
(b q)\left(x_{0}\right)=q\left(x_{0}, x_{0}\right)-q\left(x_{0}, x_{0}\right)=0
$$

vanishes identically. Thus the homology in dimension 0 is indeed $\mathbb{C}[x]$.
Suppose that $k>1$ and consider the subspace of $\mathbb{C}\left[x_{0}, x_{1}, \ldots, x_{k}\right]$ consisting of the elements which are independent of $x_{1}$. Then the first two terms in the definition of $b$ cancel and

$$
\begin{aligned}
&(b q)\left(x_{0}, x_{1}, \ldots, x_{k-1}\right)=\sum_{j=2}^{k-1}(-1)^{j} p\left(x_{0}, \ldots, x_{j-1}, x_{j}, x_{j}, x_{j+1}, \ldots, x_{k-1}\right) \\
&+(-1)^{k} q\left(x_{0}, x_{1}, \ldots, x_{k-1}, x_{0}\right), \partial_{x_{1}} q \equiv 0
\end{aligned}
$$

It follows that $b q$ is also independent of $x_{1}$. Thus there is a well-defined subcomplex on polynomails independend of $x_{1}$ given by

$$
\begin{aligned}
& \mathbb{C}\left[x_{0}, x_{2}, \ldots, x_{k}\right] \ni q \longmapsto(\tilde{b} q)\left(x_{0}, x_{2}, \ldots, x_{k-1}\right) \\
&=\sum_{j=2}^{k-1}(-1)^{j} p\left(x_{0}, x_{2}, x_{2}, x_{3} \ldots, x_{k-1}\right)+\sum_{j=3}^{k-1}(-1)^{j} \\
& p\left(x_{0}, \ldots, x_{j-1}, x_{j}, x_{j}, x_{j+1}, \ldots, x_{k-1}\right)+(-1)^{k} q\left(x_{0}, x_{2}, \ldots, x_{k-1}, x_{0}\right)
\end{aligned}
$$

The reordering of variables $\left(x_{0}, x_{2}, x_{3}, \ldots, x_{k}\right) \longrightarrow\left(x_{2}, x_{3}, \ldots, x_{k}, x_{0}\right)$ for each $k$, transforms $\tilde{b}$ to the reduced Hochschild differential $b^{\prime}$ acting in $k$ variables. Thus $\tilde{b}$ is acyclic.

Similarly consider the subspace of $\mathbb{C}\left[x_{0}, x_{1}, \ldots, x_{k}\right]$ consisting of the polynomials which vanish at $x_{1}=x_{0}$. Then the first term in the definition of $b$ vanishes and the action of the differential becomes

$$
\begin{align*}
& (b q)\left(x_{0}, x_{1}, \ldots, x_{k-1}\right)=p\left(x_{0}, x_{1}, x_{1}, x_{2}, \ldots, x_{k-1}\right)+  \tag{10.10}\\
& \quad \sum_{j=2}^{k-1}(-1)^{j} p\left(x_{0}, \ldots, x_{j-1}, x_{j}, x_{j}, x_{j+1}, \ldots, x_{k-1}\right) \\
& \quad+(-1)^{k} q\left(x_{0}, x_{1}, \ldots, x_{k-1}, x_{0}\right), \text { if } b\left(x_{0}, x_{0}, x_{2}, \ldots\right) \equiv 0
\end{align*}
$$

It follows that $b q$ also vanishes at $x_{1}=x_{0}$.
By Taylor's theorem any polynomial can be written uniquely as a sum

$$
q\left(x_{0}, x_{1}, x_{2}, \ldots, x_{k}\right)=q_{1}^{\prime}\left(x_{0}, x_{1}, x_{2}, \ldots, x_{k}\right)+q^{\prime \prime}\left(x_{0}, x_{2}, \ldots, x_{k}\right)
$$

of a polynomial which vanishes at $x_{1}=x_{0}$ and a polynomial which is independent of $x_{1}$. From the discussion above, this splits the complex into a sum of two subcomplexes, the second one of which is acyclic. Thus the Hochschild homology is the same as the homology of $b$, which is then given by (년․I0), acting on the spaces
24.93

$$
\begin{equation*}
\left\{q \in \mathbb{C}\left[x_{0}, x_{1}, \ldots, x_{k}\right] ; q\left(x_{0}, x_{1}, \ldots\right)=0\right\} . \tag{10.11}
\end{equation*}
$$

This argument can be extended iteratively. Thus, if $k>2$ then $b$ maps the subspace of ( $1 \mathrm{IT.III}$ ) of functions independent of $x_{2}$ to functions independent of $x_{2}$ and on these subspaces acts as $b^{\prime}$ in $k-2$ variables; it is therefore acyclic. Similar it acts on the complementary spaces given by the functions which vanish on $x_{2}=x_{1}$. Repeating this argument shows that the Hochschild homology is the same as the homology of $b$ acting on the smaller subspaces

$$
\begin{gather*}
\left\{q \in \mathbb{C}\left[x_{0}, x_{1}, \ldots, x_{k}\right] ; q\left(\ldots, x_{j-1}, x_{j}, \ldots\right)=0, j=1, \ldots, k\right\}, \\
\quad(b q)\left(x_{0}, x_{1}, \ldots, x_{k-1}\right)=(-1)^{k} q\left(x_{0}, x_{1}, \ldots, x_{k-1}, x_{0}\right) . \tag{10.12}
\end{gather*}
$$

Note that one cannot proceed further directly, in the sense that one cannot reduce to the subspace of functions vanishing on $x_{k}=x_{0}$ as well, since this subspace is not linearly independent of the previous ones ${ }^{2}$

$$
x_{k}-x_{0}=\sum_{j=0}^{k-1}\left(x_{j_{1}}-x_{j}\right) .
$$

It is precisely this 'non-transversality' of the remaining restriction map in ( $\frac{24.94}{10.12}$ ) which remains to be analysed.

Now, let us we make the following change of variable in each of these reduced chain spaces setting

$$
y_{0}=x_{0}, y_{1}=x_{j}-x_{j-1}, \text { for } j=1, \ldots, k .
$$

Then the differential can be written in terms of the pull-back operation

$$
\begin{gathered}
E_{P}: \mathbb{R}^{n k} \hookrightarrow \mathbb{R}^{n(k+1)}, E_{P}\left(y_{0}, y_{1}, \ldots, y_{k-1}\right)=\left(y_{0}, y_{1}, \ldots, y_{k-1},-\sum_{j=1}^{k-1} y_{j}\right), \\
b q=(-1)^{k} E_{P}^{*} q,
\end{gathered}
$$

[^15]The variable $x_{0}=y_{0}$ is a pure parameter, so gan be dropped from the notation (and restored at the end as the factor $\mathbb{C}[x]$ in (ITI.9). Also, as already noted, the degree of a polynomial (in all variables) does not increase under any of these pull-back operations, in fact they all preserve the total degree of homogeneity so it suffices to consider the differential $b$ acting on the spaces of homogeneous polynomials which vanish at the origin in each factor

$$
\begin{gathered}
Q_{k}^{m}=\left\{q \in \mathbb{C}^{m}\left[y_{1}, \ldots, y_{k}\right] ; q(s y)=s^{m} q(y), q\left(y_{1}, \ldots, y_{j-1}, 0, y_{j+1}, \ldots, y_{k}\right)=0\right\} \\
b: Q_{k}^{m} \longrightarrow Q_{k-1}^{m}, b q=(-1)^{*} E_{P}^{*} q
\end{gathered}
$$

To analyse this non-transversality further, let $J_{i} \subset \mathbb{C}\left[y_{1}, \ldots, y_{k}\right]$ be the ideal generated by the $n$ monomials $y_{i}^{l}, l=1, \ldots, n$. Thus, by Taylor's theorem,

$$
J_{i}=\left\{q \in \mathbb{C}\left[y_{1}, \ldots, y_{k}\right] ; q\left(y_{1}, y_{2}, \ldots, y_{j-1}, 0, y_{j}, y_{k}\right)=0\right.
$$

Similary set

$$
J_{P}=\left\{q \in \mathbb{C}\left[y_{1}, \ldots, y_{k}\right] ; q\left(y_{1}, \ldots,-\sum_{j=1}^{k-1} y_{j}\right)=0\right)
$$

For any two ideals $I$ and $J$, let $I \cdot J$ be the span of the products. Thus for these particular ideals an element of the product is a sum of terms each of which has a factor vanishing on the corresponding linear subspace. For each $k$ there are $k+1$ ideals and, by Taylor's theorem, the intersection of any $k$ of them is equal to the span of the product of those $k$ ideals. For the $k$ coordinate ideals this is Taylor's theorem as used in the reduction above. The general case of any $k$ of the ideals can be reduced to this case by linear change of coordinates. The question then, is structure of the intersection of all $k+1$ ideals. The proof of the theorem is therefore completed by the following result.

Lemma 10.3. The intersection $Q_{k}^{m} \cap J_{P}=Q^{m} \cdot J_{P}$ for every $m \neq k$ and

$$
\begin{equation*}
Q_{k}^{k} \cap J_{P}=\Lambda^{k}\left(\mathbb{C}^{n}\right) \tag{10.13}
\end{equation*}
$$

Proof. When $m<k$ the jdeal $Q_{k}^{m}$ vanishes, so the result is trivial.
 degree $k$ in $k$ variables (each in $\mathbb{R}^{n}$ ) which vanishes at the origin in each variable is necessarily linear in each variable, i.e. is just a $k$-multilinear function. Given such a multilinear function $q\left(y_{1}, \ldots, y_{k}\right)$ the condition that $b q=0$ is just that

$$
\begin{equation*}
q\left(y_{1}, \ldots, y_{k-1},-y_{1}-y_{2}-\cdots-y_{k-1}\right) \equiv 0 \tag{10.14}
\end{equation*}
$$

Using the linearity in the last variable the left side can be expanded as a sum of $k-1$ functions each quadratic in one variables $y_{j}$ and linear in the rest. Thus the vanishing of the sum implies the vanishing of each, so

$$
q\left(y_{1}, \ldots, y_{k-1}, y_{j}\right) \equiv 0 \forall j=1, \ldots, k-1
$$

This is the statement that the multlinear function $q$ is antisymmetric between the $j$ th and $k$ th variables for each $j<k$. Since these exchange maps generate the permutation group, $q$ is necessarily totally antisymmetric. This proves the isomorphism ( 10.13 ) since $\Lambda^{k}\left(\mathbb{C}^{n}\right)$ is the space of complex-valued totally antisymmetric $k$-linear forms. ${ }^{3}$

Thus it remains to consider the case $m \geq k+1$. Consider a general element $q \in$ $Q_{k}^{m} \cap J_{P}$. To show that it is in $Q_{k}^{m} \cdot J_{P}$ we manipulate it, working modulo $Q_{k}^{m} \cap J_{P}$,

[^16]and use induction over $k$. Decompose $q$ as a sum of terms $q_{l}$, each homogeneous in the first variable, $y_{1}$, of degree $l$. Since $q$ vanishes at $y_{1}=0$ the first term is $q_{1}$, linear in $y_{1}$. The condition $b q=0$, i.e. $q \in J_{P}$, is again just (ITI.14). Expanding in the last variable shows that the only term in $b q$ which is linear in $y_{1}$ is
$$
q_{1}\left(y_{1}, \ldots, y_{k-1},-y_{2}-\cdots-y_{k-1}\right)
$$

Thus the coefficient of $y_{1, i}$, the $i$ th component of $y_{1}$ in $q_{1}$, is an element of $Q_{k-1}^{m-1}$ which is in the ideal $J_{P}\left(\mathbb{R}^{k-1}\right)$, i.e. for $k-1$ variables. This ideal is generated by the components of $y_{2}+\cdots+y_{k}$. So we can proceed by induction and suppose that the result is true for less than $k$ variables for all degrees of homogeneity. Writing $y_{2}+\cdots+y_{k}=\left(y_{1}+y_{2}+\cdots+y_{k}\right)-y_{1}$ It follows that, modulo $Q_{k}^{m} \cdot J_{P}, q_{1}$ can be replaced by a term of one higher homogeneity in $y_{1}$. Thus we can assume that $q_{i}=0$ for $i<2$. The same argument now applies to $q_{2}$; expanded as a polynomial in $y_{1}$ the coefficients must be elements of $Q_{k-1}^{m-2} \cap J_{P}$. Thus, unless $m-2=k-1$, i.e. $m=k+1$, they are, by the inductive hypothesis, in $Q_{k-1}^{m-2} \cdot J_{P}\left(\mathbb{R}^{k-1}\right)$ and hence, modulo $Q_{k}^{m} \cdot J_{P}, q_{2}$ can be absorbed in $q_{3}$. This argument can be continued to arrange that $q_{i} \equiv 0$ for $i<m-k+1$. In fact $q_{i} \equiv 0$ for $i>m-k+1$ by the assumption that $q \in Q_{k}^{m}$.

Thus we are reduced to the assumption that $q=q_{m-k+1} \in Q_{k}^{m} \cap J_{P}$ is homogeneous of degree $m-k+1$ in the first variable. It follows that it is multilinear in the last $k-1$ variables. The vanishing of $b q$ shows that it is indeed totally antisymmetric in these last $k-1$ variables. Now for each non-zero monomial consider the map $J:\{1,2, \ldots, n\} \longrightarrow \mathbb{N}_{0}$ such that $J(i)$ is the number of times a variable $y_{l, i}$ occurs for some $1 \leq l \leq k$. The decomposition into the sum of terms for each fixed $J$ is preserved by $b$. It follows that we can assume that $q$ has only terms corresponding to a fixed map $J$. If $J(i)>1$ for any $i$ then a factor $y_{1, i}$ must be present in $q$, since it is antisymmetric in the other $k-1$ variables. In this case it can be written $y_{1, i} q^{\prime}$ where $b q^{\prime}=0$. Since $q^{\prime}$ is necessarily in the product of the indeals $J_{2} \ldots J_{k} \cdot J_{P}$ it follows that $q^{\prime} \in Q^{m} \cdot J_{P}$. Thus we may assume that $J(i)=0$ or 1 for all $i$. Since the extra variables now play no rôle we may assume that $n=m$ is the degree of homogeneity and each index $i$ occurs exactly once.

For convenience let us rotate the last $k-1$ variables so the last is moved to the first position. Polarizing $q$ in the first variable, it can be represented uniquely as an $n$-multilinear function on $\mathbb{R}^{n}$ which is symmetric in the first $n-k+1$ variables, totally antisymmetric in the last $k-1$ and has no monomial with repeated index. Let $M_{k-1}(n)$ be the set of such multilinear funtions. The vanishing of $b q$ now corresponds to the vanishing of the symmetrization of $q$ in the first $n-k+2$ variables. By the antisymmetry in the second group of variables this gives a complex

$$
M_{n}(n) \xrightarrow{b_{n}} M_{n-1}(n)^{b_{n-1}} \cdots \quad \xrightarrow{b_{2}} M_{1}(n) \xrightarrow{b_{1}} M_{0} \xrightarrow{b_{0}} 0
$$

The remaining step is to show that this is exact.
Observe that $\operatorname{dim}\left(M_{k}(n)\right)=\binom{n}{k}$ since there is a basis of $M_{k}(n)$ with elements labelled by the subsets $I \subset\{1, \ldots, n\}$ with $k$ elements. Indeed let $\omega$ be a nontrivial $k$-multilinear function of $k$ variables and let $\omega_{I}$ be this function on $\mathbb{R}^{k} \subset \mathbb{R}^{n}$ identified as the set of variables indexed by $I$. Then if $a \in M_{0}(n-k)$ is a basis of this 1-dimensional space and $a_{I}$ is this function on the complementary $\mathbb{R}^{n-k}$ the
tensor products $a_{I} \omega_{I}$ give a basis. Thus there is an isomorphism

$$
M_{k} \ni q=\sum_{I \subset\{1, \ldots, n\},|I|=k} c_{I} a_{I} \otimes \omega_{I} \longmapsto \sum_{I \subset\{1, \ldots, n\},|I|=k} c_{I} \otimes \omega_{I} \in \Lambda^{k}\left(\mathbb{R}^{n}\right) .
$$

Transfered to the exterior algebra by this isomorphism the differential $b$ is just contraction with the vector $e_{1}+e_{2}+\cdots+e_{n}$ (in the first slot). A linear transformation reducing this vector to $e_{1}$ shows immediately that this (Koszul) complex is exact, with the null space of $b_{k}$ on $\Lambda^{k}\left(\mathbb{R}^{n}\right)$ being spanned by those $\omega_{I}$ with $1 \in I$ and the range of $b_{k+1}$ spanned by those with $1 \notin I$. The exactness of this complex completes the proof of the lemma.

### 10.3. Hochschild homology of $\mathcal{C}^{\infty}(X)$

The first example of Hochschild homology that we shall examine is for the commutative algebra $\mathcal{C}^{\infty}(X)$ where $X$ is any $\mathcal{C}^{\infty}$ manifold (compact or not). As noted above we need to replace the tensor product by some completion. In the present case observe that for any two manifolds $X$ and $Y$

$$
\begin{equation*}
\mathcal{C}^{\infty}(X) \otimes \mathcal{C}^{\infty}(Y) \subset \mathcal{C}^{\infty}(X \times Y) \tag{10.15}
\end{equation*}
$$

is dense in the $\mathcal{C}^{\infty}$ topology. Thus we simply declare the space of $k$-chains for Hochschild homology to be $\mathcal{C}^{\infty}\left(X^{k+1}\right)$, which can be viewed as a natural completion ${ }^{4}$ of $\mathcal{C}^{\infty}(X)^{\otimes(k+1)}$. Notice that the product of two functions can be written in terms of the tensor product as

$$
\begin{equation*}
a \cdot b=D^{*}(a \otimes b), a, b \in \mathcal{C}^{\infty}(X), D: X \ni z \longmapsto(z, z) \in X^{2} \tag{10.16}
\end{equation*}
$$

The variables in $X^{k+1}$ will generally be denoted $z_{0}, z_{1}, \ldots, z_{k}$. Consider the 'diagonal' submanifolds

$$
\begin{equation*}
D_{i, j}=\left\{\left(z_{0}, z_{1}, \ldots, z_{k}\right) ; z_{i}=z_{j}\right\}, i, j=0, \ldots, m, i \neq j \tag{10.17}
\end{equation*}
$$

We shall use the same notation for the natural embedding of $X^{k}$ as each of these submanifolds, at least for $j=i+1$ and $i=0, j=m$,

$$
\begin{gathered}
D_{i, i+1}\left(x_{0}, \ldots, z_{m-1}\right)=\left(z_{0}, \ldots, z_{i}, z_{i}, z_{i+1}, \ldots, z_{m-1}\right) \in D_{i, i+1}, i=0, \ldots, m-1 \\
D_{m, 0}\left(z_{0}, \ldots, z_{m-1}\right)=\left(z_{0}, \ldots, z_{m-1}, z_{0}\right)
\end{gathered}
$$

Then the action of $b^{\prime}$ and $b$ on the tensor products, and hence on all chains, can be written

$$
\begin{equation*}
b^{\prime} \alpha=\sum_{i=0}^{m-1}(-1)^{i} D_{i, i+1}^{*} \alpha, b \alpha=b^{\prime} \alpha+(-1)^{m} D_{m, 0}^{*} \alpha \tag{10.18}
\end{equation*}
$$

[^17]ThEOREM 10.2. The differential $b^{\prime}$ is acyclic and the homology ${ }^{5}$ of the complex

$$
\begin{equation*}
\ldots \stackrel{b}{\longrightarrow} \mathcal{C}^{\infty}\left(X^{k+1}\right) \xrightarrow{b} \mathcal{C}^{\infty}\left(X^{k}\right) \xrightarrow{b} \ldots \tag{10.19}
\end{equation*}
$$

is naturally isomorphic to $\mathcal{C}^{\infty}\left(X ; \Lambda^{*}\right)$.
Before proceeding to the proof proper we note two simple lemmas.
HL.ciX LEMMA 10.4. ${ }^{6}$ For any $j=0, \ldots, m-1$, each function $\alpha \in \mathcal{C}^{\infty}\left(X^{k+1}\right)$ which vanishes on $D_{i, i+1}$ for each $i \leq j$ can be written uniquely in the form

$$
\alpha=\alpha^{\prime}+\alpha^{\prime \prime}, \alpha^{\prime}, \alpha^{\prime \prime} \in \mathcal{C}^{\infty}\left(X^{k+1}\right)
$$

where $\alpha^{\prime \prime}$ vanishes on $D_{i, i+1}$ for all $i \leq j+1$ and $\alpha^{\prime}$ is independent of $z_{j+1}$.
Proof. Set $\alpha^{\prime}=\pi_{j+1}^{*}\left(D_{j, j+1}^{*} \alpha\right)$ where $\pi_{j}: X^{k+1} \longrightarrow X^{k}$ is projection off the $j$ th factor. Thus, essentially by definition, $\alpha^{\prime}$ is independent of $z_{j+1}$. Moreover, $\pi_{j+1} D_{j, j+1}=$ Id so $D_{j, j+1}^{*} \alpha^{\prime}=D_{j, j+1}^{*} \alpha$ and hence $D_{j, j+1}^{*} \alpha^{\prime \prime}=0$. The decomposition is clearly unique, and for $i<j$,

$$
\begin{equation*}
D_{j, j+1} \circ \pi_{j+1} \circ D_{i, i+1}=D_{i, i+1} \circ F_{i, j} \tag{10.20}
\end{equation*}
$$

for a smooth map $F_{i, j}$, so $\alpha^{\prime}$ vanishes on $D_{i, i+1}$ if $\alpha$ vanishes there.
Lemma 10.5. For any finite dimensional vector space, $V$, the $k$-fold exterior power of the dual, $\Lambda^{k} V^{*}$, can be naturally identified with the space of functions

$$
\begin{gather*}
\left\{u \in \mathcal{C}^{\infty}\left(V^{k}\right) ; u(s v)=s^{k} v, s \geq 0, u \upharpoonright\left(V^{i} \times\{0\} \times V^{k-i-1}\right)=0 \text { for } i=0, \ldots, k-1\right.  \tag{10.21}\\
\text { and } \left.u \upharpoonright G=0, G=\left\{\left(v_{1}, \ldots, v_{k}\right) \in V^{k} ; v_{1}+\cdots+v_{k}=0\right\}\right\}
\end{gather*}
$$

Proof. The homogeneity of the smooth function, $u$, on $V^{k}$ implies that it is a homogeneous polynomial of degree $k$. The fact that it vanishes at 0 in each variable then implies that it is multlinear, i.e. is linear in each variable. The vanishing on $G$ implies that for any $j$ and any $v_{i} \in V, i \neq j$,

$$
\begin{equation*}
\sum_{i \neq j} u\left(v_{1}, \ldots, v_{j-1}, v_{i}, v_{j+1}, \ldots, v_{k}\right)=0 \tag{10.22}
\end{equation*}
$$

Since each of these terms is quadratic (and homogeneous) in the corresponding variable $v_{i}$, they must each vanish identically. Thus, $u$ vanishes on $v_{i}=v_{j}$ for each $i \neq j$; it is therefore totally antisymmetric as a multlinear form, i.e. is an element of $\Lambda^{k} V^{*}$. The converse is immediate, so the lemma is proved.

Proof $_{24 . \mathrm{FF}_{9} \text { THEOREM }}^{\frac{H H . c i X}{10.2}}$. The H-unitality ${ }^{7}$ of $\mathcal{C}^{\infty}(X)$ follows from the proof of Lemma $\frac{24.61}{10.61}$ which carries over verbatim to the larger chain spaces.

By definition the Hochschild homology in degree $k$ is the quotient

$$
\begin{equation*}
\operatorname{HH}_{k}\left(\mathcal{C}^{\infty}(X)\right)=\left\{u \in \mathcal{C}^{\infty}\left(X^{k+1}\right) ; b u=0\right\} / b \mathcal{C}^{\infty}\left(X^{k+2}\right) \tag{10.23}
\end{equation*}
$$

The first stage in identifying this quotient is to apply Lemma $\frac{\text { HL.ciX }}{10.4 \text { repeatedly. Let }}$ us carry through the first step separately, and then do the general case.

[^18]For $j=0$, consider the decomposition of $u \in \mathcal{C}^{\infty}\left(X^{k+1}\right)$ given by Lemma $\frac{\text { HL.ciX }}{10.4,}$ thus
$u=u_{0}+u_{(1)}, u_{0} \in \pi_{1}^{*} \mathcal{C}^{\infty}\left(X^{k}\right), u_{(1)} \in J_{1}^{(k)}=\left\{u \in \mathcal{C}^{\infty}\left(X^{k+1}\right) ; u \upharpoonright D_{0,1}=0\right\}$.
Now each of these subspaces of $\mathcal{C}^{\infty}\left(X^{k+1}\right)$ is mapped into the corresponding subspace of $\mathcal{C}^{\infty}\left(X^{k}\right)$ by $b$; i.e. they define subcomplexes. Indeed,

$$
\begin{aligned}
& u \in \pi_{1}^{*} \mathcal{C}^{\infty}\left(X^{k}\right) \Longrightarrow D_{0,1}^{*} u=D_{1,2}^{*} u \text { so } \\
& u=\pi_{1}^{*} v \Longrightarrow b u=\pi_{1}^{*} B v, B^{*} v=-\sum_{i=1}^{k-1}(-1)^{i} D_{i, i+1}^{*} u+(-1)^{k} D_{k-1,0}^{*} v
\end{aligned}
$$

For the other term

$$
\begin{equation*}
b u_{(1)}=\sum_{i=1}^{k-1}(-1)^{i} D_{i, i+1}^{*} u_{(1)}+(-1)^{k} D_{k, 0}^{*} u_{(1)} \Longrightarrow b u_{(1)} \in J_{1}^{(k-1)} \tag{10.25}
\end{equation*}
$$

Thus, $b u=0$ is equivalent to $b u_{0}=0$ and $b u_{(1)}=0$. From (24.35), defining an isomorphism by

$$
\begin{equation*}
E_{(k-1)}: \mathcal{C}^{\infty}\left(X^{k}\right) \longrightarrow \mathcal{C}^{\infty}\left(X^{k}\right), E_{(k-1)} v\left(z_{1}, \ldots, z_{k}\right)=v\left(z_{2}, \ldots, z_{k}, z_{1}\right) \tag{10.26}
\end{equation*}
$$

it follows that

$$
\begin{equation*}
B=-E_{(k-1)}^{-1} b^{\prime} E_{(k-1)} \tag{10.27}
\end{equation*}
$$

is conjugate to $b^{\prime}$. Thus $B$ is acyclic so in terms of ( $\left(\frac{24.37}{10.24}\right)$

$$
\begin{equation*}
b u=0 \Longrightarrow u-u_{(1)}=b w, w=\pi_{1}^{*} v^{\prime} \tag{10.28}
\end{equation*}
$$

As already noted this is the first ftep in an inductive procedure, the induction being over $1 \leq j \leq k$ in Lemma $\frac{1 L 10.4 .}{10.4}$. Thus we show inductively that

### 24.38

(10.29) $b u=0 \Longrightarrow u-u_{(j)}=b w$,

$$
u_{(j)} \in J_{j}^{(k)}=\left\{u \in \mathcal{C}^{\infty}\left(X^{k+1}\right) ; u \upharpoonright D_{i, i+1}=0,0 \leq i \leq j-1\right\}
$$

For $j=1$ this is ( $\frac{24.39}{10.28)}$. Proceeding inductively we may suppose that $u=u_{(j)}$ and take the decomposition of Lemma $i \mathbf{i 0 . 4 , \text { so }}$

$$
\begin{equation*}
u_{(j)}=u^{\prime}+u_{(j+1)}, u_{(j+1)} \in J_{j+1}^{(k)}, u^{\prime}=\pi_{j+1}^{*} v \in J_{j}^{(k)} \tag{10.30}
\end{equation*}
$$

Then, as before, $b u_{(j)}=0$ implies that $b u^{\prime}=0$. Furthermore, acting on the space $\pi_{j+1}^{*} \mathcal{C}^{\infty}\left(X^{k}\right) \cap J_{(j)}^{k}, b$ is conjugate to $b^{\prime}$ acting in $k+1-j$ variables. Thus, it is again acyclic, so $u_{(j)}$ and $u_{(j+1)}$ are homologous as Hochschild $k$-cycles.

The end point of this inductive procedure is that each $k$-cycle is homologous to an element of

$$
\begin{equation*}
J^{(k)}=J_{k}^{(k)}=\left\{u \in \mathcal{C}^{\infty}\left(X^{k+1}\right) ; D_{i, i+1}^{*} u=0, i \leq i \leq k-1\right\} \tag{10.31}
\end{equation*}
$$

Acting on this space $b u=(-1)^{k} D_{k, 0}^{*} u$, so we have shown that
24.43 (10.32)

$$
\operatorname{HH}_{k}\left(\mathcal{C}^{\infty}(X)\right)=M^{(k)} /\left(M^{(k)} \cap b \mathcal{C}^{\infty}\left(X^{k+1}\right)\right), M^{(k)}=\left\{u \in J^{(k)} ; D_{k, o}^{*} u=0\right\}
$$

Now consider the subspace

$$
\begin{align*}
\tilde{M}^{(k)} & =\left\{u \in \mathcal{C}^{\infty}\left(X^{k+1}\right) ;\right.  \tag{10.33}\\
u & \left.=\sum_{\text {finite, }, 0 \leq j \leq k-1}\left(f\left(z_{j}\right)-f\left(z_{j+1}\right)\right) u_{f, j}, u_{f, j} \in M^{(k)}, f \in \mathcal{C}^{\infty}(X) .\right\} .
\end{align*}
$$

If $u=\left(f\left(z_{j}\right)-f\left(z_{j+1}\right)\right) v$, with $v \in M^{(k)}$ set

$$
\begin{align*}
& w\left(z_{0}, z_{1}, \ldots, z_{j}, z_{j+1}, z_{j+2}, \ldots, z_{k+1}\right)  \tag{10.34}\\
& \quad=(-1)^{j}\left(f\left(z_{j}\right)-f\left(z_{j+1}\right)\right) v\left(z_{0}, \ldots, z_{j}, z_{j+2}, z_{j+3}, \ldots, z_{k}\right)
\end{align*}
$$

Then, using the assumed vanishing of $v, b w=u .^{8}$ Thus all the elements of $\tilde{M}^{(k)}$ are exact.

Let us next compute the quotient $M^{(k)} / \tilde{M}^{(k)}$. Linearizing in each factor of $X$ around the submanifold $z_{0}=z_{1}=\cdots=z_{k}$ in $V^{k}$ defines a map

$$
\begin{equation*}
\mu: M^{(k)} \ni u \longrightarrow u^{\prime} \in \mathcal{C}^{\infty}\left(X ; T X \otimes \cdots \otimes T^{*} X\right) \tag{10.35}
\end{equation*}
$$

The map is defined by taking the term of homogeneity $k$ in a normal expansion around the submanifold. The range space is therefore precisely the space of sections of the $k$-fold tensor product bundle which vanish on the subbundle defined in each fibre by $v_{1}+\cdots+v_{k}=0$. Thus, by Lemma $\frac{24.33}{10.5, ~} \mu$ actually defines a sequence

$$
\begin{equation*}
0 \longrightarrow \tilde{M}^{(k)} \hookrightarrow M^{(k)} \xrightarrow{\mu} \mathcal{C}^{\infty}\left(X ; \Lambda^{k} X\right) \longrightarrow 0 . \tag{10.36}
\end{equation*}
$$

Lemma 10.6. For any $k,\left(\frac{24.48}{10.36}\right)$ is a short exact sequence.
Proof. So far I have a rather nasty proof by induction of this result, there should be a reasonably elementary argument. Any offers?

From this the desired identification, induced by $\mu$,

$$
\begin{equation*}
\operatorname{HH}_{k}\left(\mathcal{C}^{\infty}(X)\right)=\mathcal{C}^{\infty}\left(X ; \Lambda^{k} X\right) \tag{10.37}
\end{equation*}
$$

follows, once is is shown that no element $u \in M^{(k)}$ with $\mu(u) \neq 0$ can be exact. This follows by a similar argument. Namely if $u \in M_{M L . c I X}^{(k)}$ is exact then write $u=b v$, $v \in \mathcal{C}^{\infty}\left(X^{k}\right)$ and apply the decomposition of Lemma $\frac{10.4 \text { to get } v=v_{0}+v_{(1)} \text {. Since }}{}$ $u=0$ on $D_{1,0}$ it follows that $b v_{0}=0$ and hence $u=b v_{1)}$. Proceeding inductively we conculde that $u=b v$ with $v \in M^{(k+1)}$. Now, $\mu(b v)=0$ by inspection.

### 10.4. Commutative formal symbol algebra

As a first step towards the computation of the Hochschild homology of the algebra $\mathcal{A}=\Psi^{\mathbb{Z}}(X) / \Psi^{-\infty}(X)$ we consider the formal algebra of symbols with commutative product. Thus,

$$
\begin{equation*}
\mathcal{A}=\left\{\left(a_{j}\right)_{j=-\infty}^{\infty} ; a_{j} \in \mathcal{C}^{\infty}\left(S^{*} X ; P^{(j)}\right), a_{j}=0 \text { for } j \gg 0\right\} \tag{10.38}
\end{equation*}
$$

Here $P^{(k)}$ is the line bundle over $S^{*} X$ with sections consisting of the homogeneous functions of degree $k$ on $T^{*} X \backslash 0$. The multiplication is as functions on $T^{*} X \backslash 0$, so

$$
\left(a_{j}\right) \cdot\left(b_{j}\right)=\left(c_{j}\right), c_{j}=\sum_{k=-\infty}^{\infty} a_{j-k} b_{k}
$$

[^19]using the fact that $P^{(l)} \otimes P^{(k)} \equiv P^{(l+k)}$. We take the completion of the tensor product to be
\[

$$
\begin{align*}
& \mathcal{B}^{(k)}=\left\{u \in \mathcal{C}^{\infty}\left(\left(T^{*} X \backslash 0\right)^{k+1}\right) ; u=\sum_{\text {finite }} u_{I}\right.  \tag{10.39}\\
&\left.u_{I} \in \mathcal{C}^{\infty}\left(S^{*} X ; P^{\left(I_{0}\right)} \otimes P^{\left(I_{1}\right)} \otimes \cdots \otimes P^{\left(I_{k}\right)}\right),|I|=k\right\}
\end{align*}
$$
\]

That is, an element of $\mathcal{B}^{(k)}$ is a finite sum of functions on the $(k+1)$-fold product of $T^{*} X \backslash 0$ which are homogeneous of degree $I_{j}$ on the $j$ th factor, with the sum of the homogeneities being $k$. Then the Hochschild homology is the cohomology of the subcomplex of the complex for $\mathcal{C}^{\infty}\left(T^{*} X\right)$

$$
\begin{equation*}
\cdots \xrightarrow{b} \mathcal{B}^{(k)} \xrightarrow{b} \mathcal{B}^{(k-1)} \xrightarrow{b} \cdots \tag{10.40}
\end{equation*}
$$

THEOREM 10.3. The cohomology of the complex (10.54) for the commutative product on $\mathcal{A}$ is
(10.41) $\operatorname{HH}_{k}(\mathcal{A}) \equiv\left\{\alpha \in \mathcal{C}^{\infty}\left(T^{*} X \backslash 0 ; \Lambda^{k}\left(T^{*} X\right) ; \alpha\right.\right.$ is homogeneous of degree $\left.k\right\}$.

### 10.5. Hochschild chains

The completion of the tensor product that we take to define the Hochschild homology of the 'full symbol algebra' is the same space as in ( $(\mathbb{1 0 . 3 9})$ but with the non-commutative product derived from the quantization map for some Riemann metric on $X$. Since the product is given as a formal sum of bilinear differential operators it can be take to act on an pair of factors.

$$
\begin{equation*}
\ldots \xrightarrow{b^{(*)}} \mathcal{B}^{(k)} \xrightarrow{b^{(\star)}} \mathcal{B}^{(k-1)} \xrightarrow{b^{(\star)}} \ldots \tag{10.42}
\end{equation*}
$$

The next, and major, task of this chapter is to describe the cohomology of this complex.

THEOREM 10.4. The Hochschild homolgy of the algebra, $\Psi_{\mathrm{phg}}^{\mathbb{Z}}(X) / \Psi_{\mathrm{phg}}^{-\infty}(X)$, of formal symbols of pseudodifferential operators of integral order, identified as the cohomology of the complex ( $(\mathbb{I T} .42)$, is naturally identified with two copies of the cohomology of $S^{*} X^{9}$

$$
\begin{equation*}
\operatorname{HH}_{k}(\mathcal{A} ; \circ) \equiv H^{2 n-k}\left(S^{*} X\right) \oplus H^{2 n-1-k}\left(S^{*} X\right) \tag{10.43}
\end{equation*}
$$

### 10.6. Semi-classical limit and spectral sequence

The 'classical limit' in physics, especially quantuum mechanics, is the limit in which physical variables become commutative, i.e. the non-commutative coupling between position and momentum variables vanishes in the limit. Formally this typically involves the replacement of Planck's constant by a parameter $h \rightarrow 0$. A phenomenon is 'semi-classical' if it can be understood at least in Taylor series in this parameter. In this sense the Hochschild homology of the full symbol algebra is semi-classical and (following $\frac{\operatorname{sry} 1 \text { inski-Getzier }}{[2]) \text { this is how we shall compute it. }}$

The parameter $h$ is introduced directly as an isomorphism of the space $\mathcal{A}$

$$
L_{h}: \mathcal{A} \longrightarrow \mathcal{A}, L_{h}\left(a_{j}\right)_{j=-\infty}^{*}=\left(h^{j} a_{j}\right)_{j=-\infty}^{*}, h>0
$$

[^20]Clearly $L_{h} \circ L_{h^{\prime}}=L_{h h^{\prime}}$. For $h \neq 1, L_{h}$ is not an algebra morphism, so induces a 1-parameter family of products

$$
\begin{equation*}
\alpha \star_{h} \beta=\left(L_{h}^{-1}\right)\left(L_{h} \alpha \star L_{h} \beta\right) . \tag{10.44}
\end{equation*}
$$

In terms of the differential operators, associated to quantization by a particular choice of Riemann metric on $X$ this product can be written

$$
\begin{equation*}
\alpha \star_{h} \beta=\left(c_{j}\right)_{j=-\infty}^{*}, c_{j}=\sum_{k=0}^{*} \sum_{l=-*}^{*} h^{k} P_{k}\left(a_{j-l-k}, b_{l}\right) . \tag{10.45}
\end{equation*}
$$

It is important to note here that the $P_{k}$, as differential operators on functions on $T^{*} X$, do only depend on $k$, which is the difference of homogeneity between the product $a_{j-l+k} b_{l}$, which has degree $j+k$ and $c_{j}$, which has degree $j$.

Since $\mathcal{A}$ with product $\star_{h}$ is a 1-parameter family of algebras, i.e. a deformation of the algebra $\mathcal{A}$ with product $\star=\star_{1}$, the Hochschild homology is 'constant' in $h$. More precisely the map $L_{h}$ induces a canonical isomorphism

$$
L_{h}^{*}: \operatorname{HH}_{k}\left(\mathcal{A} ; \star_{h}\right) \equiv \operatorname{HH}_{k}(\mathcal{A} ; \star)
$$

The dependence of the product on $h$ is smooth, so it is reasonable to expect the cycles to have smooth representatives as $h \rightarrow 0$. To investigate the consider Taylor series in $h$ and define

$$
\begin{gather*}
F_{p, k}=\left\{\alpha \in \mathcal{B}^{(k)} ; \exists \alpha(h) \in \mathcal{C}^{\infty}\left([0,1)_{h} ; \mathcal{B}^{(k)}\right) \text { with } \alpha(0)=\alpha\right. \text { and }  \tag{10.46}\\
b_{h} \alpha \in h^{p} \mathcal{C}^{\infty}\left([0,1)_{h} ; \mathcal{B}^{k-1)}\right\}, \\
G_{p, k}=\left\{\alpha \in \mathcal{B}^{(k)} ; \exists \beta(h) \in \mathcal{C}^{\infty}\left([0,1)_{h} ; \mathcal{B}^{(k+1)}\right)\right. \text { with } \\
b_{h} \beta(h) \in h^{p-1} \mathcal{C}^{\infty}\left([0,1)_{h} ; \mathcal{B}^{(k)} \text { and }\left(t^{-p+1} b_{h} \beta\right)(0)=\alpha\right\} .
\end{gather*}
$$

Here $b_{h}$ is the differential defined by the product $\star_{h}$.
Notice that the $F_{p, k}$ decrease with increasing $p$, since the condition becomes stronger, while $G_{p, k}$ increases with $p$, the condition becoming weaker. ${ }^{10}$ We define the 'spectral sequence' corresponding to this filtration by

$$
E_{p, k}=F_{p, k} / G_{p, k}
$$

These can also be defined successively, in the sense that if

$$
\begin{aligned}
& F_{p, k}^{\prime}=\left\{u \in E_{p-1, k} ; u=\left[u^{\prime}\right], u^{\prime} \in F_{p, k}\right\} \\
&\left.G_{p, k}^{\prime}=\left\{e \in E_{p-1, k} ; u=\left[u^{\prime}\right],\right] u^{\prime} \in G_{p, k}\right\} \\
& \text { then } E_{p, k} \equiv F_{p, k}^{\prime} / G_{p, k}^{\prime}
\end{aligned}
$$

The basic idea ${ }^{11}$ of a spectral sequence is that each $E_{p}=\bigoplus_{k} E_{p, k}$, has defined on it a differential such that the next spaces, forming $E_{p+1}$, are the cohomology space for the complex. This is easily seen from the definitions of $F_{p, k}$ as follows. If $\alpha \in F_{p, k}$ let $\beta(t)$ be a 1-parameter family of chains as in the defintion. Then consider

$$
\begin{equation*}
\gamma\left(t^{-p} b_{h} \beta\right)(0) \in \mathcal{B}^{(k-1)} \tag{10.48}
\end{equation*}
$$

[^21]This depends on the choice of $\beta$, but only up to a term in $G_{p, k-1}$. Indeed, let $\beta^{\prime}(t)$ is another choice of extension of $\alpha$ satisfying the condition that $b_{h} \beta^{\prime} \in$ $h^{p} \mathcal{C}^{\infty}\left([0,1) ; \mathcal{B}^{(k-1)}\right.$ and let $\gamma^{\prime}$ be defined by (10.48) with $\beta$ replaced by $\beta^{\prime}$. Then $\delta(t)=t^{-1}\left(\beta(h)-\beta^{\prime}(h)\right)$ satisfies the requirements in the definition of $G_{p, k-1}$, i.e. the difference $\gamma^{\prime}-\gamma \in G_{p, k-1}$. Similarly, if $\alpha \in G_{p, k}$ then $\gamma \in G_{p, k} .{ }^{12}$ The map so defined is a differential

$$
b_{(p)}: E_{p, k} \longrightarrow E_{p, k-1}, b_{(p)}^{2}=0
$$

This follows from the fact that if $\mu=b_{(p)} \alpha$ then, by definition, $\mu=\left(t^{-p} b_{h} \beta\right)(0)$, where $\alpha=\beta(0)$. Taking $\lambda(t)=t^{-p} b_{h} \beta(t)$ as the extension of $\mu$ it follows that $b_{h} \lambda=0$, so $b_{(p)} \mu=0$.

Now, it follows directly from the definition that $F_{0, k}=E_{0, k}=\mathcal{B}^{(k)}$ since $G_{0, k}=\{0\}$. Furthermore, the differential $b_{(0)}$ induced on $E_{0}$ is just the Hochschild differential for the limiting product, $\star_{0}$, which is the commutative product on the algebra. Thus, Theorem $\frac{10.3 \text { just states that }}{}$

$$
E_{1, k}=\bigoplus_{k=-\infty}^{*}\left\{u \in \mathcal{C}^{\infty}\left(T^{*} X \backslash 0 ; \Lambda^{k}\right) ; u \text { is homogeneous of degree } k\right\}
$$

To complete the proof of Theorem $\frac{24.56}{10.4}$ t therefore suffices to show that

$$
\begin{gather*}
E_{2, k} \equiv H^{2 n-k}\left(S^{*} X\right) \oplus H^{2 n-1-k}\left(S^{*} X\right)  \tag{10.49}\\
E_{p, k}=E_{2, k}, \forall p \geq 2, \text { and }  \tag{10.50}\\
\operatorname{HH}_{k}\left(\Psi_{\mathrm{phg}}^{\mathbb{Z}}(X) / \Psi_{\mathrm{phg}}^{-\infty}(X)\right)=\lim _{p \rightarrow \infty} E_{p, k} \tag{10.51}
\end{gather*}
$$

The second and third of these results are usually described, respectively, as the 'degeneration' of the spectral sequence (in this case at the ' $E_{2}$ term') and the 'convergence' of the spectral sequence to the desired cohomology space.

### 10.7. The $E_{2}$ term

As already noted, the $E_{1, k}$ term in the spectral sequence consists of the formal sums of $k$-forms, on $T^{*} X \backslash 0$, which are homogeneous under the $\mathbb{R}^{+}$action. The $E_{2}$ term is the cohomology of the complex formed by these spaces with the differential $b_{(1)}$, which we proceed to compute. For simplicity of notation, consider the formal tensor prodoct rather than its completion. As already noted, for any $\alpha \in \mathcal{B}^{(k)}$ the function $b_{h} \alpha$ is smooth in $h$ and from the definition of $b$,

$$
\begin{align*}
\frac{d}{d h} b_{h} \alpha(0)= & \sum_{i=0}^{k-1}(-1)^{i} a_{0} \otimes \cdots \otimes a_{i-1} \otimes P_{1}\left(a_{i+1}, a_{i}\right) \otimes a_{i+2} \otimes \cdots \otimes a_{k}  \tag{10.52}\\
& +(-1)^{k} P_{1}\left(a_{0}, a_{k}\right) \otimes a_{1} \otimes \cdots \otimes a_{k-1}, \quad \alpha=a_{0} \otimes \cdots \otimes a_{k}
\end{align*}
$$

The general case is only more difficult to write, not different. ${ }^{13}$ This certainly determines $b_{1} \alpha$ if $\alpha$ is a superposition of such terms with $b_{0} \alpha=0$. Although ( 110.52 ) is explicit, it is not given directly in terms of the representation of $\alpha$, assumed to satisfy $b_{0} \alpha=0$ as a form on $T^{*} X \backslash 0$.

[^22]To get such an explicit formula we shall use the symplectic analogue of the Hodge isomorphism. Recall that in any local coordinates on $X, x_{i}, i=1, \ldots, n$, induce local coordinates $x_{i}, \xi_{i}$ in the part of $T^{*} X$ lying above the coordinate patch. In these canonical coordinates the symplectic form (which determines the Poisson bracket) is given by

$$
\begin{equation*}
\omega=\sum_{k=1}^{n} d \xi_{k} \wedge d x_{k} \tag{10.53}
\end{equation*}
$$

This 2 -form is non-degenerate, i.e. the $n$-fold wedge product $\omega^{n} \neq 0$. In fact this volume form fixes an orientation on $T^{*} X$. The symplectic form can be viewed as a non-degenerate antisymmetric bilinear form on $T_{q}\left(T^{*} X\right)$ at each point $q \in T^{*} X$, and hence by duality as a bilineear form on $T_{q}^{*}\left(T^{*} X\right)$. We denote this form in the same way as the Poisson bracket, since with the convention

$$
\{a, b\}(q)=\{d a, d b\}_{q}
$$

they are indeed the same. As a non-degenerate bilinear form on $T^{*} Y, Y=T^{*} X$ this also induces a bilinear form on the tensor algebra, by setting

$$
\left\{e_{1} \otimes \cdots \otimes e_{k}, f_{1} \otimes \cdots \otimes f_{k},\right\}=\prod_{j}\left\{e_{j}, f_{j}\right\}
$$

These bilinear forms are all antisymmetric and non-degenerate and restrict to be non-degnerate on the antisymmetric part, $\Lambda^{k} Y$, of the tensor algebra. Thus each of the form bundles has a bilinear form defined on it, so there is a natural isomorhism
24.72 (10.54) $W_{\omega}: \Lambda_{q}^{k} Y \longrightarrow \Lambda_{q}^{2 n-k} Y, \alpha \wedge W_{\omega} \beta=\{\alpha, \beta\} \omega^{n}, \alpha, \beta \in \mathcal{C}^{\infty}\left(Y, \Lambda^{k} Y\right)$,
for each $k$.
24.74 Lemma 10.7. In canonical coordinates, as in (14.73), consider the basis of $k$ forms given by all increasing subsequences of length $k$,

$$
I:\{1,2, \ldots, k\} \longrightarrow\{1,2, \ldots, 2 n\}
$$

and setting
(10.55) $\quad \alpha_{I}=d z_{I(1)} \wedge d z_{I(2)} \wedge \cdots \wedge d z_{I(k)}$,

$$
\left(z_{1}, z_{2}, \ldots, z_{2 n}\right)=\left(x_{1}, \xi_{1}, x_{2}, \xi_{2}, \ldots, x_{n}, \xi_{n}\right)
$$

In terms of this ordering of the coordinates

$$
\begin{equation*}
W_{\omega}\left(\alpha_{I}\right)=(-1)^{N(I)} \alpha_{W(I)} \tag{10.56}
\end{equation*}
$$

where $W(I)$ is obtained from $I$ by considering each pair $(2 p-1,2 p)$ for $p=1, \ldots, n$, erasing it if it occurs in the image of $I$, inserting it into $I$ if neither $2 p-1$ nor $2 p$ occurs in the range of $I$ and if exactly one of $2 p-1$ and $2 p$ occurs then leaving it unchanged; $N(I)$ is the number of times $2 p$ appears in the range of $I$ without $2 p-1$.

Proof. The Poisson bracket pairing gives, on 1-forms,

$$
-\left\{d x_{j}, d \xi_{j}\right\}=1=\left\{d \xi_{j}, d x_{j}\right\}
$$

with all other pairings zero. Extending this to $k$-forms gives

$$
\begin{array}{r}
\left\{\alpha_{I}, \alpha_{J}\right\}=0 \text { unless }(I(j), J(j))=(2 p-1,2 p) \text { or }(2 p, 2 p-1) \forall j \text { and } \\
\left\{\alpha_{I}, \alpha_{J}\right\}=(-1)^{N}, \text { if }(I(j), J(j))=(2 p-1,2 p) \text { for } N \text { values of } j \\
\quad \text { and }(I(j), J(j))=(2 p-1,2 p) \text { for } N-k \text { values of } j .
\end{array}
$$

From this, and (24.72), (10.54.75 $(10.56)$ follows.
From this proof if also follows that $N(W(I))=N(I)$, so $W_{\omega}^{2}=$ Id. We shall let

$$
\begin{equation*}
\delta_{\omega}=W_{\omega} \circ d \circ W_{\omega} \tag{10.57}
\end{equation*}
$$

denote the differential operator obtained from $d$ by conjugation,

$$
\delta_{\omega}: \mathcal{C}^{\infty}\left(T^{*} X \backslash 0 ; \Lambda^{k}\right) \longrightarrow \mathcal{C}^{\infty}\left(T^{*} X \backslash 0, \Lambda^{k-1}\right)
$$

By construction $\delta_{\omega}^{2}=0$. The exterior algebra of a symplectic manifold with this differential is called the Koszul complex. ${ }^{14}$ All the $\alpha_{I}$ are closed so

$$
\begin{gather*}
\delta_{\omega}\left(a \alpha_{I}\right)=W_{\omega}\left(\sum_{j} \frac{\partial a}{\partial z_{j}} d z_{j}\right) \wedge(-1)^{N(I)} \alpha_{W(I)} \\
=\sum_{j} \frac{\partial a}{\partial z_{j}}(-1)^{N(I)} W_{\omega}\left(d z_{j} \wedge \alpha_{W(I)}\right), \tag{10.58}
\end{gather*}
$$

Observe that ${ }^{15}$

$$
\begin{aligned}
& W_{\omega}\left(d z_{2 p-1} \wedge \alpha_{W(I)}\right)=\iota_{\partial / \partial z_{2 p}} \alpha_{I} \\
& W_{\omega}\left(d z_{2 p} \wedge \alpha_{W(I)}\right)=\iota_{\partial / \partial z_{2 p-1}} \alpha_{I}
\end{aligned}
$$

where, $\iota_{v}$ denotes contraction with the vector field $v$. We therefore deduce the following formula for the action of the Koszul differential

$$
\begin{equation*}
\delta_{\omega}\left(a \alpha_{I}\right)=\sum_{i=1}^{2 n}\left(H_{z_{i}} a\right) \iota_{\partial / \partial z_{i}} \alpha_{I} \tag{10.59}
\end{equation*}
$$

Lemma 10.8. With $E_{1}$ identified with the formal sums of homogeneous forms on $T^{*} X \backslash 0$, the induced differential is

$$
\begin{equation*}
b_{(1)}=\frac{1}{i} \delta_{\omega} . \tag{10.60}
\end{equation*}
$$

Proof. We know that the bilinear differential operator $2 i P_{1}$ is the Poisson bracket of functions on $T^{*} X$. Thus (10.52) can be written

$$
\begin{align*}
& 2 i b_{1} \alpha=\sum_{i=0}^{k-1}(-1)^{i} a_{0} \otimes \cdots \otimes a_{i-1} \otimes\left\{a_{i+1}, a_{i}\right\} \otimes a_{i+2} \otimes \cdots \otimes a_{k}  \tag{10.61}\\
& +(-1)^{k}\left\{a_{0}, a_{k}\right\} \otimes a_{1} \otimes \cdots \otimes a_{k-1}, \alpha=a_{0} \otimes \cdots \otimes a_{k} .
\end{align*}
$$

The form to which this maps under the identification of $E_{2}$ is just
(10.62) $21 b_{1} \alpha=\sum_{i=0}^{k-1}(-1)^{i} a_{0} \wedge d a_{i-1} \wedge \cdots \wedge d\left\{a_{i+1}, a_{i}\right\} \wedge d a_{i+2} \wedge a_{k}$

$$
+(-1)^{k}\left\{a_{0}, a_{k}\right\} \wedge d a_{1} \wedge \cdots \wedge d a_{k-1}
$$

[^23]Consider the basis elements $\alpha_{I}$ for $k$-forms. These arise as the images of the corresponding functions in local coordinates on $X^{k+1}$

$$
\begin{aligned}
\tilde{\alpha}_{I}\left(z_{0}, z_{1}, \ldots, z_{k}\right)=\sum_{\sigma}(-1)^{\operatorname{sgn} \sigma} & \left(z_{1, \sigma I(1)}-z_{0, \sigma I(1)}\right) \\
& \left.\times z_{2, \sigma I(1)}-z_{1, \sigma I(1)}\right) \ldots\left(z_{1, \sigma I(m)}-z_{0, \sigma I(m-1)}\right)
\end{aligned}
$$

Since these functions are defined in local coordinates they are not globally defined on $\left(T^{*} X \backslash 0\right)^{k+1}$. Nevertheless they can be localized away from $z_{0}=\cdots=z_{m}$ and then, with a coefficient $\left(a_{j}\left(z_{0}\right)\right)_{j=-\infty}^{*}, a_{j} \in \mathcal{C}^{\infty}\left(T^{*} X \backslash 0\right)$ homogeneous of degree $j$ with support in the coordinate patch, unambiguously define elements of $E_{1}$ which we can simply denote as $a\left(z_{0}\right) \tilde{\alpha}_{I} \in E_{1}$. These elemments, superimposed over a coordinate cover, span $E_{1}$. Consider $b_{(1)} \tilde{\alpha}$ given by (ITV.62). In the sum, the terms with $P_{1}$ contracting between indices other than 0,1 or $m, 0$ must give zero because the Poisson bracket is constant in the 'middle' variable. Futhermore, by the antisymmetry of $\tilde{\alpha}$, the two remaining terms are equal so

$$
\begin{gathered}
i b_{(1)}\left(a \tilde{\alpha}_{I}\right)=\sum_{\sigma \in \mathcal{P}_{k}}\left(H_{z_{\sigma I(1)}} a\right)(-1)^{\operatorname{sgn}(\sigma)} d z_{\sigma I(2)} \wedge \cdots \wedge d z_{\sigma I(k)} \\
=\sum_{i}\left(H_{z_{i}} a\right) \iota_{\partial / \partial_{i}} \alpha_{I}
\end{gathered}
$$

Since this is just $\left(\frac{24.81}{10.59}\right)$ the lemma follows.
With this lemma we have identified the differential on the $E_{1}$ term in the spectral sequence with the exterior differential operator. To complete the identification (10.49) we need to compute the corresponding deRham groups.

Proposition 10.1. The cohomology of the complex

$$
\ldots \xrightarrow{d} \sum_{j=-\infty}^{*} \mathcal{C}_{\text {hom }(j)}^{\infty}\left(T^{*} X \backslash 0 ; \Lambda^{k}\right) \xrightarrow{d} \sum_{j=-\infty}^{*} \mathcal{C}_{\text {hom }(j)}^{\infty}\left(T^{*} X \backslash 0 ; \Lambda^{k+1}\right) \xrightarrow{d} \ldots
$$

in dimension $k$ is naturally isomorphic to $H^{k}\left(S^{*} X\right) \oplus H^{k-1}\left(S^{*} X\right)$.
Proof. Choose a metric on $X$ and let $R=|\xi|$ denote the corresponding length function on $T^{*} X \backslash 0$. Thus, identifying the quotient $S^{*} X=\left(T^{*} X \backslash 0\right) / \mathbb{R}^{+}$with $\{R=1\}$ gives an isomorphism $T^{*} X \backslash 0 \equiv S^{*} X \times(0, \infty)$. Under this map the smooth forms on $T^{*} X \backslash 0$ which are homogeneous of degree $j$ are identified as sums

$$
\begin{array}{r}
\mathcal{C}_{\mathrm{hom}(j)}^{\infty}\left(T^{*} X \backslash 0, \Lambda^{k}\right) \ni \alpha_{j} \\
=R^{j}\left(\alpha_{j}^{\prime}+\alpha_{j}^{\prime \prime} \wedge \frac{d R}{R}\right), \alpha_{j}^{\prime} \in \mathcal{C}^{\infty}\left(S^{*} X ; \Lambda^{k}\right), \alpha_{j}^{\prime \prime} \in \mathcal{C}^{\infty}\left(S^{*} X ; \Lambda^{k-1}\right) \tag{10.63}
\end{array}
$$

The action of the exterior derivative is then easily computed

$$
\begin{aligned}
d \alpha_{j} & =\beta_{j}, \quad \beta_{j}=R^{j}\left(\beta_{j}^{\prime}+\beta_{j}-^{\prime \prime} \wedge \frac{d R}{R}\right) \\
\beta_{j}^{\prime} & =d \alpha_{j}^{\prime}, \quad \beta_{j}^{\prime \prime}=d \alpha_{j}^{\prime \prime}+j(-1)^{k-1} \alpha_{j}^{\prime}
\end{aligned}
$$

Thus a $k$-form $\left(\alpha_{j}\right)_{j=-\infty}^{*}$ is closed precisely if it satisfies

$$
\begin{equation*}
j \alpha_{j}^{\prime}=(-1)^{k} d \alpha_{j}^{\prime \prime}, d \alpha_{j}^{\prime}=0 \forall j \tag{10.64}
\end{equation*}
$$

It is exact if there exists a $(k-1)$-form $\left(\gamma_{j}\right)_{j=-\infty}^{*}$ such that
24.85

$$
\begin{equation*}
\alpha_{j}^{\prime}=d \gamma_{j}^{\prime}, \quad \alpha_{j}^{\prime \prime}=d \gamma_{j}^{\prime \prime}+j(-1)^{k} \gamma_{j}^{\prime} \tag{10.65}
\end{equation*}
$$

Since the differential preserves homogeneity it is only necessary to analyze these equations for each inteqral $j$. For $j \neq 0$, the second equation in (10.64) follows from the first and $\left(\frac{24.85}{10.65}\right)$ then holds with $\gamma_{j}^{\prime}=\frac{1}{j}(-1)^{k} \alpha_{j}^{\prime \prime}$ and $\gamma_{j}^{\prime \prime}=0$. Thus the coh. 24.840 logy lies only in the subcomplex of homogeneous forms of degree 0 . Then (10.64) and (10.65) become

$$
d \alpha_{0}^{\prime}=0, d \alpha_{0}^{\prime \prime}=0 \text { and } \alpha_{0}^{\prime}=d \gamma_{0}^{\prime}, \alpha_{0}^{\prime \prime}=d \gamma_{0}^{\prime \prime}
$$

respectively. This gives exactly the direct sum of $H^{k}\left(S^{*} X\right)$ and $H^{k-1}\left(S^{*} X\right)$ as the cohomology in degree $k$. The resulting isomorphism is independent of the choice of the radial function $R$, since another choice replaces $R$ by $R a$, where $a$ is a smooth positive function on $S^{*} X$. In the decomposition (10.63), for $j=0, \alpha_{0}^{\prime \prime}$ is unchanged whereas $\alpha_{0}^{\prime}$ is replaced by $\alpha_{0}^{\prime}+\alpha_{0}^{\prime \prime} \wedge d \log a$. Since the extra term is exact whenever $\alpha_{0}^{\prime \prime}$ is closed it has no effect on the identification of the cohomology.

Combining Proposition $\frac{24.82}{10.1}$ and Lemma $\frac{24.67}{10.8}$ completes the proof of ( $\frac{24.63}{10.49)}$. We make the identification a little more precise by locating the terms in $E_{2}$.

Proposition 10.2. Under the identification of $E_{1}$ with the sums of homogeneous forms on $T^{*} X \backslash 0, E_{2}$, identified as the cohomology of $\delta_{\omega}$, has a basis of homogeneous forms with the homogeneity degree $j$ and the form degree $k$ confined to

$$
\begin{equation*}
k-j=\operatorname{dim} X,-\operatorname{dim} X \leq j \leq \operatorname{dim} X, \operatorname{dim} X \geq 2 \tag{10.66}
\end{equation*}
$$

Proof. Provided $\operatorname{dim} X \geq 2$, the cohomology of $S^{*} X$ is isomorphic to two copies of the cohomology of $X$, one in the same degree and one shifted by $\operatorname{dim} X-$ 1. ${ }^{16}$ The classes in the first copy can be taken to be the lifts of deRham classes from $X$, while the second is spanned by the wedge of these same classes with the Todd class of $S^{*} X$. This latter, $n-1$, class restricts to each fibre to be non-vanishing. Thus in local representations the first forms involve only the base variable and in the second each terms has the maximum number, $n-1$, of fibre forms. The cohomology of the complex in Proposition IIU.I therefore consists of four copies of $H^{*}(X)$ consisting of these forms and the same forms wedged with $d R / R$.

With this decomoposition of the cohomology consider the effect on it of the map $W_{\omega}$. In each case the image forms are again homogeneous. A deRham class on $X$ in degree $l$ therefore has four images in $E_{2}$. One is a form of degree $k_{1}=2 n-l$ which is homogeneous of degree $j_{1}=n-l$. The second is a form of degree $k_{2}=2 n-l-1$ which is homogeneous of degree $j_{2}=n-l-1$. The third image is of form degree $k_{3}=n-l+1$ and homogeneous of degree $j_{3}=-l+1$ and the final image is of form degree $k_{4}=n-l$ and is homogeneous of degree $j_{4}=-l$. This gives the relations (10.66).

### 10.8. Degeneration and convergence

Now that the $E_{2}$ term in the spectral sequence has been explicitly computed, consider the induced differential, $b_{(2)}$ on it. Any homogeneous form representing a class in $E_{2}$ can be represented by a Hochshild chain $\alpha$ of the same homogeneity. Thus an element of $E_{2}$ in degree $k$ corresponds to a function on $\left.\mathcal{C}^{\infty}\left(\left(T^{*} X\right) \backslash\right)^{k+1}\right)$ which is separately homogeneous in each variable and of total homogeneity $k-n$. Furthermore it has an extension $\beta(t)$ as a function of the parameter $h$, of the same

[^24]homogeneity, such that $b_{t} \beta(t)=t^{2} \gamma(t)$. Then $b_{(2)} \alpha=[\gamma(0)]$, the class of $\gamma(0)$ in $E_{2}$. Noting that the differential operator, $P_{j}$, which is the $j$ th term in the Taylor series of the product $\star_{h}$ reduces homogeneity by $j$ and that $b_{h}$ depends multilinearly on $\star_{h}$ it follows tha $b_{(24)}$ must decrease homogeneity by $r$. Thus if the class $[\gamma(0)]$ must vanish in $E_{2}$ by (10.66). We have therefore shown that $b_{(2)} \equiv 0$, so $E_{3}=E_{2}$. The same argument applies to the higher differentials, defining the $E_{r} \equiv E_{2}$ for $r \geq 2$, proving the 'degeneration' of the spectral sequence. ( 24.65 .50 ).

The 'convergence' of the spectral sequence, (110.51), follows from the same analysis of homogneities. Thus, we shall define a map from $E_{2}$ to the Hochschild homology and show that it is an isomorphism.

### 10.9. Explicit cohomology maps

10.10. Hochschild holomology of $\Psi^{-\infty}(X)$
10.11. Hochschild holomology of $\Psi^{\mathbb{Z}}(X)$

### 10.12. Morita equivalence

CHAPTER 11

The index formula

HHcoId (11.1)

## APPENDIX A

## Bounded operators on Hilbert space

Some of the main properties of bounded operators on a complex Hilbert space, $H$, are recalled here; they are assumed at various points in the text.
(1) Boundedness equals continuity, $\mathcal{B}(H)$.
(2) $\|A B\| \leq\|A\|\|B\|$
(3) $(A-\lambda)^{-1} \in \mathcal{B}(H)$ if $|\lambda| \geq\|A\|$.
(4) $\left\|A^{*} A\right\|=\left\|A A^{*}\right\|=\|A\|^{2}$.
(5) Compact operators, defined by requiring the closure of the image of the unit ball to be compact, form the norm closure of the operators of finite rank.
(6) Fredholm operators have parametrices up to compact errors.
(7) Fredholm operators have generalized inverses.
(8) Fredholm operators for an open subalgebra.
(9) Hilbert-Schmidt operators?
(10) Operators of trace class?
(11) General Schatten class?

## Bibliography




[^0]:    ${ }^{1}$ I suggest Friedlander's little book $\frac{\text { MR86h:46002 }}{[3] \text { (there is }}$ lso newer edition) as a good introduction to distributions. Volume 1 of Hörmander's treatise $[7]$ has all that you would need and a good deal more; it is a good general reference.

[^1]:    ${ }^{2}$ This is the standard action, but it is potentially confusing since it means that for the basis elements $e_{j} \in \mathbb{R}^{n}, L e_{j}=\sum_{k=1}^{n} L_{k j} e_{k}$.
    ${ }^{3}$ So $D_{j}$ transforms as a basis of $\mathbb{R}^{n}$ as it should, despite the factors of $i$.

[^2]:    ${ }^{4}$ See Problem $\frac{1.2 \cdot 2000.278}{1.9 .}$

[^3]:    ${ }^{1}$ Note it is required that $\epsilon$ be chosen to be independent of $z$ here, so this is a notion of uniform ellipticity.

[^4]:    ${ }^{2}$ See Problem $\frac{1.2 .2000 .276}{2.5 .}$

[^5]:    ${ }^{3}$ Émile Borel

[^6]:    ${ }^{4}$ Schur

[^7]:    ${ }^{5}$ For a somewhat more general class of polyhomogeneous symbols, see problem $\frac{\text { pr }}{2.8}$

[^8]:    ${ }^{6}$ Polyhomogeneous symbols may seem to be quite sophisticated ${ }^{\text {obidects }}$ but they are really smooth functions on manifolds with boundary; see Problems 12.872 .7.
    ${ }^{7}$ See Problem $\frac{\text { prob }}{1.9}$

[^9]:    ${ }^{8}$ See Problem $\frac{\text { prob:CC }}{2.10 .}$

[^10]:    1.2 .2000 .328

[^11]:    ${ }^{1}$ We need to show that $\|\tilde{B} f\|$ is bounded when $f \in R_{1}$ and $\|f\|=1$. This is just the boundedness of $u \in R_{0}$ when $f=(\operatorname{Id}+A) u$ is bounded in $R_{1}$.

[^12]:    ${ }^{2}$ Namely the trace of a finite rank projection, such as either $\Pi_{0}$ or $\Pi_{1}$, is its rank, hence the dimension of the space onto which it projects. From the identity satisfied by the generalized inverse we see that

    $$
    \operatorname{Ind}(\operatorname{Id}+A)=\operatorname{Tr}\left(\Pi_{0}\right)-\operatorname{Tr}\left(\Pi_{1}\right)=\operatorname{Tr}((\operatorname{Id}+B)(\operatorname{Id}+A)-(\operatorname{Id}+A)(\operatorname{Id}+B))=\operatorname{Tr}([B, A])=0
    $$

    from the basic property of the trace.

[^13]:    ${ }^{3}$ This is an essentially microlocal proof.

[^14]:    ${ }^{1}$ The method used here to compute the homology of a polynomial algebra is due to Sergiu Moroianu; thanks Sergiu.

[^15]:    ${ }^{2}$ Hence Taylor's theorem cannot be applied.

[^16]:    ${ }^{3}$ Really on the dual but that does not matter at this stage.

[^17]:    ${ }^{4}$ One way to justify this is to use results on smoothing operators. For finite dimensional linear spaces $V$ and $W$ the tensor product can be realized as

    $$
    V \otimes W=\operatorname{hom}\left(W^{\prime}, V\right)
    $$

    the space of linear maps from the dual of $W$ to $V$. Identifying the topological dual of $\mathcal{C}^{\infty}(X)$ with $\mathcal{C}_{c}^{-\infty}(X ; \Omega)$, the space of distributions of compact support, with the weak topology, we can identify the projective tensor product $\mathcal{C}^{\infty}(X) \hat{\otimes} \mathcal{C}^{\infty}(X)$ as the space of continuous linear maps from $\mathcal{C}_{c}^{-\infty}(X ; \Omega)$ to $\mathcal{C}^{\infty}(X)$. These are precisely the smoothing operators, corresponding to kernels in $\mathcal{C}^{\infty}(X \times X)$.

[^18]:    ${ }^{5}$ This homology is properly referred to as the continuous Hochschild homology of the topological algebra $\mathcal{C}^{\infty}(X)$.
    ${ }^{6}$ As pointed out to me by Maciej Zworski, this is a form of Hadamard's lemma.
    ${ }^{7}$ Meaning here the continuous H-unitality, that is the acyclicity of $b^{\prime}$ on the chain spaces $\mathcal{C}^{\infty}\left(X^{k+1}\right)$.

[^19]:    ${ }^{8}$ Notice that $v\left(z_{0}, \ldots, z_{j}, z_{j+2}, \ldots, z_{k+1}\right)$ vanishes on $z_{i+1}=z_{i}$ for $i<j$ and $i>j+1$ and also on $z_{0}+z_{1}+\cdots+z_{k+1}=0$ (since it is independent of $z_{j+1}$ and $b v=0$.

[^20]:    ${ }^{9}$ In particular $44^{\text {th }} 58$ Hochsfhildd homology vanishes for $k>2 \operatorname{dim} X$. For a precise form of the identification in (110.43) see (ㅍ1.1).

[^21]:    ${ }^{10}$ If $\alpha \in G_{p, k}$ and $\beta(h)$ is the 1-parameter family of chains whose existence is required for the definition then $\beta^{\prime}(h)=h \beta(h)$ satisfies the same condition with $p$ increased to show that $\alpha \in G_{p+1, k}$.
    ${ }^{11}$ Of Leray I suppose, but I am not really sure.

[^22]:    ${ }^{12}$ Indeed, $\alpha$ is then the value at $h=0$ of $\beta(t)=t^{-p+1} b_{h} \phi(t)$ which is by hypothesis smooth; clearly $b_{h} \beta \equiv 0$.
    ${ }^{13}$ If you feel it necessary to do so, resort to an argument by continuity towards the end of this discussion.

[^23]:    ${ }^{14} \mathrm{Up}$ to various sign conventions of course!
    ${ }^{15}$ Check this case by case, as the range of $I$ meets the pair $\{2 p-1,2 p\}$ in $\{2 p-1,2 p\}$, $\{2 p-1\},\{2 p\}$ or $\emptyset$. Both sides of the first equation are zero in the second and fourth case as are both sides of the second equation in the third and fourth cases. In the remaining four individual cases it is a matter of checking signs.

[^24]:    ${ }^{16}$ That is, just as though $S^{*} X=\mathbb{S}^{n-1} \times X$, where $n=\operatorname{dim} X$.

