

Evaluation and Comparison of Continuum Models for Dense Granular Flow

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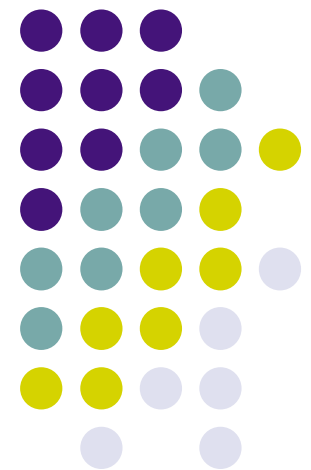
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[Camilo Guaqueta, Ken Weaver]
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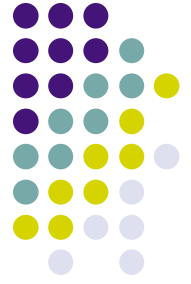
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Dense Flow Models:

1. Mohr-Coulomb Plasticity

[Sokolovskii (1965), Jenike & Johanson (1962)]

2. Hourglass Theory

[Savage (1967), Sullivan (1972), Davidson & Nedderman (1973)]

3. Kinematic Model

[Litwiniszyn (1963), Nedderman & Tuzun (1979)]

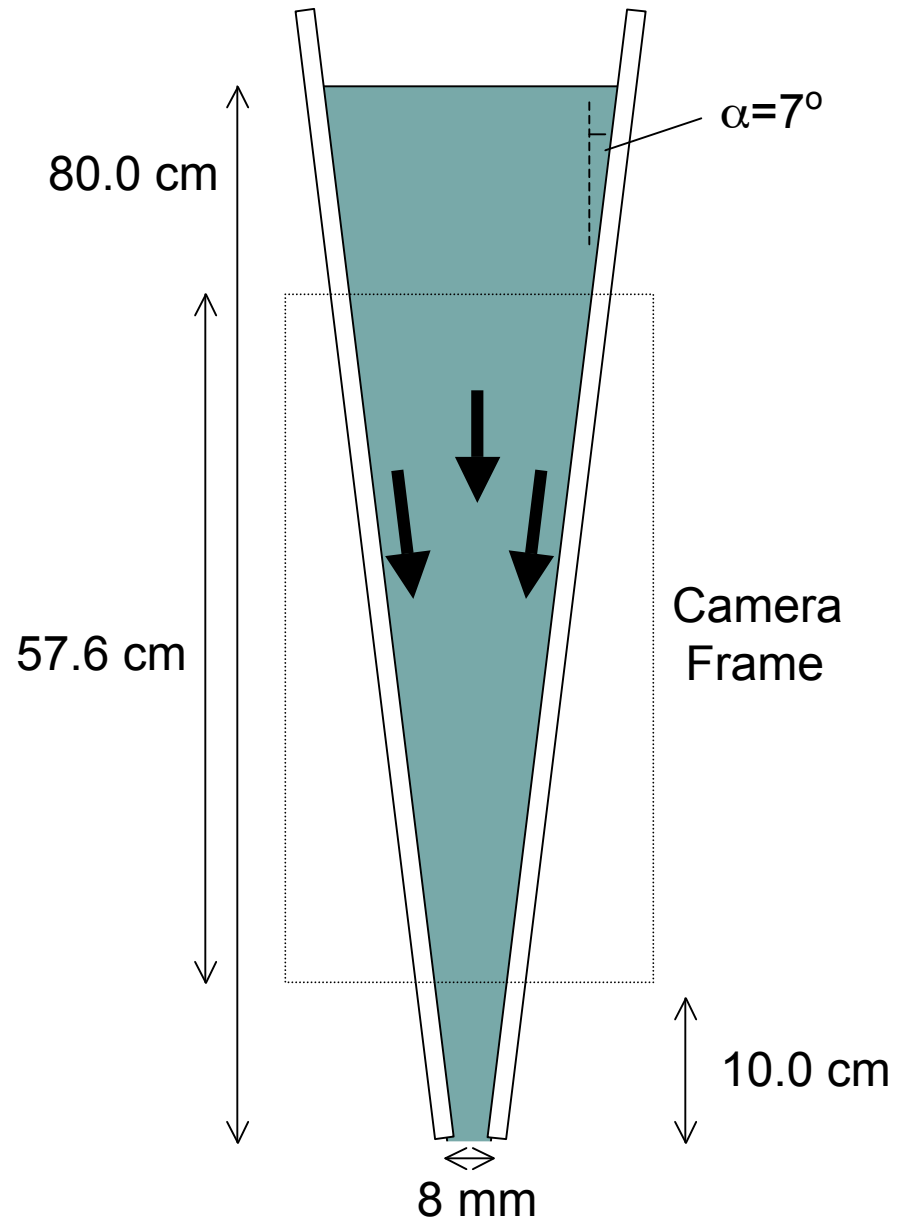
Setup

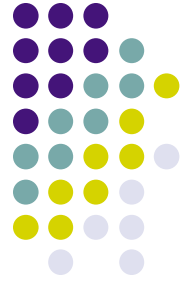
Grains: 3 mm glass spheres

Side Walls: smooth and fully rough

Depth: 2.5 cm

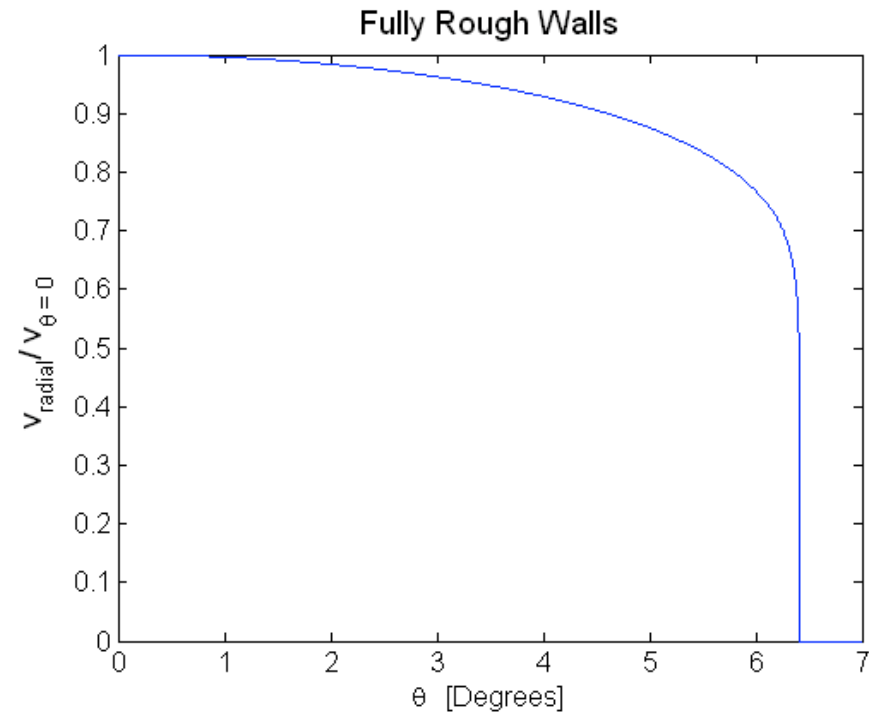
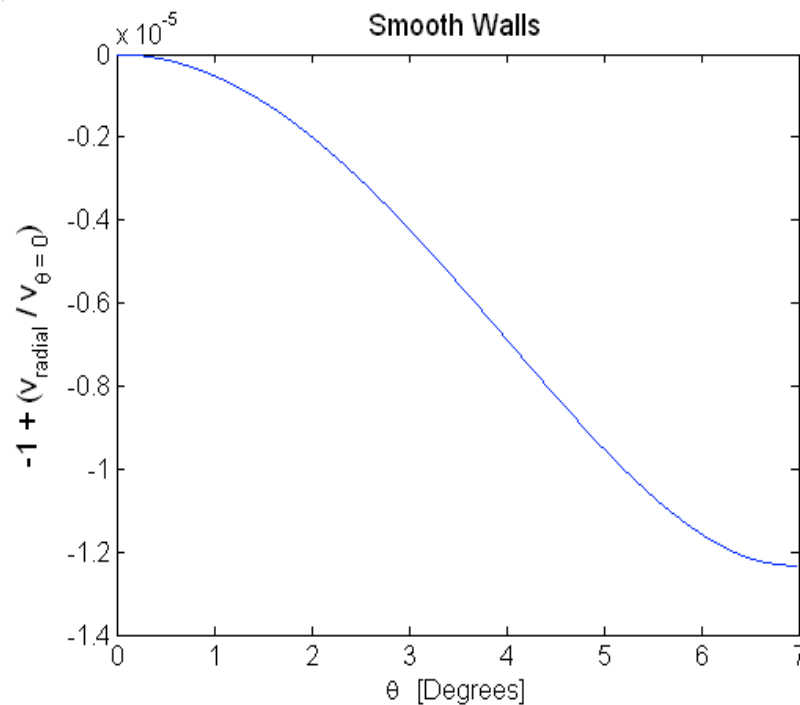
Camera: 250 fps



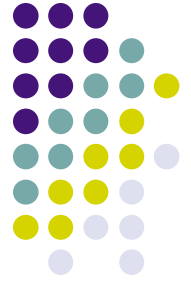


Mohr-Coulomb Plasticity

- Yield occurs only when $\tau = \mu_{\text{internal}} \sigma = (\tan \phi) \sigma$. For us, $\phi = 24.5^\circ$.
- Material assumed to be at incipient yield everywhere and quasi-static.
- Upholds Levy flow rule.
- Jenike found a similarity solution operative in a wedge geometry which gives radial flow varying in θ :



Hourglass Theory



- Least complicated flow model to apply.
- Based on M-C Plasticity with simplifications to enable convective solution.
- Only for wedge or conical hopper geometry with small apex angle.
- Requires hopper walls to be frictionless. Asserts that flow is radial and constant in θ .

Result for narrow wedge:

$$\vec{v} = -\frac{A}{r}\hat{r}$$

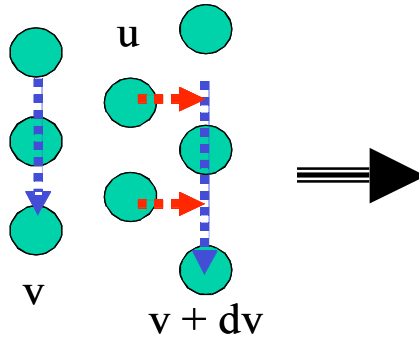
$$\text{for } A = \sqrt{\frac{(K+1)(r_0^{K+1}r_1^3 - r_0^3r_1^{K+1})g}{(K-2)(r_0^{K+1} - r_1^{K+1})}}, \quad K = \frac{1 + \sin \phi}{1 - \sin \phi}.$$

Thus we have, $A = 370.0 \text{ cm}^2\text{s}^{-1}$.

Kinematic Model

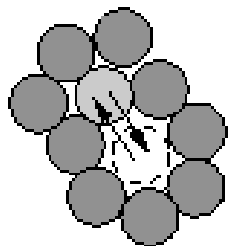


General principle:

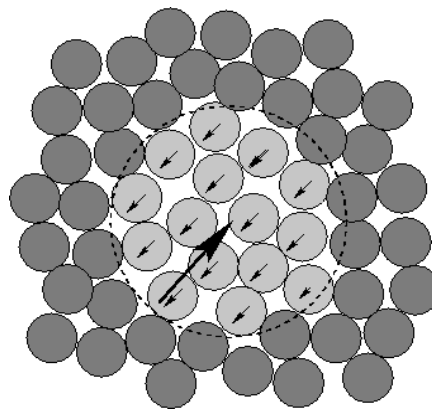


$$u = b \frac{\partial v}{\partial x} \quad , \quad \frac{\partial v}{\partial z} = b \frac{\partial^2 v}{\partial x^2}$$

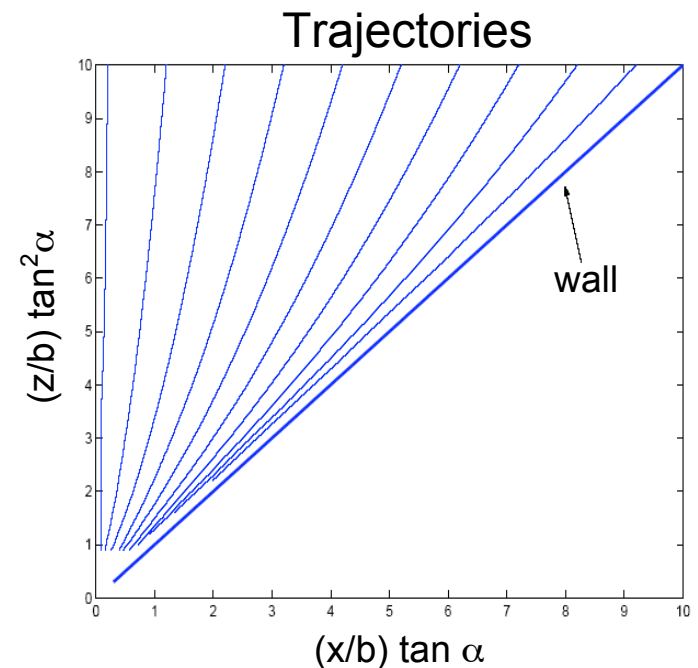
- Only for use in silos and hoppers.
- No dependence on internal or wall friction.
- Flow governed by boundary conditions and empirically determined diffusion length b .



Void model



Spot model



Experiment: Smooth Walls

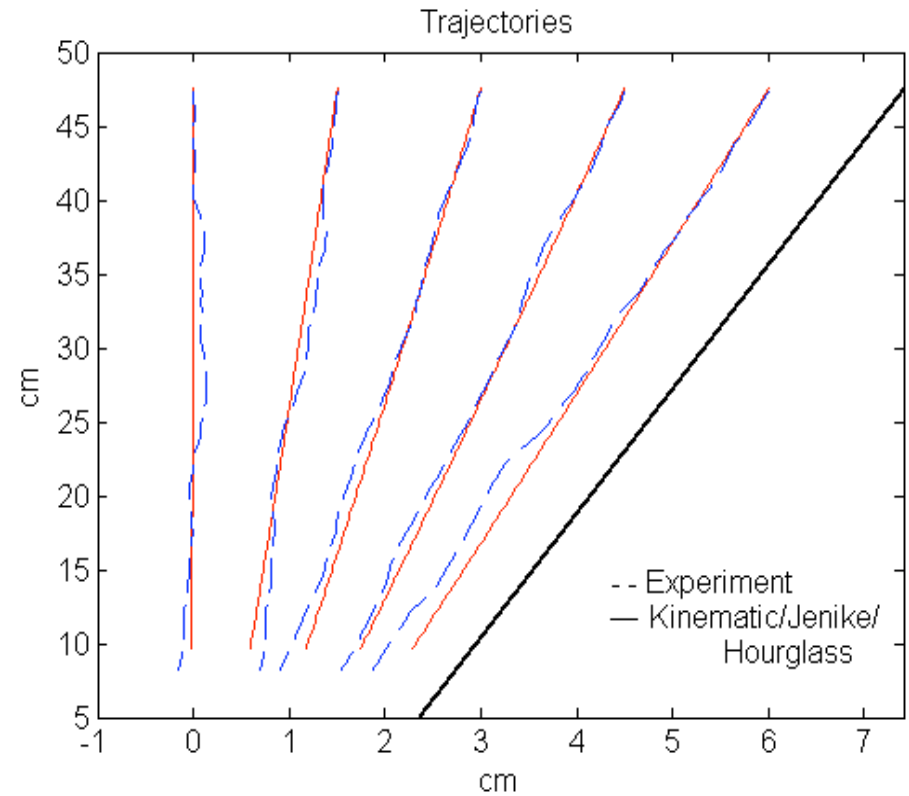
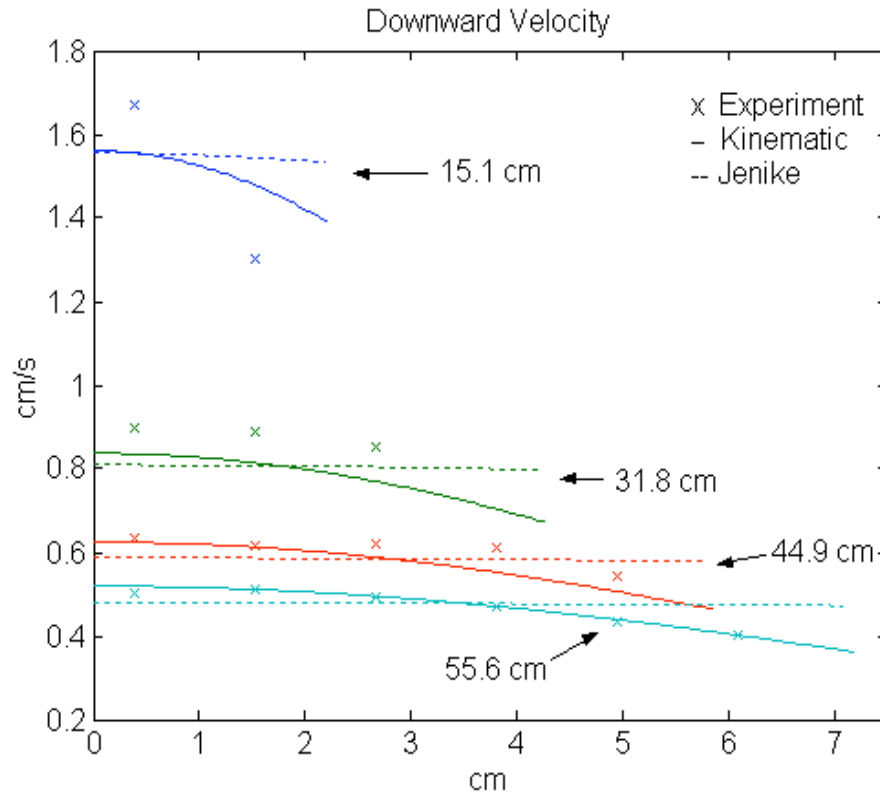


Hourglass Constant:

$A = 370.0 \text{ cm}^2/\text{s}$ theoretically
 $A = 25.0 \text{ cm}^2/\text{s}$ experimentally

Kinematic Diffusion Length:

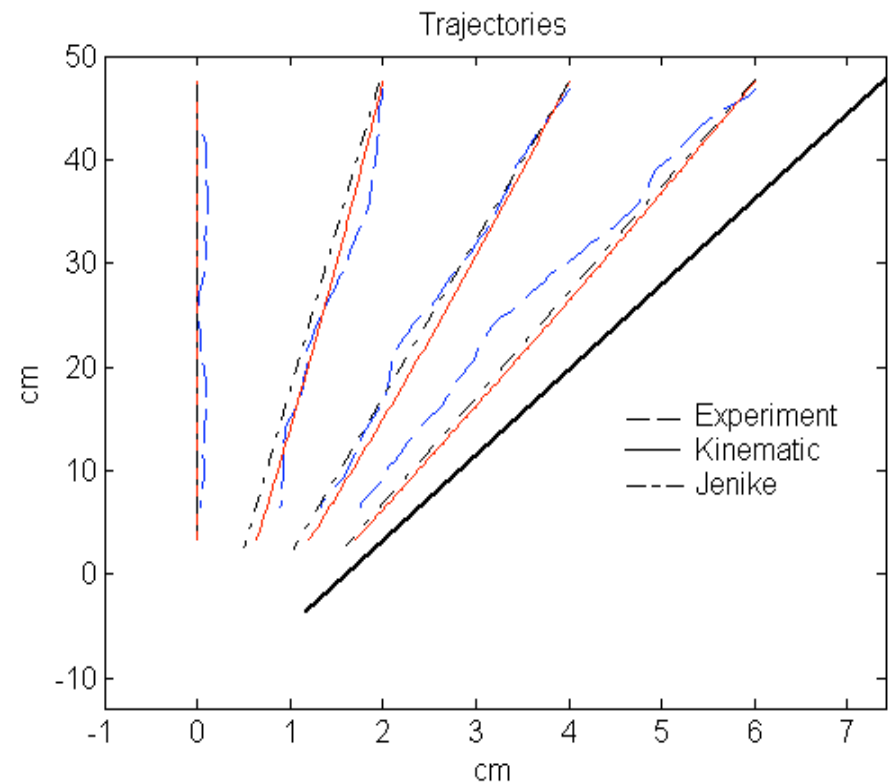
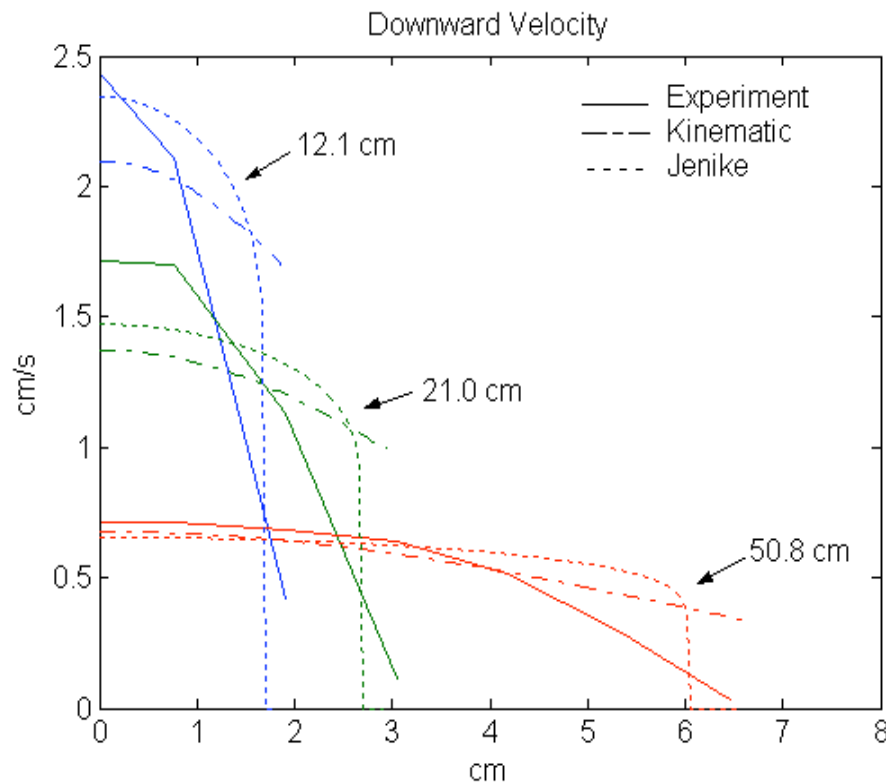
$b = 3.86d$



Experiment: Fully Rough Walls



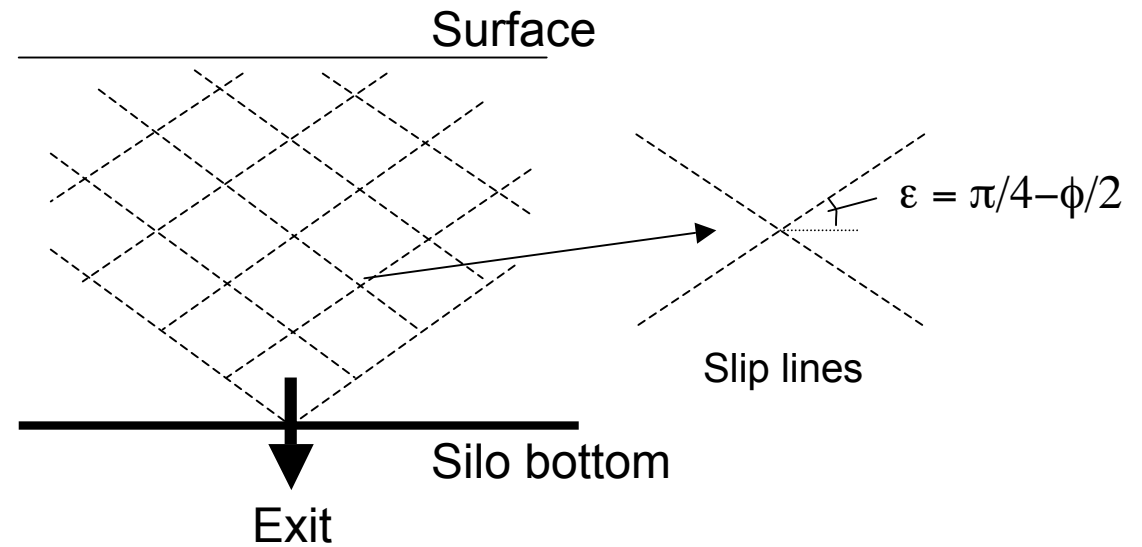
Kinematic Diffusion Length: $b = 1.77d$



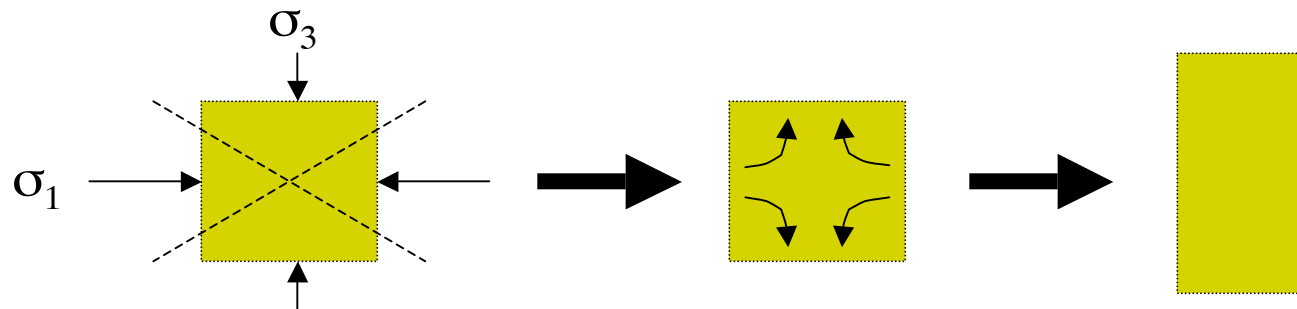


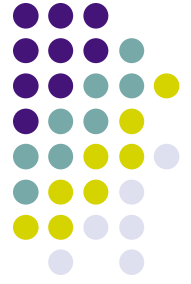
Possible Connections

In a wide, approximately 2D silo, solution to Mohr-Coulomb slip lines turn out to be straight lines along $|\theta| = \epsilon$ extending down from the surface.

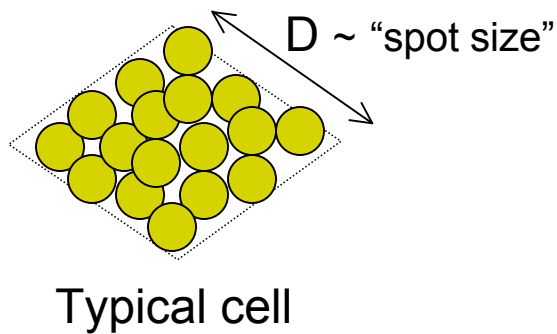
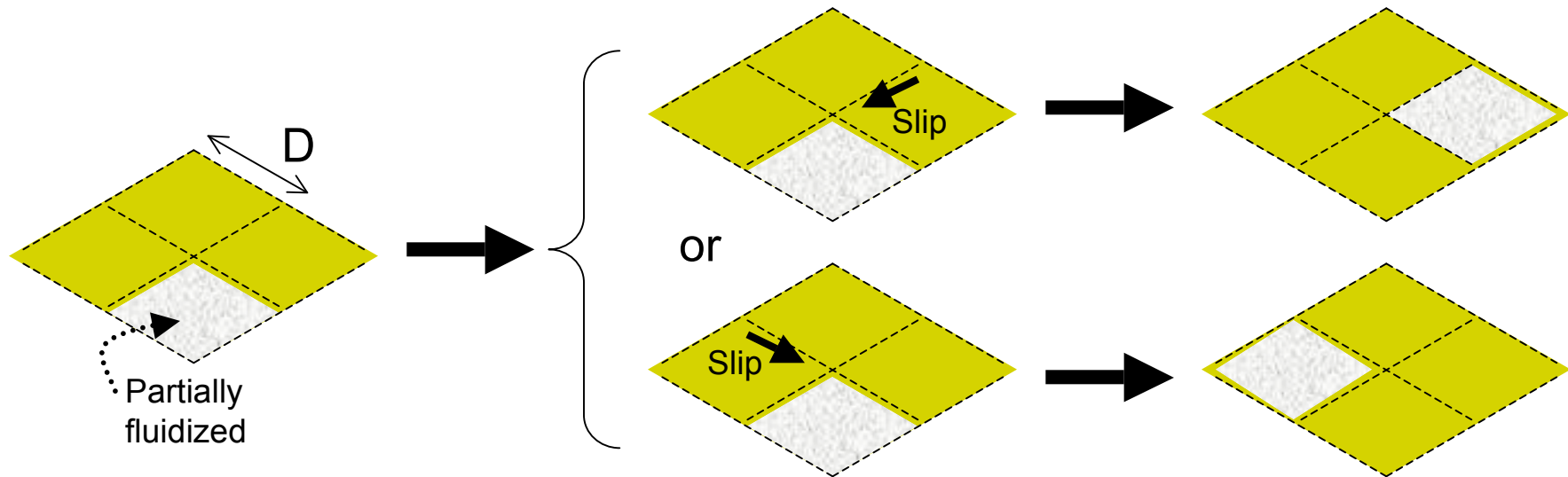


**Co-axiality
(Levy Flow Rule):**





Stochastic Flow Rule



Rule produces analytic formula for Kinematic b :

$$b = D \cos^2 \epsilon / (2 \sin \epsilon)$$

For glass beads: $20^\circ < \phi < 25^\circ$, $3d < D < 5d$, predict $1.75 d < b < 3.31 d$. (Experiments: $b = 1.3d - 3.4d$.)

Conclusions



- 1) While each model accounts well for certain effects, none alone can accurately predict a dense flow with generality.**
- 2) The Kinematic Model, successful thus far as an empirical model, may be mechanically derivable from the Mohr-Coulomb stress equations but with a stochastic flow rule. May enable us to generalize the use of Kinematic Modeling beyond drainage.**