

Symplectic topology and q -intersection numbers
joint work with R. Bekturkavnikov and J. Solomon

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Horse lecture # 2
April 2010

intersection number
 $L_0 \cdot L_1$, an integer*

* $(-1)^{n(n-1)/2}$ times the
 standard convention

$q=1$
 improved or q-intersection
 number $L_0 \cdot_q L_1$, a function
 of a formal variable q

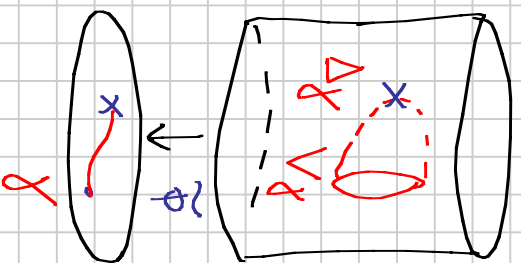
Here, M^{2n} is an oriented smooth manifold, and L_0, L_1 are closed oriented submanifolds. In many situations, one would like to enhance the ordinary intersection number to a q -version. Obviously, this doesn't always work!

There is a classical topological approach based on infinite cyclic coverings, where $\mathbb{Z}[q, q^{-1}]$ is the group ring of the covering group.

Sample application (Givental) Take a polynomial $p: \mathbb{C}^{n+1} \rightarrow \mathbb{C}$ with an isolated critical point at the origin. To simplify the exposition, assume p is weighted homogeneous. The Milnor fibre is

$$M = p^{-1}(t)$$

for $t \neq 0$. $H_n(M) \cong \mathbb{Z}^n$ is generated by vanishing cycles. To see those, replace p with a perturbation \tilde{p} .



γ path from t to a critical value of \tilde{p}

\forall sphere in $\tilde{p}^{-1}(t) = M$

Δ_X ball in (\mathbb{C}^{n+1}, M) , the Lefschetz thimble

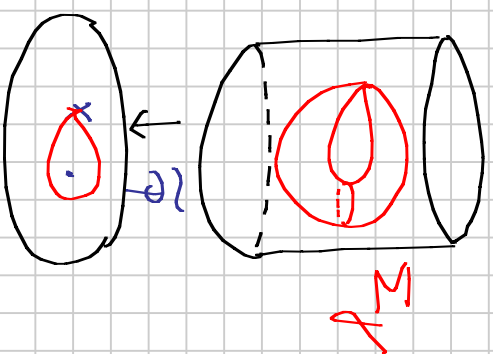
Take $C = \mathbb{R}^{n+1} \setminus M$, and $\tilde{C} \rightarrow C$ its obvious \mathbb{Z} -covering.
 Given X , we have two ways of getting an associated cycle in \tilde{C} :

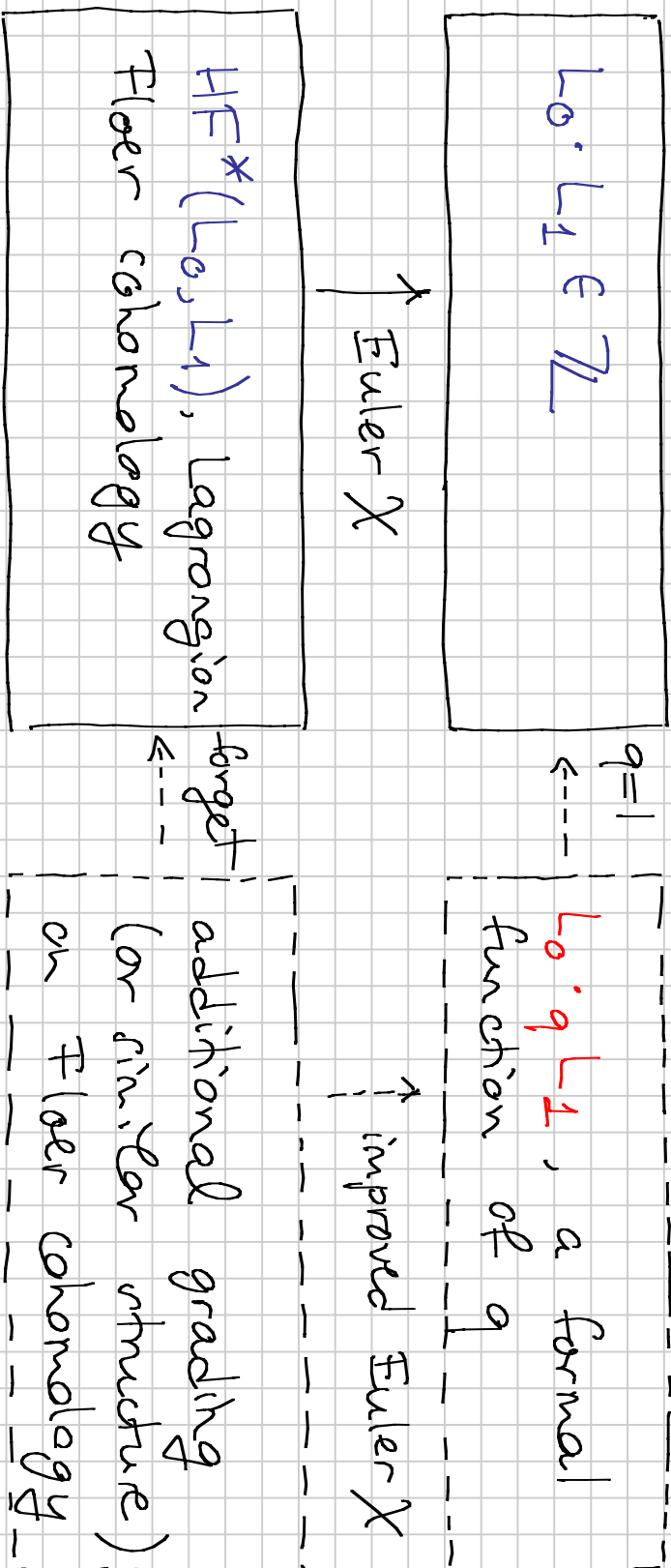
- Lift $\Delta_X \setminus V_X$ to $\tilde{\Delta}_X$ (yields a properly embedded open ball)
- double the path, which yields an immersed sphere $\Sigma_X \subset C$, and lift that to $\tilde{\Sigma}_X$:

If $\sigma: \tilde{C} \rightarrow \tilde{C}$ is the deck transformation, then Givental defines

$$\forall x_0 \cdot q \forall t_1 = \sum_{k \in \mathbb{Z}} q^k \sigma^k(\tilde{\Delta}_{x_0}) \cdot \tilde{\Sigma}_{t_1}$$

He proves that this is well-defined and refines the standard intersection form.





This is what we really want (for M a symplectic manifold, and L_0, L_1 Lagrangian submanifolds). One could follow Gromov's approach for Hilbert fibres, but his construction leaves M and its asymmetry because problematic in the categorified world.

Mirror symmetry inspiration: Take X a smooth algebraic variety, and consider compactly supported coherent sheaves (or more generally objects of $\mathcal{D}_c^b(X)$).
By Hodgebruch-Riemann-Roch,

$$\chi(\mathrm{Ext}_X^*(\mathcal{E}_0, \mathcal{E}_1)) = \langle \mathcal{E}_0 | \mathcal{E}_1 \rangle = \int_X \mathrm{ch}(\mathcal{E}_0^*) \mathrm{ch}(\mathcal{E}_1) \mathrm{Td}(X)$$

If X is Calabi-Yau, the right hand side is the graded symmetric Mukai pairing. This can be seen directly or follows from Serre duality

$$\mathrm{Ext}^*(\mathcal{E}_1, \mathcal{E}_0) \cong \mathrm{Ext}^{n-*}(\mathcal{E}_0, \mathcal{E}_1 \otimes \mathcal{K}_X)^\vee$$

coherent with compact support

Now suppose that X carries a holomorphic \mathbb{C}^* -action, and consider equivariant sheaves. Then $\mathrm{Ext}^*(\mathcal{E}_0, \mathcal{E}_1)$ is a representation of \mathbb{C}^* , hence decomposes into weight spaces $\mathrm{Ext}^*(\mathcal{E}_0, \mathcal{E}_1)^k$. Define

$$\langle \mathcal{E}_0 | \mathcal{E}_1 \rangle_q = \sum_k q^k \chi(\mathrm{Ext}^*(\mathcal{E}_0, \mathcal{E}_1)^k).$$

Modifications of the "equivariant" picture

Instead of a \mathbb{C}^* -action, let's take a holomorphic vector field Z . This induces elements $\Phi_Z^0 \in \text{Ext}_X^1(\mathcal{E}, \mathcal{E})$ for each \mathcal{E} . Instead of asking for equivariance, we only assume that these are zero. Then, there is an endomorphism

$$\tilde{\Phi}_{\mathcal{E}_0, \mathcal{E}_1}^1 : \text{Ext}_X^*(\mathcal{E}_0, \mathcal{E}_1) \rightarrow \text{Ext}_X^*(\mathcal{E}_0, \mathcal{E}_1)$$

This may not be diagonalizable, but its eigenvalues still decompose $\text{Ext}_X^*(\mathcal{E}_0, \mathcal{E}_1) = \bigoplus_{\lambda \in \mathbb{C}} \text{Ext}_X^*(\mathcal{E}_0, \mathcal{E}_1)^\lambda$. We have q -Mukai pairings

$$\begin{aligned} \langle \mathcal{E}_0 | \mathcal{E}_1 \rangle_q &= \sum_{\lambda} q^\lambda \cdot \mathcal{N}(\text{Ext}_X^*(\mathcal{E}_0, \mathcal{E}_1)^\lambda) \\ &= \text{Str}(\exp(\log(q) \tilde{\Phi}_{\mathcal{E}_0, \mathcal{E}_1}^1)) \end{aligned}$$

Actually, $\tilde{\Phi}_{\mathcal{E}_0, \mathcal{E}_1}^1$ is not strictly unique.

The difference between two choices is the sum of composition maps with elements of $\text{Ext}_X^0(\mathcal{E}_k, \mathcal{E}_k)$ (on the left and right, respectively).

Example If $\text{Ext}_X^0(\mathcal{E}_k, \mathcal{E}_k) = \mathbb{C} \cdot \text{Id}$, then $\tilde{\Phi}_{\mathcal{E}_0, \mathcal{E}_1}^1$ is unique up to constant multiples of the identity. If in addition $\mathcal{E}_0 = \mathcal{E}_1 = \mathcal{E}$, then $\tilde{\Phi}_{\mathcal{E}, \mathcal{E}}^1$ is independent of any choices.

From now on, fix a holomorphic volume form η . The situation where $L_{Z\eta} = 0$ is not so interesting from a mirror symmetry viewpoint, so let's assume that $L_{Z\eta} = \eta$. Then, Serre duality shows that

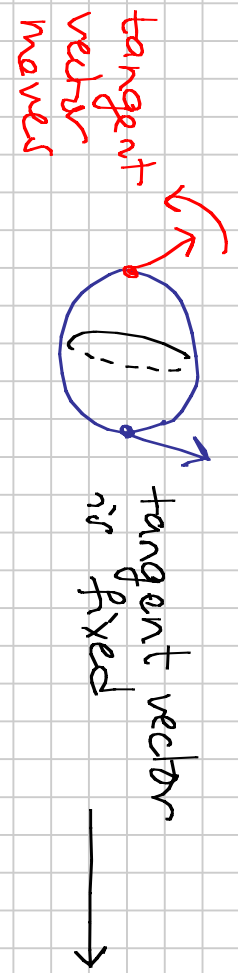
$$\langle \mathcal{E}_1 | \mathcal{E}_0 \rangle_g = \bar{g}'(-1)^n \langle \mathcal{E}_0 | \mathcal{E}_1 \rangle_{g-1}$$

Example If \mathcal{E} is spherical, then $\langle \mathcal{E} | \mathcal{E} \rangle_g = 1 + g(-1)^n$.

Reformulation as extended TCFIT associated to (X, η)
 (Costello, Hopkins-Lurie, Caldararu-Willett)



where by HKR $HH^*(X) \cong H^*(X, \Lambda^* T_X)$, so $Z \in HH^{-1}(X, X)$.
 One also has operations for families of surfaces:



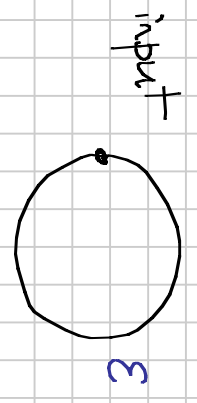
$L_Z \eta = \eta$ means $\text{diz} \eta = \eta$,
 hence $\Delta Z = 1$.

BV operator Δ of degree -1
 $HH^*(X) \cong H^*(X, \Omega_X^*)$
 $\Delta \downarrow$
 $HH^*(X) \cong H^*(X, \Omega_X^*)$
 $\downarrow d$
 algebraic de Rham

One further considers surfaces with boundaries marked with objects of $\mathcal{D}_\mathbb{C}^b(X)$.

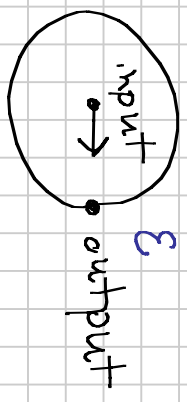


$$1 \in \text{Ext}_X^0(\mathcal{E}, \mathcal{E})$$

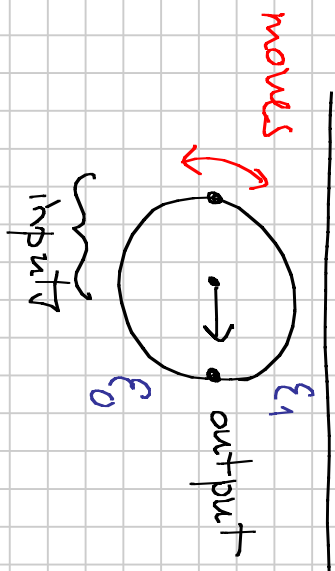


$$\int : \text{Ext}_X^n(\mathcal{E}, \mathcal{E}) = \text{Ext}_X^n(\mathcal{E}, \mathcal{E} \otimes R_X) \rightarrow \mathbb{C}$$

(or since \mathcal{E} has compact support)



$$\text{HH}^*(X, X) \rightarrow \text{Ext}_X^*(\mathcal{E}, \mathcal{E})$$

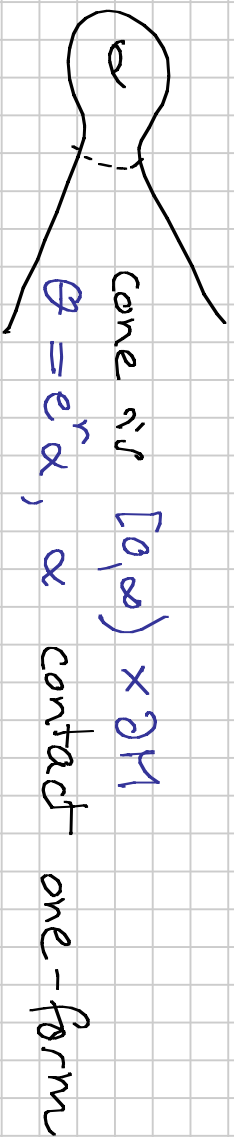


Assuming we have killed boundary contributions, this yields

$$\text{HH}^*(X, X) \rightarrow \text{End}_{\mathbb{C}}(\text{Ext}_X^*(\mathcal{E}_0, \mathcal{E}_1)[\pm 1])$$

(not strictly unique)

Symplectic topology $(M, \omega = d\theta)$ a finite type complete
 Liouville manifold



Take $H \in C^\infty(M, \mathbb{R})$, $H(r, x) = e^r$ on the cone. Symplectic
cohomology (Cieliebak-Floer-Hofer-Hyockki, Viterbo) \leftarrow used here

$$SH^*(M) = \varinjlim_{\lambda} HF^*(\lambda H)$$

We additionally assume that $c_1(M) = 0$, so that SH^* is \mathbb{Z} -graded.

Together with the standard Lagrangian Floer cohomology groups $HF^*(L_0, L_1)$ (LFCM compact, exact, spin, graded) this forms a naturally occurring example of an extended TCF. In particular, we have $1 \in SH^0(M)$ as well as $\Delta: SH^*(M) \rightarrow SH^{*-1}(M)$.

Let $SC^*(M)$ be the chain complex underlying $SH^*(M)$.

Definition A Liouville element is a cocycle $b \in SC^1(M)$ such that $B = [b] \in SH^1(M)$ satisfies $\Delta B = 1$.

We have open-closed string maps

$$\phi^0 : SC^*(M) \longrightarrow CF^*(L, L)$$

$$\phi^{\pm} : SC^*(M) \longrightarrow \text{End}(CF^*(L_0, L_1)) [-1]$$

satisfying

ϕ^0 is a chain map

ϕ^{\pm} is a chain homotopy between left and

right multiplication with ϕ^0

(in fact, all of them together form a chain map from $SC^*(M)$ to the Hochschild cochain complex of the Fukaya category).

Definition A b-equivariant Lagrangian submanifold is a pair (L, β) where $\beta \in CF^0(L, L)$ satisfies $d\beta = \phi^0(b)$.

More precisely, two choices of β are considered equivalent if they differ by something nullhomologous: so, the obstruction is $[\phi^0(L)] \in HF^1(L, L) = H^1(L; \mathbb{C})$, and the "freedom of choice" is $HF^0(L, L) = H^0(L; \mathbb{C})$. To simplify things, we will assume from now on that L is connected.

Given (L_0, β_0) and (L_1, β_1) , we can correct $\phi^1(b)$ to get an endomorphism $[\tilde{\phi}^1]$ of $HF^*(L_0, L_1)$, induced by

$$a \mapsto \phi^1(b)(a) \pm \beta_1 a \pm a \beta_0$$

This yields $HF^*(L_0, L_1) = \bigoplus_{\lambda \in \mathbb{C}} HF^*(L_0, L_1)^\lambda$

Definition

$$\log q L_1 = \text{Str} \left(e^{\log(q) [\tilde{\phi}^1]} \right) = \sum_{\lambda} q^\lambda \chi(HF^*(L_0, L_1)^\lambda).$$

Easy properties

• $L_1 \cdot q L_0 = (-1)^n q (L_0 \cdot q^{-1} L_1)$.

use $\Delta[b] = 1$

• If $H^*(L; \mathbb{C})$ has a single generator (as a ring) of degree k/n , then $L \cdot q L = 1 + (-1)^k q^{k/n} + (-1)^{2k/n} q^{2k/n} + \dots + (-1)^n q$.

• Change the equivariant structure $\beta_k \mapsto \beta_k + n_k \cdot 1_{L_k}$, denoted by $L_k \mapsto L_k \langle n_k \rangle$, yields $L_0 \langle n_0 \rangle \cdot L_1 \langle n_1 \rangle = q^{n_1 - n_0} (L_0 \cdot q L_1)$.

• Invariant under Hamiltonian (exact Lagrangian) isotopies.

Picard-Lefschetz If T_{L_1} is the Dehn twist along a Lagrangian sphere L_1 ,

$$T_{L_1} (L_0) \cdot q L_2 = L_0 \cdot q L_2 + (-1)^{n+1} q^{-1} (L_0 \cdot q L_1) (L_1 \cdot q L_2)$$

ordinary grading

equivariant structure

Note $T_L(L) = L \langle [1-n] \rangle \langle 1 \rangle$, so q -intersection numbers feel the non-triviality of T_L^2 in even dimensions.

Existence of Liouville elements (one could use more general classes in SFT as well, but these may not give rise to nontrivial q -intersection numbers)? On the negative side:

Lemma (Viterbo, Abouzaid-Schwarz) (i) If L is a $K(\pi, 1)$, then $M = T^*L$ does not admit a Liouville element.

(ii) If M admits a Liouville element, it can't contain any exact Lagrangian $K(\pi, 1)$ submanifolds.

Lemma (Bourgain-Gancs) Let $\mathcal{G}M$ be the contact boundary at ∞ of M . If M admits a Liouville element, then $[M] \in H_{2n}(M, \mathcal{G}M; \mathbb{C})$ lies in the image of the SFT map

$$CH_*^{lin}(\mathcal{G}M) \longrightarrow H_*(M, \mathcal{G}M; \mathbb{C}) [u^{-1}]$$

(The converse is probably not true)

On the positive side:

Lemma (same people) (i) T^*S^r , $r \geq 2$, has a Liouville element.

(ii) If L is so that there is a non-zero degree map $\underline{L} \times S^r \rightarrow L$, $r \geq 2$, then it has a Liouville element.

This applies for instance to $L = \mathbb{C}P^r = \text{Sym}^r(S^2)$.

To construct other examples, use any approach which (partially) computes SH^* , such as Dannea, McLean, ...

Lemma Let $\pi: M \rightarrow \mathbb{C}$ be an exact Lagrangian fibration.

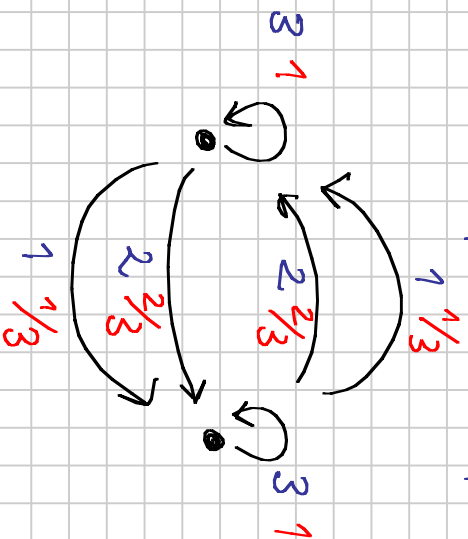
If the fibre admits a Liouville element, then so does the total space.

Example $M = \{x_1^2 + x_2^2 + x_3^2 + x_4^2 + 1 = 0\} \subset \mathbb{C}^4$, $m \geq 2$

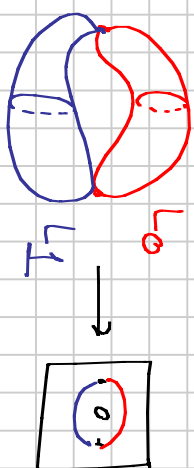
Example (This is the mirror of $\text{Tot}(\mathcal{O}(-1) \oplus \mathcal{O}(-1)) \rightarrow \mathbb{CP}^1$)
Take

$$M = \{x_1^2 + x_2^2 + x_3^2 + x_4^2 = 1, x_4 \neq 0\}$$

Draw $H\mathbb{F}^*(L; \mathbb{C})$:



numbers in blue
 are degrees
 numbers in red
 are "weights"



Such homogeneity properties have strong implications for the underlying cochain level structures:

Lemma \mathcal{A} on A_∞ -algebra. If there is $b \in HH^1(A, \mathcal{A})$ such that $[b^0] = 0 \in H^1(A)$ and which induces the Euler derivation (weight k on the degree k part) $\Rightarrow \mathcal{A}$ is formal.